## NEW RESULTS FOR NON-PERTURBATIVE **QUANTUM FIELD THEORY**



# **Carleton** University

**Department of Physics** 

PRISMA+ Colloquium, JGU Mainz

**Collaborators: Yang Bai, Hassan Easa,** Carlos de Lima, Jonathan Ponnudurai, and Cyrus Robertson Orkish.

#### DANIEL STOLARSKI

April 30, 2025



#### Grew up in Dallas, TX, USA.

Liked, chemistry, physics, and math in high school.



## ABOUT ME

## ABOUT ME

### Went to Caltech for undergrad.

Found chemistry, boring, math difficult.

Majored in physics.

Advisor: David Politzer.



## ABOUT ME

Undergrad research working on CMS for LHC.

Went to Berkeley wanting to do particle physics.

Experiment was hard, chose thoery.





Theoretical Advanced Study Institute was a great experience.



### TAS 2009

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|             |   |      |  |

#### Theoretical Advanced Study Institute in Elementary Particle Physics: Physics of the Large and the Small (TASI 2009)

| Papers<br>5<br>4<br>3<br>2<br>1<br>0<br>Inflation | - Citeable   | - Published  | 5<br>10-49   | <b>3</b><br>50-99           | 1 100-249         | 1<br>250-499 | 1<br>500+<br>Citations            | 4 contribu |
|---|--|--|--|-----------------------------|-------------------|--------------|-----------------------------------|------------|
| Daniel Bau  | mann (Princ  | eton, Inst. Advar                                  | nced Study and H   | larvard U., Phys.           | Dept.) (Jul, 2009 | )            |                                   |            |
| Contributio                                       | on to: TASI 2                                      | 009, 523-686 •                                     | e-Print: 0907.54   | 24 [hep-th]                 |                   |              |                                   |            |
| 🖾 pdf   | ℓ∂ DOI   | 🖸 cite   |  |                             | c referer         | nce search   | → 1,244 citations                 |            |
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| Tony Gherg  | hetta (Melb  | ourne U.) (Aug,                                    | 2010)  |                             |                   |              |                                   |            |
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1-26 June 2009. Boulder, CO, United States (C09-06-01.3)

Part of the TASI series

Contact: tasi@colorado.edu

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| Papers<br>5<br>4<br>3<br>2<br>1<br>0<br>Inflation | - Citeable   | - Published  | 5<br>10-49   | <b>3</b><br>50-99           | 1 100-249         | 1<br>250-499 | 1<br>500+<br>Citations            | 4 contribu |
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#### iNSPIRE:∎₽₽

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|             |   |      |  |



### POSTDOCS

First postdoc joint at University of Maryland and Johns Hopkins.

Second postdoc at CERN.



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Published in: JHEP 05 (2015) 059 Published: May 12, 2015 e-Print: 1502.05409 [hep-ph] DOI: 10.1007/JHEP05(2015)059 Report number: CERN-PH-TH-2015-031, DESY-15-026 View in: ADS Abstract Service, CERN Document Server

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Mainz

To appear



## PREVIJS V ST

#### DANIEL STOLARSKI WITH PEDRO SCHWALLER AND ANDREAS WEILER

December 16, 2014

### Faculty at Carleton University.





### **QUANTUM FIELD THEORY**

Combination of quantum mechanics and special relativity.

All objects represented as fields (e.g. electric and magnetic field).

Particles are excitations in the field.



**David Tong** 



### PERTURBATIVE QET

#### QFT is hard.

Most successful tool is perturbation theory.

Terms in series represented by Feynman diagrams.

Electron scattering:

 $\mathcal{M} \sim e^2$ 

 $\sigma_{\rm Born} \sim e^4$ 





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#### Electron magnetic moment:

$$g = 2$$
 Dirac, 1928

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#### Electron magnetic moment:





Schwinger, 1948

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Electron magnetic moment:





Electron magnetic moment:

g = 2**Dirac**, 1928  $g = 2\left(1 + \frac{\alpha}{2\pi}\right)$  Schwinger, 1948

 $g_{\rm th} = 2(1 + 0.0011596521816)$ Kinoshita et al, 2014

### $(i\gamma^{\mu}\partial_{\mu} - m)\psi(x) = 0$



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#### $g_{\rm exp} = 2(1 + 0.0011596521806)$ Fan et. al. 2023

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### $(i\gamma^{\mu}\partial_{\mu} - m)\psi(x) = 0$



## RUNNING COUPLINGS

Quantum electrodynamics characterized by charge of electron: *e*.

Often use dimensionless fine structure constant:  $\alpha = \frac{e^2}{4\pi} \frac{1}{\epsilon_0 \hbar c} \approx \frac{1}{137}$ 

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Fun fact: the fine structure constant is not a constant, it "runs" with energy.

 $\frac{d\alpha}{EE} \approx \frac{2\alpha^2}{3\pi}$ 

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 $3\pi$ 







### NON-PERTURBATIVE QFT

QED (electromangetic force) expansion parameter is

$$\frac{e^2}{4\pi} \equiv \alpha \approx \frac{1}{137} \ll 1$$

What about theories where expansion parameter is ~ 1?

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## Strong force (QCD) is such a theory!







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What about theories where expansion parameter is ~ 1?

## Strong force (QCD) is such a theory!



#### Feynman diagrams are unhelpful.

Most of our understanding of strong force is from data.



Things we know about the strong force from data:

 Confinement: do not see free quarks, only see heavy bound states.



#### Millennium Problems

Yang-Mills and Mass Gap Riemann Hypothesis P vs NP Problem Navier-Stokes Equation Hodge Conjecture <del>Poincaré Conjecture</del> Birch and Swinnerton-Dyer Conjecture

A Martin State Constant Alter and a state



π

Things we know about the strong force from data:

 Spontaneous symmetry breaking: symmetry of bound states different than that of physical laws.



Two things are actually deeply connected.

 Confinement: do not see free quarks, only see heavy bound states.





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### RUNNING STRONG COUPLING

Like electromagnetism, strong coupling  $\alpha_s$  changes with energy.



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Like electromagnetism, strong coupling  $\alpha_s$  changes with energy.



Coupling gets weaker at high energy: asymptotic freedom.

 $g_0$ 

### VERFED BY DATA

Like electromagnetism, strong coupling  $\alpha_s$  changes with energy.



Coupling gets weaker at high energy: 0.05 asymptotic freedom.



### TRPTO STOCKHOLM

EDISH ACADEMY OF SCIENCES HAS DECIDED TO AW PIZE IN PHYSICS FOR 2004 "FOR THE DISCOVERY OF A DOM IN THE THEORY OF THE STRONG INTERACTION" JO GROSS, H. DAVID POLITZER AND FRANK WILCZEK

#### A colourful connection

The scientists awarded this year's Nobel Prize in Physics have solved a mystery surrounding the strongest of nature's four fundamental forces. The three quarks within the proton can sometimes appear to be free, although no free quarks have ever been observed. The quarks have a quantum mechanical property called colour and interact with each other through the exchange of gluons – nature's glue.

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FUCHER READING Information on the Nobel Prize in Physics 2004: www.nobelprize.org CERN: www.cern.ch Hands-on-CERN: http://hands-on-cern.physto.se/ DESY: www.desy.de
 The particle adventure: http://particleadventure.org/particleadventure/
 QCD Made Simple, Physics Today August 2000, p. 22
 Joining up the dots with the

strong orce, by C. Davies, CERN Courler June 2004, p. 23 The W and Z at LEP, by C. Sutton and F. Zerwas, CERN Courier May 2004, p. 21 Lattice Quantum Chro Age, by C. De Tar and S. Gottlieb, Physics Today February 2004, p. 45 In search of the ultimate quilding blocks, by G. 't Hooft, Cambridge University Press 1997

5



#### **The Standard Model** and the four forces

The quarks and gluons of the strong (or colour) force are the third piece in the puzzle of nature's four forces. The first piece, the electromagnetic force, is similar to the strong force but instead of gluons, particles of light, photons, are the force carriers. The gluons carry colour charge while the photons are electrically neutral. The second piece in the puzzle is the weak force, which controls some radioactive decays and energy production in the sun. This force differs from the 0 other two because the force----carrying particles are very heavy. The fourth force, gravi-ty, is the least understood even though it is experienced by us

all. Gravitons are thought to be the force-carrying particles, but they have yet to be discovered. The Standard Model provides a description of all the forces apart from gravity.

dynamics Comes of

If a quark is knocked out of the proton in a high-energy colli-

sion, it appears to behave as a free particle for an instant.



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1949 Hister Yukawa, The theory of nuclear forces 1957 CHEN Ning Yang and violation in particle physics 1965 Sin-itieo Tomonaga, Julian violation in particle physics 1965 Sin-itieo Tomonaga, Julian Privman, QED - the quantum theory of electromagnetic interactions 1969 Symmetry properties of electromag Editors: Lars Bergström and Per Carlson, Secretary and Member of O The Royal Swedish Academy of Sciences 
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 the Nobel Committee for Physics. Mark Pearce. The Royal Institute of Technology, Stockholm, Jonas Foran, Anna Lindquist and Eva Krutmeijer. The Royal Swedish Academy of Sciences
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EARLIER NOBEL LAUREATES WHOSE WORK WAS OF GREAT CONSEQUENCE FOR THIS YEAR'S AWARDS

Inside the proton

The three guarks within the proton are held together by the powerful

force mediated by the gluons, depicted here as coiled springs. As the distance between the quarks increases, so does the force between them



David J. Gross ute for Theoretical Physics. California | of California, Santa Barbara, USA (Caltech), i

H. David Politzer California Institute of Technology dena, USA

> HE THEORY SHOWS ITS TRUE COLOURS The aftermath of a high-energy collision between a proton and an electron, as seen by the H1 experiment at the DESY laboratory in Hamburg. The experiment is shown in cross-section, perpendicular to col-

### THE NOBEL PRIZE IN PHYSICS 2004





Frank Wilczek Massachusetts Institute of Technology (MIT), Cambridge, USA



Many tried, but failed, to find a theory in which the strength of the strong force decreases as the energy increases. This year's Nobel laureates found a theory with the required minus sign.

$$\beta(g) = \frac{g^3}{16\pi^2} \left( \frac{11}{3} N_c - \frac{4}{3} \frac{N_F}{2} \right)$$



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1990 JEROME L. FRIEDMAN, HENRY W. KENDALL AND RICKARD E. TAYLOR. The discovery of quarks through electron-scattering experiments The quantum structure of the electro-weak interaction

Layout and illustrations: Typoform Printing: Billes Tryckeri 2005 priered free of charge by phone, fax

liding beams of protons and electrons. The electron has struck one of the quarks in a proton. An impressive shower of particles – providing

information about the struck quark - is spontaneously produced from the energy stored in the gluon force-field. The charged particles in the shower bend in the experiment's strong magnetic field.

When the quarks are very close to each other, i.e. when the distance between them is *asymptotically* approaching zero, the force is so weak that they behave almost as free particles.

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1982 KENNETH G WILSON. The theory of phase Editors: Lars Bergström and Per Carison, Secretary and Member or O The Royal Swedish Academy of Sciences. the Nobel Committee for Physics. Mark Pearce. The Royal Institute:
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#### A unified theory for all forces?

This year's prize paves the way for a more fundamental future description of the forces in nature. The electromagnetic, weak and strong forces have much in common and are perhaps different aspects of a single force. They also appear to have the same strength at very high energies, especially if 'supersym metric' particles exist. It may even be possible to include gravity if theories which treat matter as small vibrating strings are correct.

KUNGL VETENSKAPSAKADEMIEN THE ROYAL SWEDISH ACADEMY OF SCIENCES

1990 JEROME L. FRIEDMAN, HENRY W. KENDALL AND RICHARD E. Taylor, The discovery of quarks through electron-scattering experiments

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# **SUPERSYMMETRY**

## SPHERICAL COW

Supersymmetric theories are the spherical cows of quantum field theory.

Double number of particles.

SUSY version of many QFTs can be solved. Seiberg, mostly.

Deform away from SUSY.



### Keenan Crane, Wikipedia.



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## SYMETRY AND ENERGY

Does a spin up electron have the same energy as a spin down electron?

 $? < \uparrow |H| \uparrow > = < \downarrow |H| \downarrow >$ 

## SYMMETRY AND ENERGY

Does a spin up electron have the same energy as a spin down electron?

## $<\uparrow |H|\uparrow>=<\downarrow |H|\downarrow>$

Of course it does. If I just rotate the physicist, I turn  $|\uparrow\rangle$  into  $|\downarrow\rangle$  .

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Laws of physics are invariant under rotation.

## SYMMETRY AND ENERGY

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Laws of physics are invariant under rotation.

Can separate the energy with a magnetic field, but that would be an explicit breaking of rotational symmetry.

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# SUPERSYMMETRY AND ENERGY

- Supersymmetry relates bosons to fermions.
  - Q|fermion > = |boson >
  - Q|boson > = |fermion >
- In supersymmetric theories, for every fermion, there is a boson with the exact same mass (energy) and charge.

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- Supersymmetric electromagnetisr

|      |      | spin | mass    |
|------|------|------|---------|
| n.   | e    | 1/2  | 511 keV |
| •••• | se   | 0    | 511 keV |
|      | 8    | 1    | 0       |
|      | yino | 1/2  | 0       |

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Not our universe!

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## BREAKING SUPERSYMMETRY

To break supersymmetry, give mass to some spins and not others.

Electromagnetism with broken supersymmetry.

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## BREAKING SUPERSYMMETRY

To break supersymmetry, give mass to some spins and not others.

Electromagnetism with broken supersymmetry.

Could be our universe!



|      | spin | mass    |
|------|------|---------|
| e    | 1/2  | 511 keV |
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| X    | 1    | 0       |
| yino | 1/2  | 1 TeV?  |

## ANOMALY MEDIATED SUSY BREAKING

In general, new mass parameter for every particle.

A simple mechanism to break SUSY is called anomaly mediation (AMSB). Randall, Sundrum, hep-th/9810155. Giudice, Luty, Murayama, Rattazzi, hep-ph/9810442.

All new masses controlled by one parameter:  $m_{3/2}$ .

Breaking mechanism is extremely predictive!

 $\mathcal{X} = -rac{1}{3!}\left(\gamma_i+\gamma_j+\gamma_k
ight)y^{ijk}\phi_i\phi_j\phi_k,$  $\gamma_i^j = -rac{1}{32\pi^2} [y_{ikl}^* y^{jkl} - 4g^2 \delta_i^j C_A(\phi_i)],$ 

 $A_{ijk} = -(\gamma_i + \gamma_j + \gamma_k)y^{ijk}m_{3/2}.$ 

 $L=-rac{1}{2}\dot{\gamma}^i_i\phi_i^+\phi_i,$ 

 $m_i^2 = \frac{1}{\Lambda} \dot{\gamma}_i^i m_{3/2}^2.$ 

 $\dot{\gamma}_i^j = \frac{\partial \gamma_i^j}{\partial \ln \mu}.$ 



## ASIDE AMSB FOR SOED

Consider supersymmetric electromagnetism coupled to AMSB:

$$m_{se}^2 \propto -\frac{d\alpha}{dE}$$

 $\frac{d\alpha}{dE} > 0 \text{ for electromagnetism.}$ 

Partner of the electron becomes a tachyon and gives mass to the photon!

SSM + AMSB does not describe our universe.



## ANSB -- NON-PERTURBATIVE OF

V > hep-th > arXiv:2104.01179

### **High Energy Physics – Theory**

[Submitted on 2 Apr 2021 (v1), last revised 19 Jun 2021 (this version, v3)] Some Exact Results in QCD-like Theories

### Hitoshi Murayama

may be in the same universality class.

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I propose a controlled approximation to QCD-like theories with massless quarks by employing supersymmetric QCD perturbed by anomaly-mediated supersymmetry breaking. They have identical massless particle contents. Thanks to the ultravioletinsensitivity of anomaly mediation, dynamics can be worked out exactly when  $m \ll \Lambda$ , where *m* is the size of supersymmetry breaking and  $\Lambda$  the dynamical scale of the gauge theory. I demonstrate that chiral symmetry is dynamically broken for  $N_f \leq \frac{3}{2}N_c$  while the theories lead to non-trivial infrared fixed points for larger number of flavors. While there may be a phase transition as m is increased beyond  $\Lambda$ , qualitative agreements with expectations in QCD are encouraging and suggest that two limits  $m \ll \Lambda$  and  $m \gg \Lambda$ 

## ANSB -- NON-PERTURBATIVE OF

V > hep-th > arXiv:2104.01179

### **High Energy Physics – Theory**

[Submitted on 2 Apr 2021 (v1), last revised 19 Jun 2021 (this version, v3)]

### Hitoshi Murayama

may be in the same universality class.

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## Some Exact Results in QCD-like Theories

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[Submitted on 2 Apr 2021 (v1), last revised 19 Jun 2021 (this version, v3)]

### Some Exact Results in QCD-like Theories

Hitoshi Murayama

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### **High Energy Physics – Theory**

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### Some Exact Results in Chiral Gauge Theories

Csaba Csáki, Hitoshi Murayama, Ofri Telem

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### **High Energy Physics – Theory**

[Submitted on 7 May 2021]

### More Exact Results on Chiral Gauge Theories: the Case of the Symmetric Tensor

Csaba Csáki, Hitoshi Murayama, Ofri Telem





### **High Energy Physics – Theory**

### [Submitted on 18 Jun 2021 (v1), last revised 9 Sep 2021 (this version, v2)] **Demonstration of Confinement and Chiral Symmetry Breaking in** $SO(N_c)$ **Gauge Theories**

Csaba Csáki, Andrew Gomes, Hitoshi Murayama, Ofri Telem





[Submitted on 6 Jul 2021 (v1), last revised 9 Sep 2021 (this version, v2)]

### The Phases of Non-supersymmetric Gauge Theories: the $SO(N_c)$ Case Study

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### **Dynamics of Simplest Chiral Gauge Theories**

netry



[Submitted on 6 Jul 2021 (v1), last revised 9 Sep 2021 (this version, v2)]

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[Submitted on 6 Jul 2021 (v1), last revised 9 Sep 2021 (this version, v2)]

### **High Energy Physics – Phenomenology** The Phases of Non-supersymmetric [Submitted on 8 Jul 2024 (v1), last revised 16 Dec 2024 (this version, v2)] the $SO(N_c)$ Case Study Spontaneous CP Breaking in a QCD-like Theory

Csaba Csáki, Andrew Gomes, Hitoshi Murayama, Ofri To Csaba Csáki, Maximilian Ruhdorfer, Taewook Youn

Csaba Csáki, Andrew Gomes, Hitoshi Murayama, Bea Noether, Digvijay Roy Varier, Ofri



## SMALL VS. LARGE SUSY BREAKING

Two dimensionful parameters:

- $\Lambda$  mass of bound states
- $m_{3/2}$  size of SUSY breaking



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# GENERALZED (CD)

Quantum Chromodynamics (QCD) described by  $N_c = 3$  colours and  $N_f = 3$  light flavours. Promote  $N_c$  and  $N_f$  to variables.  $N_f = N_c + 1$  has a particularly nice SUSY

description.





# $V_{f} = 2 \text{ AND } N_{f} = 3$

QCD with 2 colours is different. Can calculate potential for Meson field:  $V = |M|^{4} + \frac{27}{1024\pi^{4}} |M|^{2} - \frac{9}{32\pi^{2}} M^{3}$ 

This potential has a minimum at M = 0.

Theory does not have spontaneous symmetry breaking!









Hassan Easa PhD thesis, Csaki et. al., arXiv:2212.03260. de Lima, DS, arXiv:2307.13154.



# **THOOFTANOMALY MATCHING**

# 

Most robust tool to analyze strong interactions is 't Hooft anomaly matching. 't Hooft '80.



Weakly coupled high energy theory

### Strong dynamics



# 

Most robust tool to analyze strong interactions is 't Hooft anomaly matching. 't Hooft '80.

If theory at very high energy is tractable (perturbative), get consistency condition for low energy spectrum.



Weakly coupled high energy theory



Weakly coupled low energy theory



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Most robust tool to analyze strong interactions is 't Hooft anomaly matching. 't Hooft '80.

If theory at very high energy is tractable (perturbative), get consistency condition for low energy spectrum.

NB: 't Hooft anomaly matching  $\neq$ anomaly mediated SUSY breaking



Weakly coupled high energy theory



Weakly coupled low energy theory



## ANOTHER THEORY

Simplest Grand Unified Theory of the SM is SU(5) GUT. Georgi, Glashow, PRL'74.

Can describe all (gauge) forces and all matter of SM.

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What if we do not break it to SM, and just let it evolve?





# ANOTHER THEORY

Simplest Grand Unified Theory of the SM is SU(5) GUT. Georgi, Glashow, PRL'74.

Can describe all (gauge) forces and all matter of SM.

What if we do not break it to SM, and just let it evolve?  $\frac{\Lambda_{\rm QCD}}{43} = e^{-\frac{8\pi^2}{g_0^2 b}}, \ \Lambda_{\rm QCD} \leq \frac{43}{43} \frac{\alpha^2}{\alpha^2} < 0 \ b = 7$ 





## THOOFT ANOMALY MATCHING

Compute anomalies at high energy.

 $A \left[ SU(3)_{A}^{3} \right] = 10$  $A \left| SU(3)_{\bar{F}}^3 \right| = 5$  $A\left[\operatorname{grav}^2 \times U(1)_B\right] = -15$  $A\left[U(1)_{B}^{3}\right] = -375$  $A \left[ SU(3)_{A}^{2} \times U(1)_{B} \right] = 10$  $A \left| SU(3)_F^2 \times U(1)_B \right| = -15$ 

## THOOFT ANOMALY MATCHING

Compute anomalies at high energy.

Can use 't Hooft anomaly matching to determine light bound states (baryons) of the theory.

Boils down to solving linear equations over integers.

 $A \left[ SU(3)_{A}^{3} \right] = 10$  $A \left| SU(3)_{\bar{F}}^3 \right| = 5$  $A\left[\operatorname{grav}^2 \times U(1)_B\right] = -15$  $A\left[U(1)_{R}^{3}\right] = -375$  $A \left[ SU(3)_{A}^{2} \times U(1)_{B} \right] = 10$  $A \left| SU(3)_F^2 \times U(1)_B \right| = -15$ 

# **RANOMALY MATCHING**

### Solutions are quite complicated.

A relatively simple example:

Massless baryons

Bai, DS, arXiv:2111.11214.

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|   | [SU(5)]        | $SU(3)_A$ | $SU(3)_{\overline{F}}$ | U(1) |
|---|----------------|-----------|------------------------|------|
|   | 10             | 3         | 1                      | 1    |
| B<br>F  | $\overline{5}$ | 1         | 3                      | -3   |
|   | [SU(5)]        | $SU(3)_A$ | $SU(3)_{\overline{F}}$ | U(1) |
| $(A\overline{F}\overline{F})^{\dagger}$             | 1              | 3         | 3                      | 5    |
| $A\overline{F}\overline{F}$                         | 1              | 3         | 6                      | -5   |
| $A^5$   | 1              | 6         | 1                      | 5    |
| $\overline{F}^5$                                    | 1              | 1         | 15                     | -15  |
| $A^3 \overline{F}^{\dagger 4}$                      | 1              | 1         | 6                      | 15   |
| $A^3 \overline{F}^{\dagger 4}$                      | 1              | 1         | 15                     | 15   |
| $2 \times (A^3 \overline{F}^{\dagger 4})^{\dagger}$ | 1              | 1         | 3                      | -15  |
|   |                |           |                        |      |

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### IR ANOMALY MATCHING

#### Another example:



#### Bai, DS, arXiv:2111.11214.

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|                               |                           | [SU(5)] | $SU(3)_A$         | $SU(3)_{\overline{F}}$ | $U(1)_I$ |
|-------------------------------|---------------------------|---------|-------------------|------------------------|----------|
|                               | A                         | 10      | 3                 | 1                      | 1        |
|                               | $\overline{F}$            | 5       | 1                 | 3                      | -3       |
|                               |                           | [SU(5)] | SU(3)             | SU(3)                  | II(1)    |
|                               |                           |         | $DU(\mathbf{O})A$ | $DU(0)_F$              |          |
| $(\overline{F}^5)^\dagger$    |                           | 1       | 1                 | 3                      | 15       |
| $^{5}$ or $(A^{4}]$           | $\overline{F}^3)^\dagger$ | 1       | 6                 | 1                      | 5        |
| $\overline{F}(A^2)^{\cdot}$   | †                         | 1       | 3                 | 3                      | -5       |
| $(A^3)^{\dagger}\overline{F}$ | 4                         | 1       | 1                 | 3                      | -15      |
|                               |                           | 1       |                   |                        | •        |



# IR ANOMALY MATCHING

Another example:

# Looks simpler, but $\left(\bar{F}^{5} ight)^{\dagger}$ state is problematic.



Bai, DS, arXiv:2111.11214. 39 DANIEL STOLARSKI April 30, 2025 PRISMA+

| Г                             |                               |                                    |           |                        |             |
|-------------------------------|-------------------------------|------------------------------------|-----------|------------------------|-------------|
|                               |                               | [SU(5)]                            | $SU(3)_A$ | $SU(3)_{\overline{F}}$ | $U(1)_{I}$  |
|                               | A                             | 10                                 | 3         | 1                      | 1           |
|                               | $\overline{F}$                | 5                                  | 1         | 3                      | -3          |
|                               |                               |                                    |           |                        |             |
|                               |                               | $\left\lfloor SU(5) \right\rfloor$ | $SU(3)_A$ | $SU(3)_{\overline{F}}$ | $\mid U(1)$ |
| $(\overline{F}^5)^\dagger$    |                               | 1                                  | 1         | 3                      | 15          |
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# IR ANOMALY MATCHING

A

Another example:

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#### Bai, DS, arXiv:2111.11214.

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|                               | $\overline{F}$                | 5                                  | 1         | 3                      | -3          |
|                               |                               |                                    |           |                        |             |
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| $(A^3)^{\dagger}\overline{F}$ | ·4                        | 1       | 1           | 3                      | -15  |
|                               |                           |         | •           | -                      | •    |

Can work with orbital angular momentum, but weird.



#### WHAT HAPPENS?

If there are no solutions to anomaly matching conditions, theory **must** exhibit spontaneous symmetry breaking.

#### WHATHAPPENS2

exhibit spontaneous symmetry breaking.

This theory probably exhibits spontaneous symmetry breaking. Difficulty of 't Hooft anomaly matching is one piece of evidence.

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Breaking pattern:

 $SU(3)_L \times SU(3)$ 

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$$)_R \rightarrow SO(3)_V$$

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Breaking pattern:

 $SU(3)_L \times SU(3)$ 

Different than QCD:  $SU(3)_L \times SU(3)_L$ 

- If there are no solutions to anomaly matching conditions, theory **must**

$$)_R \rightarrow SO(3)_V$$
  
 $(3)_R \rightarrow SU(3)_V$ 

#### SUPERSYMMETRIZE THE THEORY

This theory also has a nice supersymmetric description. Seiberg, hep-th/9402044, hep-th/9411149. Csaki, Schmaltz, Skiba, hep-th/9610139.

Can compute the scalar potential:

$$V_{\text{susy}} + V_{\text{susy}} = \left| \frac{dW_{\lambda}}{dM^{ai}} + \frac{dW_{\zeta}}{dM^{ai}} + \frac{1}{3}A_1^*M + \frac{1$$

#### Bai, DS, arXiv:2111.11214.

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### SIPERSYMMETR/FILETHETHEORY/

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Can draw qualitative phase diagram in theory space.

Have some control at the corners.

Must have a phase transition in going to large SUSY breaking.

Bai, DS, arXiv:2111.11214.

# PHASEDJAGRAM



 $m_{\widetilde{q}}/\Lambda$ 



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Bai, DS, arXiv:2111.11214.

# PHASEDJAGRAM



 $m_{\widetilde{q}}/\Lambda$ 



# LGHTNING ROUND

## DIFFERENT NUMBER OF FLAVOURS?

Analyzed SU(5) model with 3 flavours.

1 flavour dynamics are well known. Dimopoulos, Raby, Susskind, NPB '80.

What about 2 flavours?

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## DIFFERENT NUMBER OF FLAVOURS?

Analyzed SU(5) model with 3 flavours.

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What about 2 flavours?

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#### Analysis of the 2 Flavour SU(5)

#### Georgi–Glashow Model

by

#### Jonathan Ponnudurai

A thesis submitted to the Faculty of Graduate and

Postdoctoral Affairs in partial fulfillment of the

requirements for the degree of

**Master of Science** 

in

Physics

### 2 FLAVOUR RESULTS

#### Non-SUSY model: anomaly matching solutions

| Symmetries  | UV Anomaly | $A^{5}_{4,1,5}$ | $A^4\overline{F}^3_{3,2,-5}$ |
|---|------------|-----------------|------------------------------|
| $\mathbf{SU}(2)^2_{\mathbf{A}} \otimes \mathbf{U}(1)_{\mathbf{B}}$            | 5          | 25              | -20                          |
| $\mathbf{SU}(2)^2_{\overline{\mathbf{F}}} \otimes \mathbf{U}(1)_{\mathbf{B}}$ | -15/2      | 0               | -15/2                        |
| $U(1)_{B}^{3}$  | -250       | 500             | -750                         |
| ${f grav.}^{2}\otimes {f U(1)_{B}}$   | -10        | 20              | -30                          |
| $SU(2)^3_A$   | 0          | 0               | 0                            |
| $SU(2)\frac{3}{F}$  | 1          | 0               | 1                            |

#### J. Ponnudurai MSc thesis.

# SUSY model: proof of spontaneous symmetry breaking.



#### What is the phase diagram of QCD?



Conjectured QCD phase diagram with boundaries that define various states of QCD matter based on  $S \chi B$  patterns.

# PHASED AGRAM

#### Longstanding difficult problem with only partial results.

#### Focus on T = 0 for now.

**Fukushima and Hatsuda**, '11.

#### Steps:

- Supersymmetrize QCD 1.
- 2. Break SUSY with AMSB
- 3. Turn on baryon density

0.2 Find some new phases that do not exist in the literature! 0.0

# PHASED AGRAM

#### Preliminary



Work in progress with Bai and de Lima.



### HOLDGRAPHY2

AdS/CFT correspondence says certain 4 dimensional theories are dual to 5 dimensional theories.

Maldacena, hep-th/9711200. Witten, hep-th/9802150.

Very supersymmetric theories fall in this category.

What is gravitational dual of Anomaly mediation?

Work in progress with Cyrus Robertson Orkish.





#### SUMMARY

Non-perturbative quantum field theories pose interesting and important open problems.

Supersymmetry + anomaly mediation gives useful new tool to analyze these theories.

Found two theories that violate conjecture that this tool can be used to analyze original theory.

Various ongoing directions of this research.



