Mesogenesis

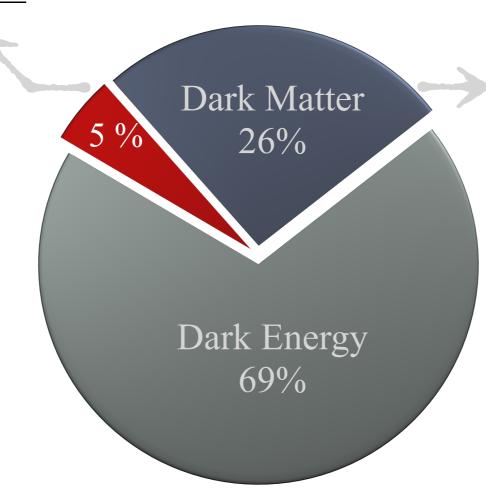
Gilly Elor

MITP Particle Theory Seminar June 30 2021

What is the Universe Made of?

From cosmological measurements we know:

The stuff we understand stars, planets, you (baryonic matter) Only 5 %

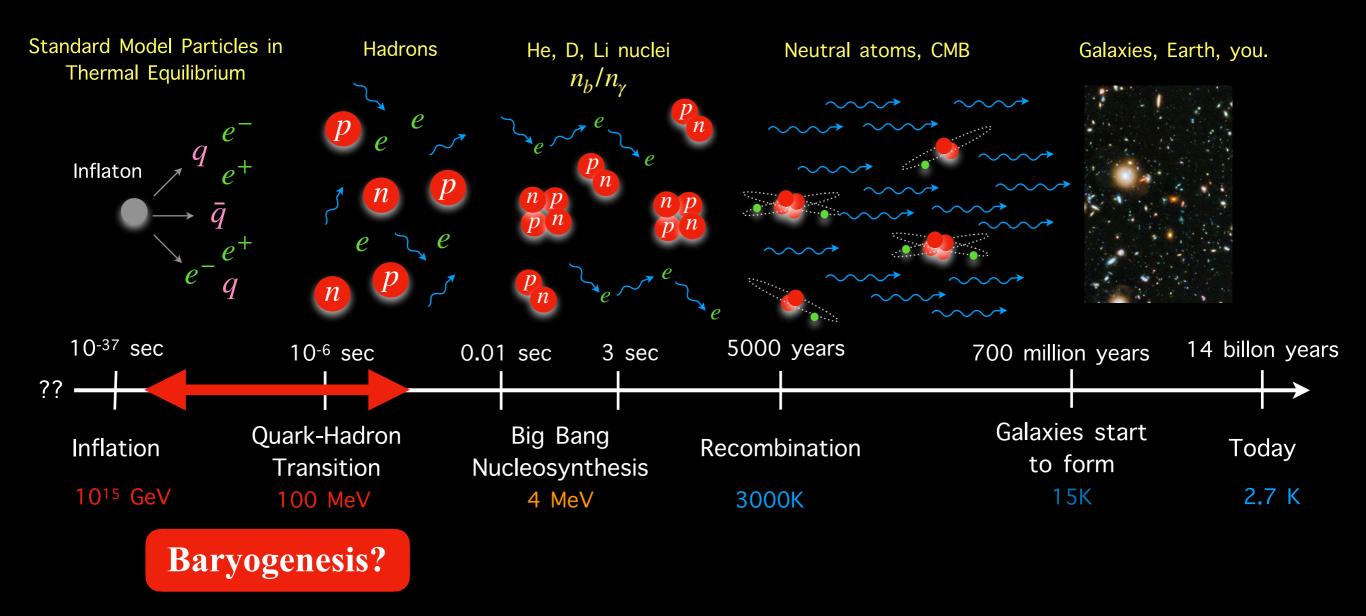


What is the nature and origin of dark matter?

Energy density today

We don't know where the 5 % of baryonic matter came from.

The History of the Universe

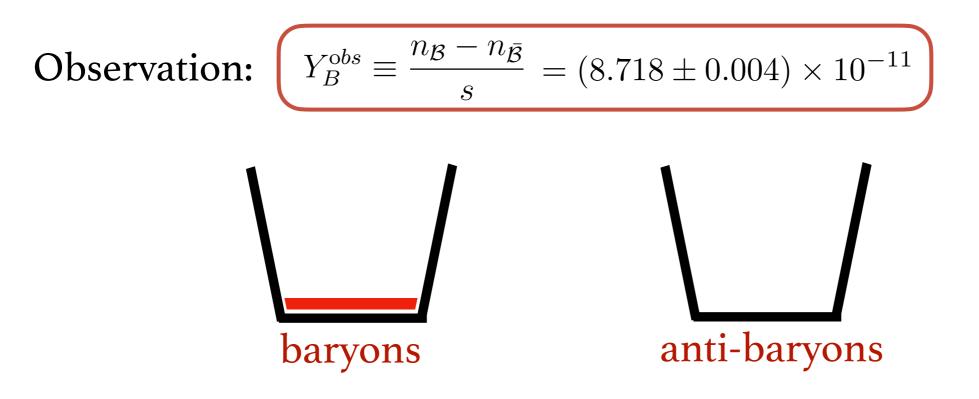


What mechanism generated the initial asymmetry? Observed to be (BBN, CMB):

$$Y_B^{\text{obs}} \equiv \frac{n_{\mathcal{B}} - n_{\bar{\mathcal{B}}}}{s} \sim 8 \times 10^{-11}$$

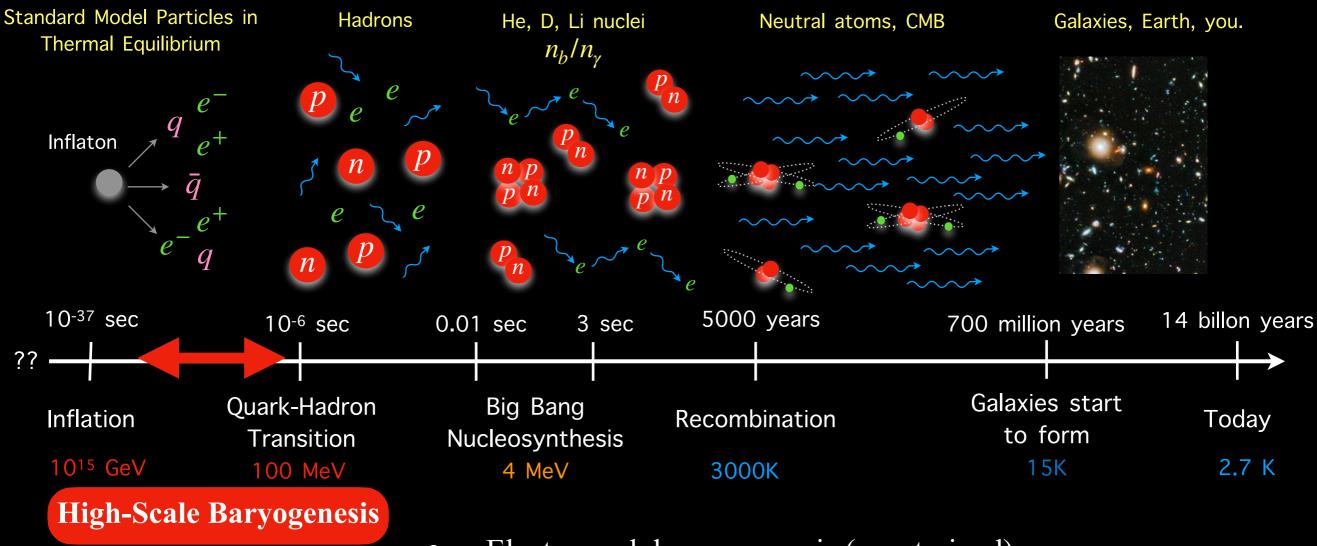
Baryogenesis

Review: how to generate a baryon asymmetry?



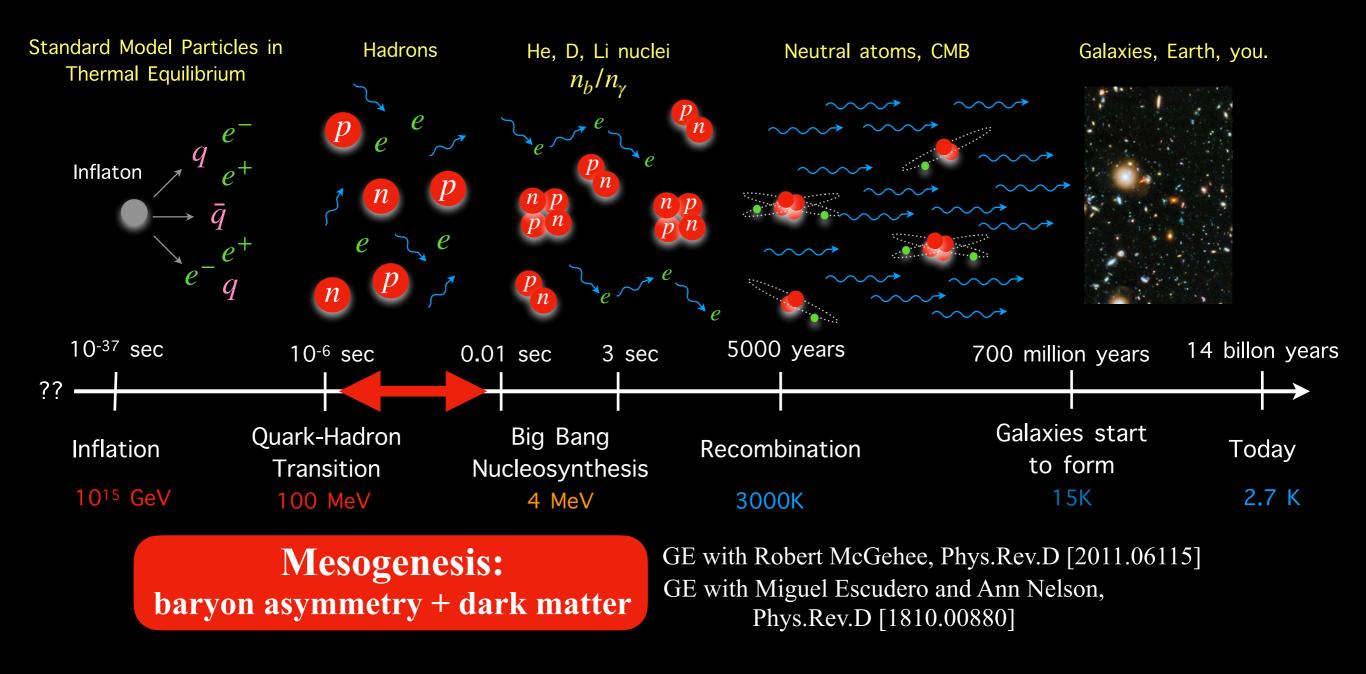
- Baryon number violation.
- Conjugate rates must be different.
- Out of thermal equilibrium.

"Traditional" Baryogenesis



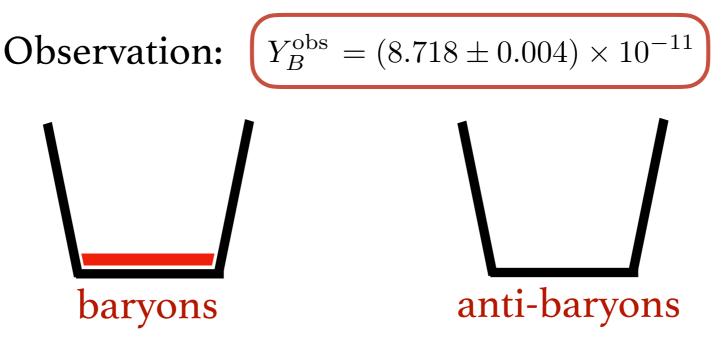
- Electroweak baryogengesis (constrained)
- Leptogenesis (hard to test)
- Affleck-Dine (very hard to test)
- • • • •

Making the Universe at 20MeV



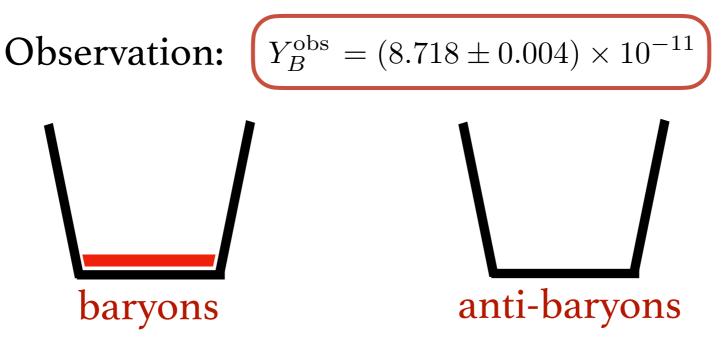
- Controlled by experimental observables. Signals!
- Theoretically appealing e.g. Relaxion and Nnaturalness require low scale baryogenesis.





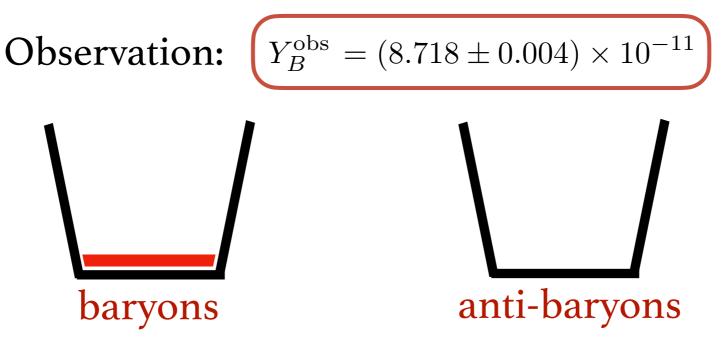
- Out of thermal equilibrium:
- CP Violation:
- "Baryon number violation":





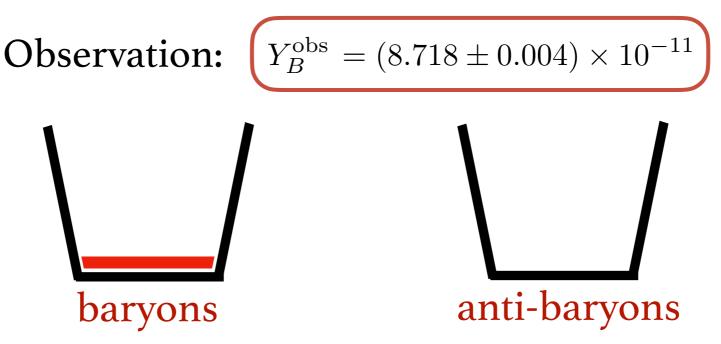
- Out of thermal equilibrium: Late decays of "inflaton" field to SM Mesons.
- CP Violation:
- "Baryon number violation":





- Out of thermal equilibrium: Late decays of "inflaton" field to SM Mesons.
- CP Violation: In SM Meson systems.
- "Baryon number violation":





- Out of thermal equilibrium: Late decays of "inflaton" field to SM Mesons.
- CP Violation: In SM Meson systems.
- "Baryon number violation": SM Meson decays to dark leptons or baryons.

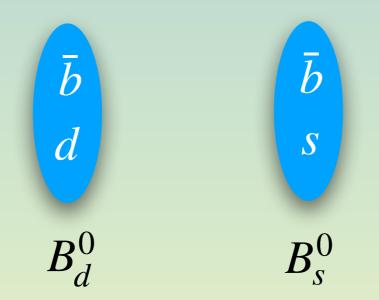
Today's Talk

Baryogenesis and Dark Matter production using the CP Violation of Standard Model Meson Systems

Observation: $Y_B^{\text{obs}} = (8.718 \pm 0.004) \times 10^{-11}$

- Part I. *B*-Mesogenesis: CP violation from neutral B meson oscillations [GE with Miguel Escudero and Ann Nelson arXiv:1810.00880].
- Part II. Discovering *B*-Mesogenesis: Details of signals and searches [GE with Gonzalo Alonso-Alvarez and Miguel Escudero arXiv:2101.02706, ongoing work, both theory and experimental].
- Part III. *D*-Mesogenesis: CP violation from charged *D* meson decays [GE with Robert McGehee arXiv:2011.06115].

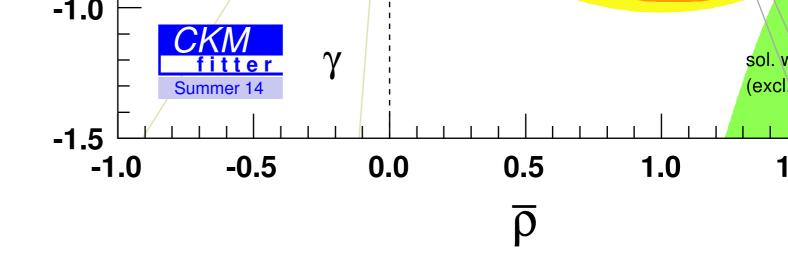
Part I. B-Mesogenesis



Based on:

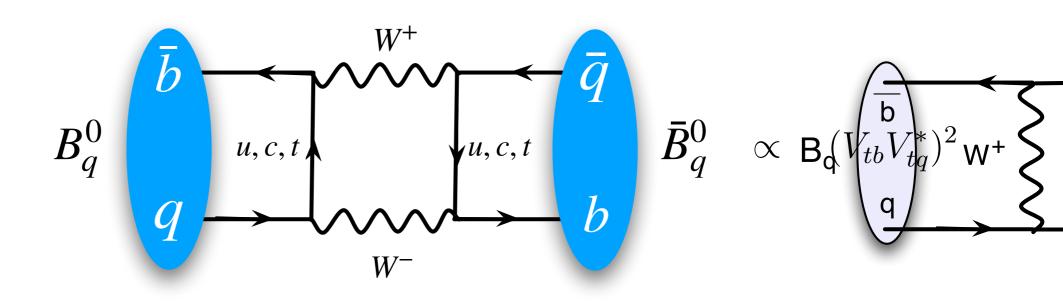
[GE with Miguel Escudero and Ann Nelson, Phys. Rev. D, arXiv:1810.00880] [GE with Gonzalo Alonso-Alvarez, Ann Nelson and Huangyu Xiao, JHEP, arXiv:1907.10612]

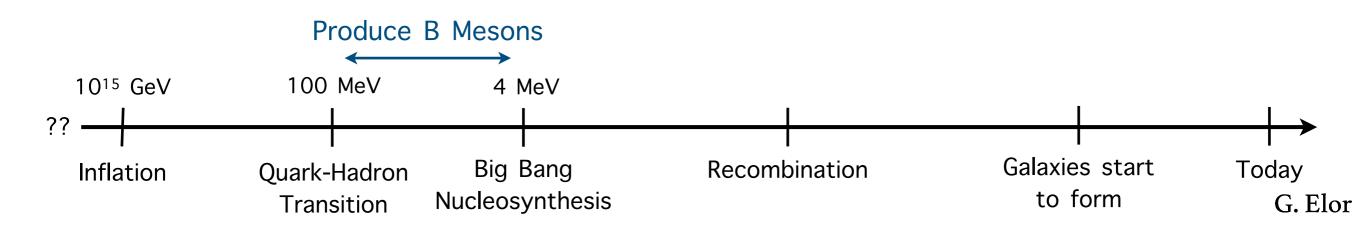
Neutral B]





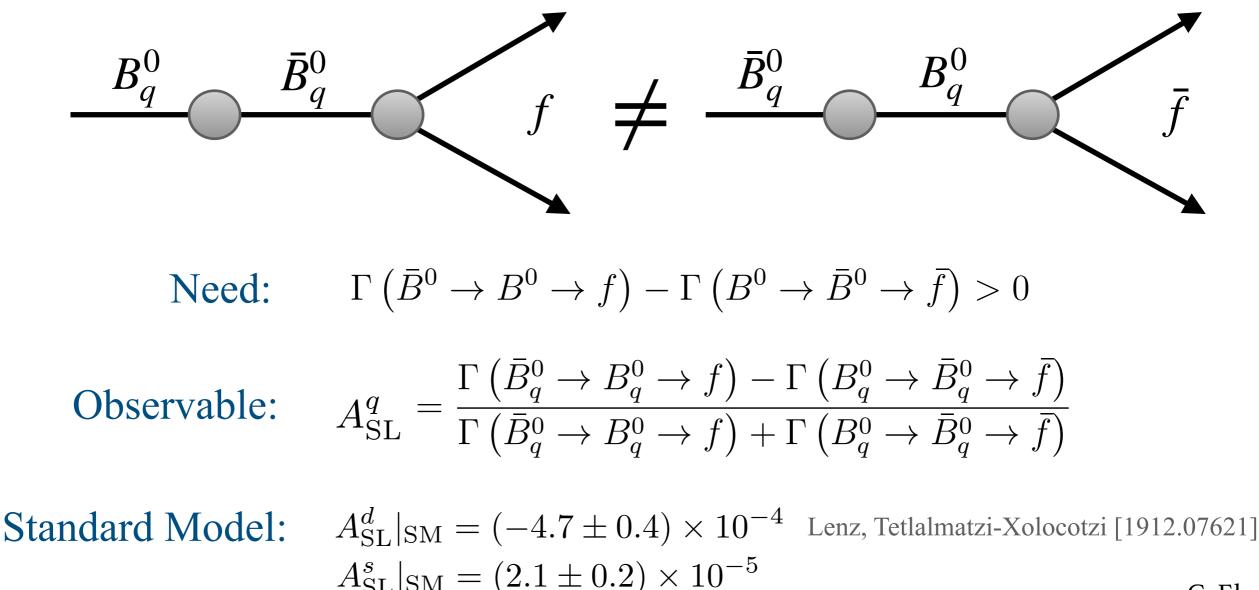
At low energies we can use CPV in *B* meson mixing e.g. from CKM phases in the case of the Standard Model (but new physics contributions are also not excluded)





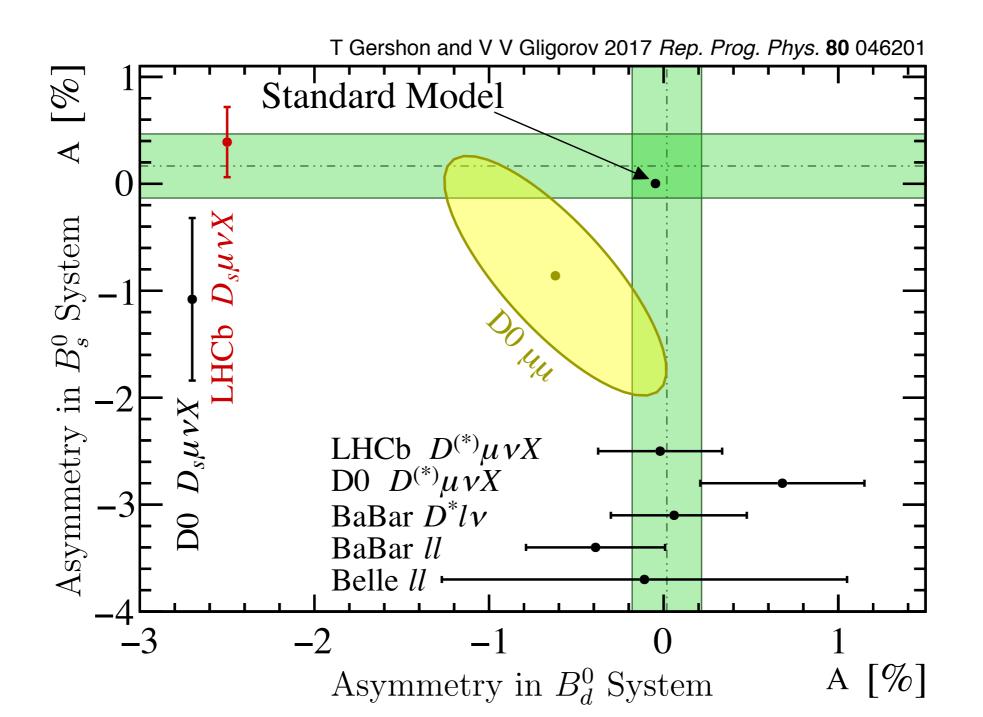
CP Violation

B meson/anti-meson mixing has sizable CP violation



Asymmetry in B Meson Mixing

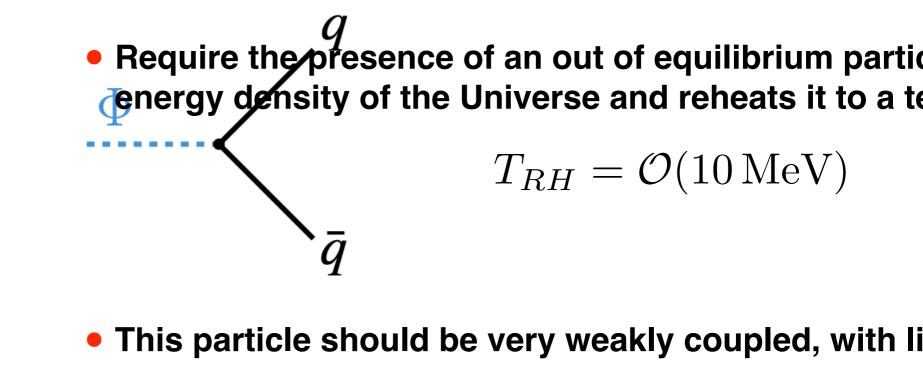
Can accommodate contributions from new physics



Sakharov I. Out of Equilibrium

Late deczy of an "inflaton-like" field

Decays at: $\Gamma_{\Phi} = 4H(T_R)$ to quarks $m_{\Phi} \in [5 \text{ GeV}, 100 \text{ GeV}]$

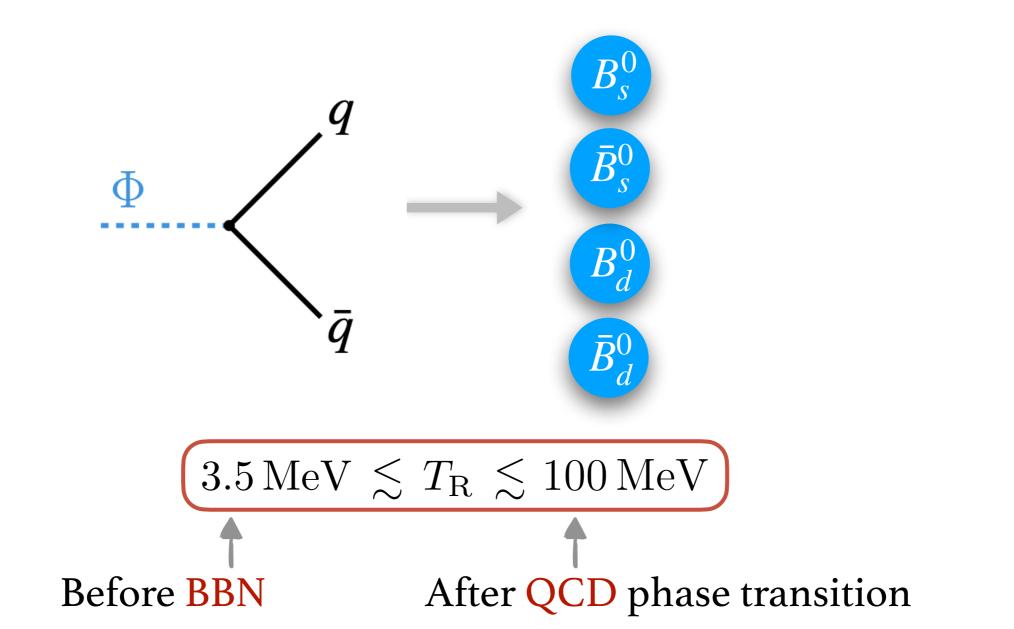


$$\begin{array}{ccc} 3.5\,\mathrm{MeV}\,\lesssim\,T_\mathrm{R}\,\lesssim\,100\,\mathrm{MeV} & \tau_\Phi=\mathcal{O}(10^{-3}\,\mathrm{s}) \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ &$$

Sakharov I. Out of Equilibrium

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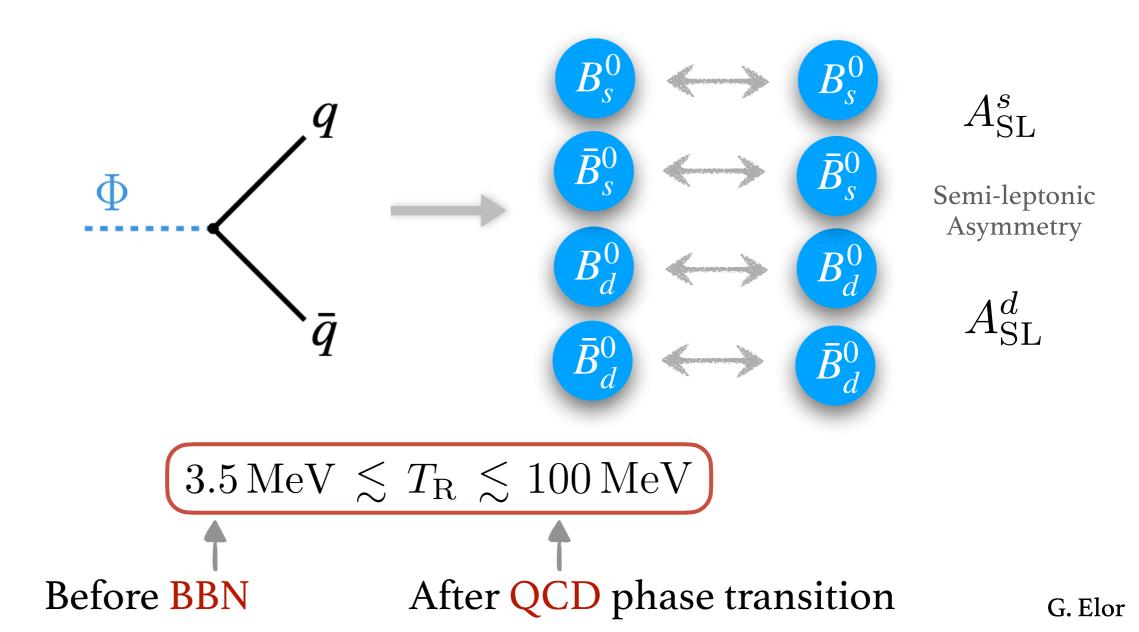
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Sakharov II. CP Violation

Late decay of an "inflaton-like" field

Decays at: $\Gamma_{\Phi} = 4H(T_R)$ to quarks $m_{\Phi} \in [5 \text{ GeV}, 100 \text{ GeV}]$



Sakharov III. B Violation?

Need a way to change baryon number

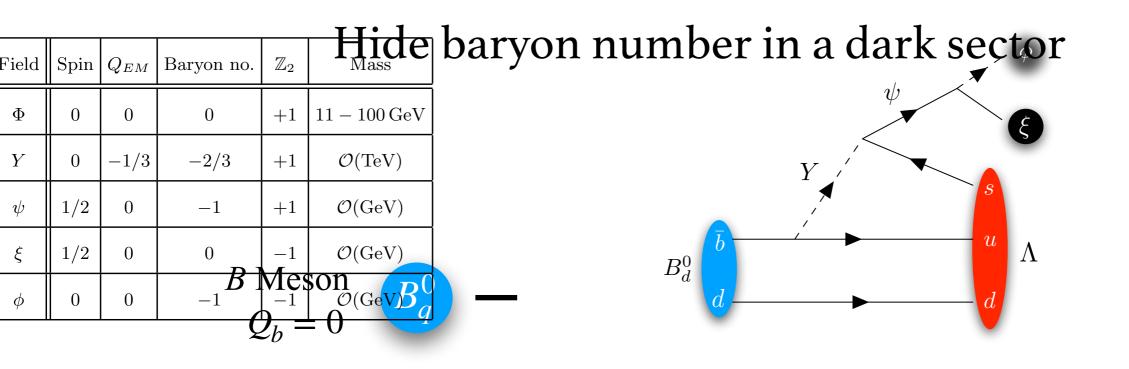


Hide baryon number in a dark sector rather than violate it



Will minimally need four new fields

Dark Sector Baryon



Kinematics: $m_{\psi} < m_B - m_{\text{Baryon}} < 4.3 \,\text{GeV}$ Proton stability: $m_{\psi} > m_p - m_e \simeq 937.8 \,\text{MeV}$

Equal and opposite dark and visible baryon $A^{\alpha}s\bar{y}n\bar{h}\bar{t}\bar{t}\bar{r}\bar{t}\bar{s}^{10}\bar{f}\bar{s}^{36}$ GeVed. $Y_{\mathcal{B}} - Y_{\bar{\mathcal{B}}} = -(Y_{\psi} - Y_{\bar{\psi}})$

Cchreat Appled Scale

	ſ	Field Sp	in Q_{EM} Baryon	1 no. 2	\mathbb{Z}_2 Mass	
Field	Spin	Q_{EM}	Baryon no.	\mathbb{Z}_2	Mass	
Φ	0	0	0	+1	$11-100\mathrm{GeV}$	
Y	0	-1/3	-2/3	+1	$\mathcal{O}({ m TeV})$	
ψ	1/2	0	-1	+1	$\mathcal{O}({ m GeV})$	

Allowed by all the symmetries:

$$\mathcal{L} \supset -y_{ub} Y^* \, \bar{u} \, b^c - y_{\psi s} Y \bar{\psi}$$

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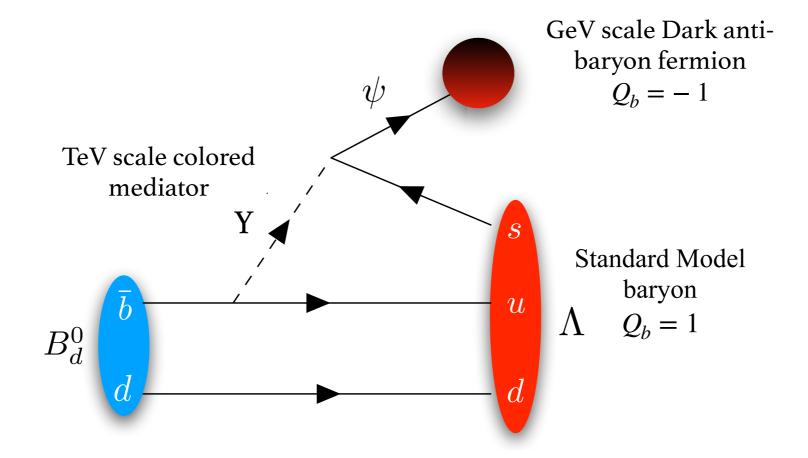
Colored triplet scalar must be TeV state to be consistent with dollider search so we will integrate it out: $m_Y^{2b}y_{1/2}$ so m_Y^{2}

$$\mathcal{H} = \bigcup_{Y}^{\kappa} b \, u \, s \, \psi$$

Note: this interaction does not change baryon nomber $\frac{m_B}{2 \text{ GeV}}$

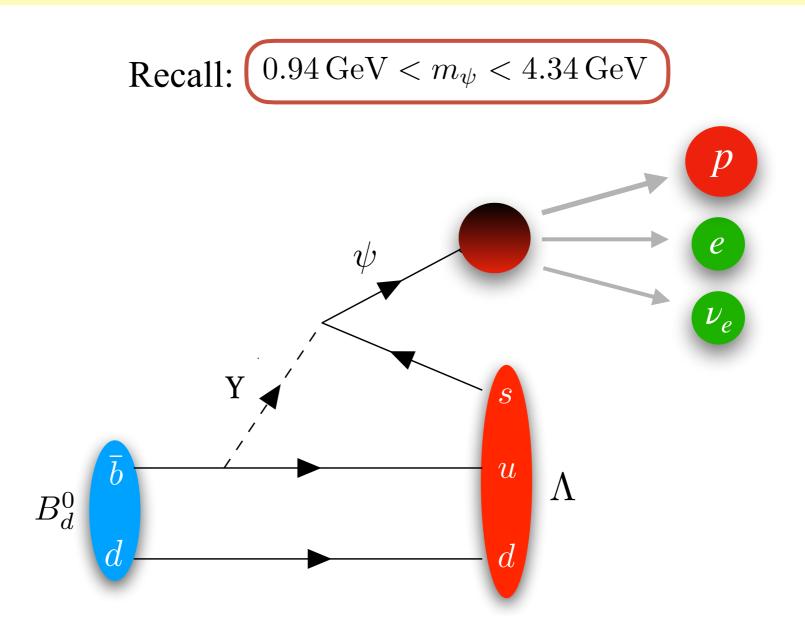
New decay of the B Meson

For instance, $\mathcal{O}_{us} = \psi b \, u \, s$ mediates the decay $\bar{b} \to \psi + u + s$



Decay modes of a *B* meson to baryon and missing energy have historically not been searched for. Relatively unconstrained.

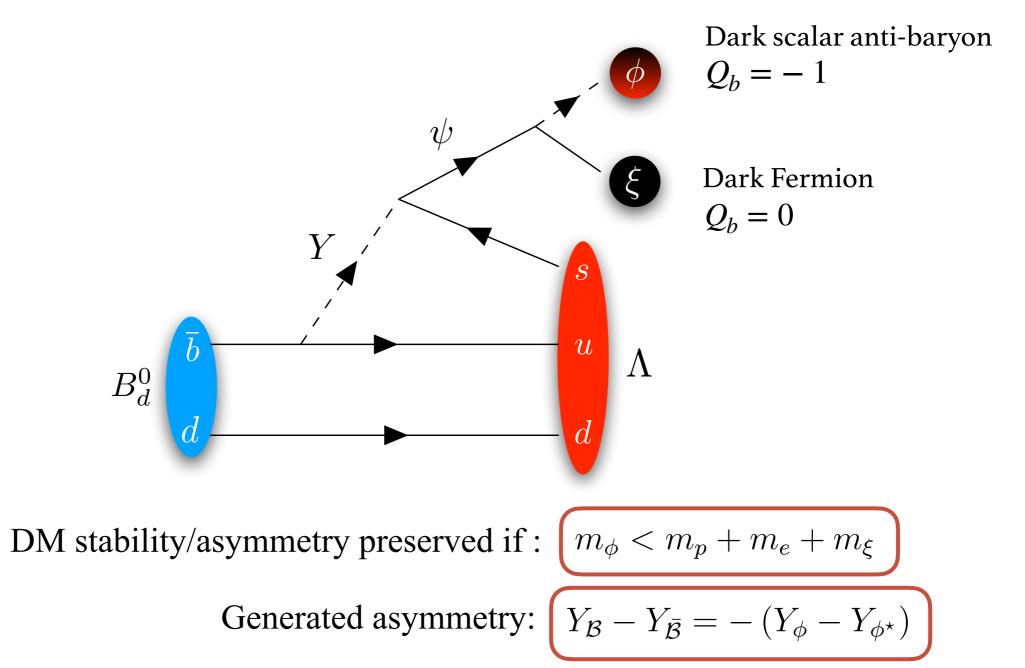
Dark Matter



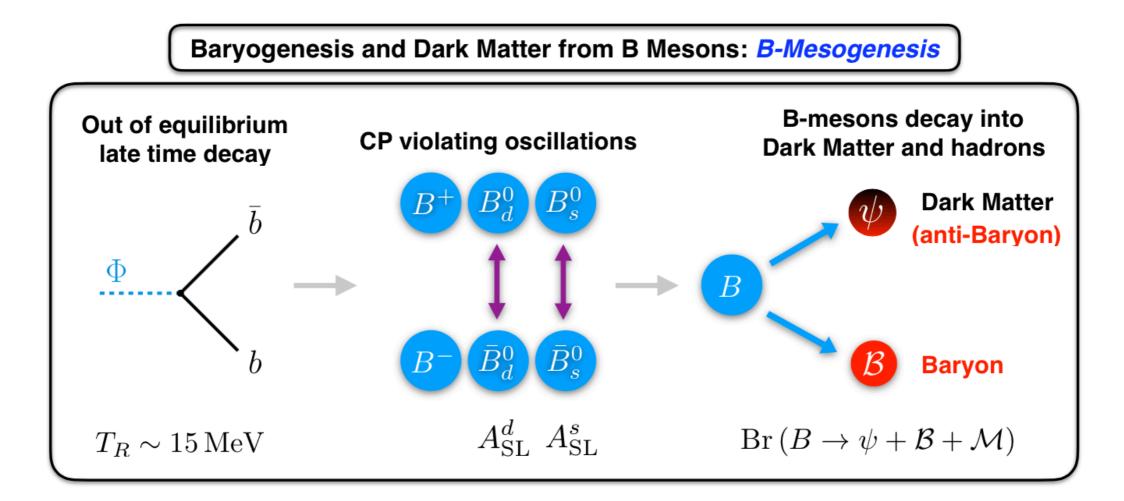
New dark baryon is unstable and will decay to baryonic matter, washing out the asymmetry in the process. It cannot be the dark matter.

Dark Matter

Dark fermion must quickly decay within the dark sector into two new dark sector fields



Summary: The Mechanism



Baryon asymmetry is related to experimental observables:

- Branching fraction for inclusive decay of B mesons into baryons, missing energy, and any number of final state mesons.
- Semi-leptonic asymmetries $Y_B \propto \sum_{q=s,d} A_{SL}^q \times Br\left(B^0 \to \psi \mathcal{BM}\right)$ G. Elor

Summary: Mirainal Mudal

B - Mesogenesis requires the minimal field content:

	Field	Spin	Q_{EM}	Baryon no.	\mathbb{Z}_2	Mass
Inflaton:	Φ	0	0	0	+1	$11 - 100 \mathrm{GeV}$
Colored Triplet Scalar:	Y	0	-1/3	-2/3	+1	$\mathcal{O}({ m TeV})$
Dark Dirac Baryon:	ψ	1/2	0	-1	+1	$\mathcal{O}({ m GeV})$
Dark Fermion:	ξ	1/2	0	0	-1	$\mathcal{O}({ m GeV})$
Dark Scalar Baryon:	ϕ	0	0	-1	-1	$\mathcal{O}({ m GeV})$

Sumary: Mainal Madel

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Aside: this can be embedded in e.g. a supersymmetric theory

 $Y \leftrightarrow \tilde{q}_R$ squark, $\psi \leftrightarrow \begin{bmatrix} \tilde{B} \\ \lambda_s^{\dagger} \end{bmatrix}$ Dirac bino, $\phi, \xi \leftrightarrow$ sterile neutrino multiplet

[GE with G. Alonso-Alvarez, A. E. Nelson, H. Xiao JHEP [arXiv:1907.10612]]

However, we do not need a UV model for numerics or general signals

Sumary: Mainal Madel

B - Mesogenesis requires the minimal field content:

	Field	Spin	Q_{EM}	Baryon no.	\mathbb{Z}_2	Mass	
Inflaton:	Φ	0	0	0	+1	$11-100\mathrm{GeV}$	Next:
Colored Triplet Scalar:	Y	0	-1/3	-2/3	+1	$\mathcal{O}({ m TeV})$	Numerically study the evolution of the
Dark Dirac Baryon:	ψ	1/2	0	-1	+1	$\mathcal{O}({ m GeV})$	Boltzmann equations describing these fields
Dark Fermion:	ξ	1/2	0	0	-1	$\mathcal{O}({ m GeV})$	deserroting these fields
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However, we *do not need a UV model* for numerics or general signals

Late time decay of Inflaton

 $\Gamma_{\Phi} = 4H(T_R)$

• Inflaton: $\frac{dn_{\Phi}}{dt} + 3Hn_{\Phi} = -\Gamma_{\Phi}n_{\Phi}$ • Radiation: $\frac{d\rho_{\text{rad}}}{dt} + 4H\rho_{\text{rad}} = +\Gamma_{\Phi}m_{\Phi}n_{\Phi}$ • Hubble: $H^{2} = \frac{8\pi}{3M_{\text{Pl}}^{2}}\left(\rho_{\text{rad}} + m_{\Phi}n_{\Phi}\right)$

Dark Matter

• Symmetric component of the dark scalar baryon

$$\frac{dn_{\phi+\phi^*}}{dt} + 3Hn_{\phi+\phi^*} = 2\Gamma_{\Phi}^B n_{\Phi} - 2\langle\sigma v\rangle_{\phi} \left(n_{\phi+\phi^*}^2 - n_{\text{eq},\phi+\phi^*}^2\right)$$

• The dark Majorana fermion

$$\frac{dn_{\xi}}{dt} + 3Hn_{\xi} = 2\Gamma_{\Phi}^{B}n_{\Phi} - \langle \sigma v \rangle_{\xi} \left(n_{\xi}^{2} - n_{\text{eq},\xi}^{2} \right)$$

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$$\Gamma_{\Phi}^{B} \equiv \Gamma_{\Phi} \times \operatorname{Br} \left(B \to \psi \mathcal{B} \mathcal{M} \right)$$

Simplification: For the (low) temperature range of interest we can check that the *B* mesons decay more quickly than they annihilate

Dark Matter

• Symmetric component of the dark scalar baryon

$$\frac{dn_{\phi+\phi^*}}{dt} + 3Hn_{\phi+\phi^*} = 2\Gamma_{\Phi}^B n_{\Phi} - 2\langle\sigma v\rangle_{\phi} \left(n_{\phi+\phi^*}^2 - n_{\text{eq},\phi+\phi^*}^2\right)$$

• The dark Majorana fermion

Overproduced particle must be depleted by additional dark interactions.

$$\frac{dn_{\xi}}{dt} + 3Hn_{\xi} = 2\Gamma_{\Phi}^{B}n_{\Phi} - \langle \sigma v \rangle_{\xi} \left(n_{\xi}^{2} - n_{eq,\xi}^{2}\right)$$
$$\Gamma_{\Phi}^{B} \equiv \Gamma_{\Phi} \times \operatorname{Br}\left(B \to \psi \mathcal{B} \mathcal{M}\right)$$

Simplification: For the (low) temperature range of interest we can check that the *B* mesons decay more quickly than they annihilate

Numerics

The Baryon Asymmetry

• Anti-symmetric dark sector baryon makes up the baryon asymmetry

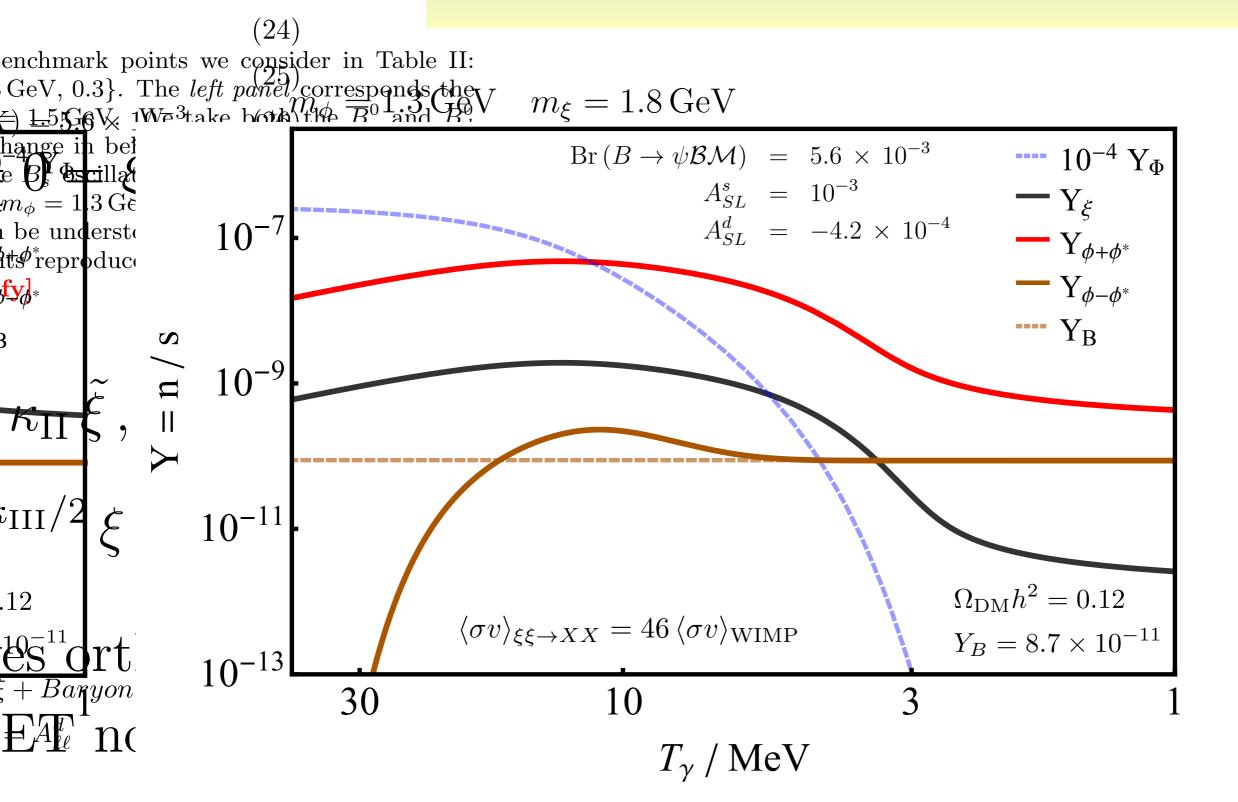
eV

$$\underbrace{\frac{\partial n_{\phi-\phi^*}}{\partial t} + 3Hn_{\phi-\phi^*}}_{b} = 2\Gamma^B_{\Phi}\sum_q \operatorname{Br}\left(\bar{b} \to B^0_q\right)A^q_{\mathrm{SL}}f^q_{\mathrm{deco}}n_{\Phi}$$

 $200 \,\mathrm{MeV}$ Coherent *B* meson oscillations maintained for 20 MeV scales and below

 $e^{\pm} \qquad \Gamma\left(e^{\pm}B_{q}^{0} \to e^{\pm}B_{q}^{0}\right) = 10^{-11} \text{GeV}\left(\frac{T}{20 \text{MeV}}\right)^{5}$ $f_{\text{deco}}^{q} = e^{-\Gamma\left(e^{\pm}B_{q}^{0} \to e^{\pm}B_{q}^{0}\right)/\Delta m_{B_{q}}}$ $B^{0} \qquad T_{B_{s}} \leq 20 \text{ MeV and } T_{B_{d}} \leq 10 \text{ MeV}$

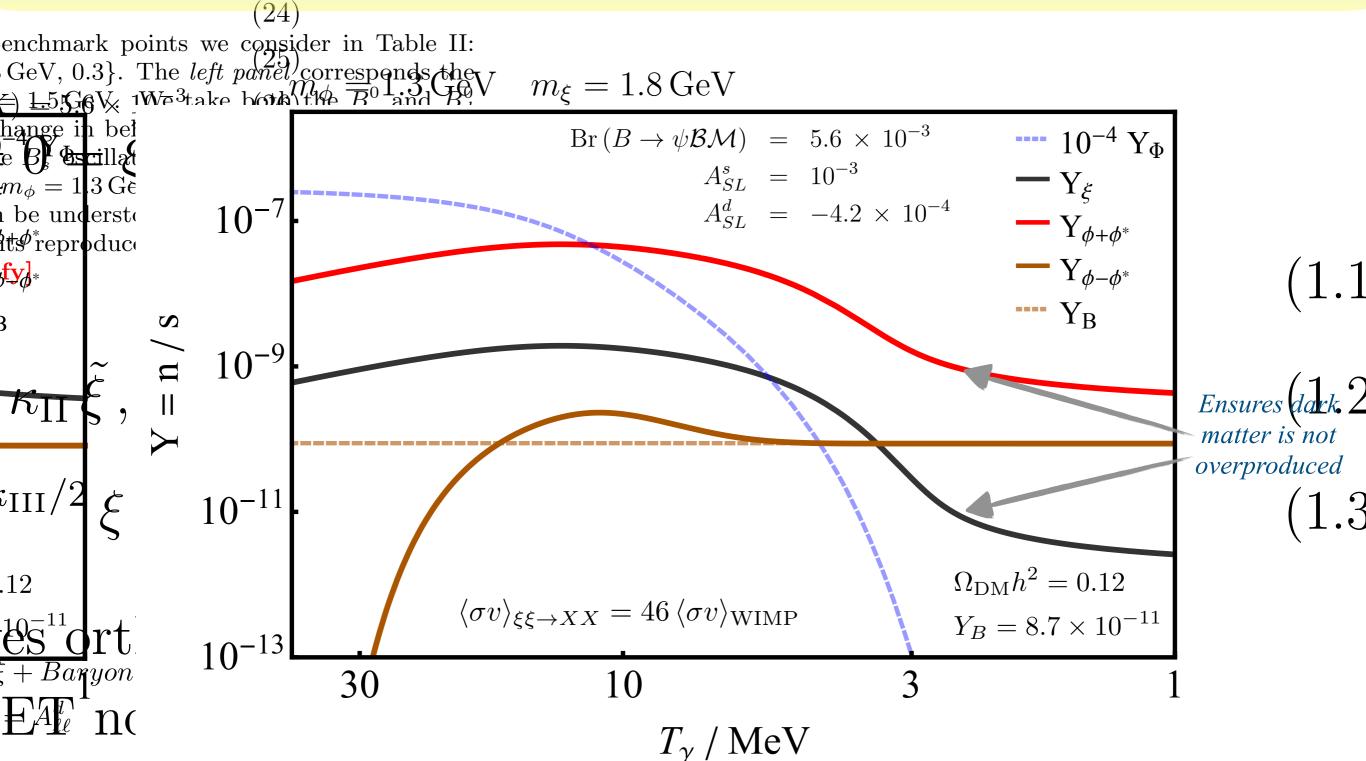
(20) **Example Benchmark Point** (23)



(1.1)(1.2)(1.3)

lS

(20) Example Benchmark Point

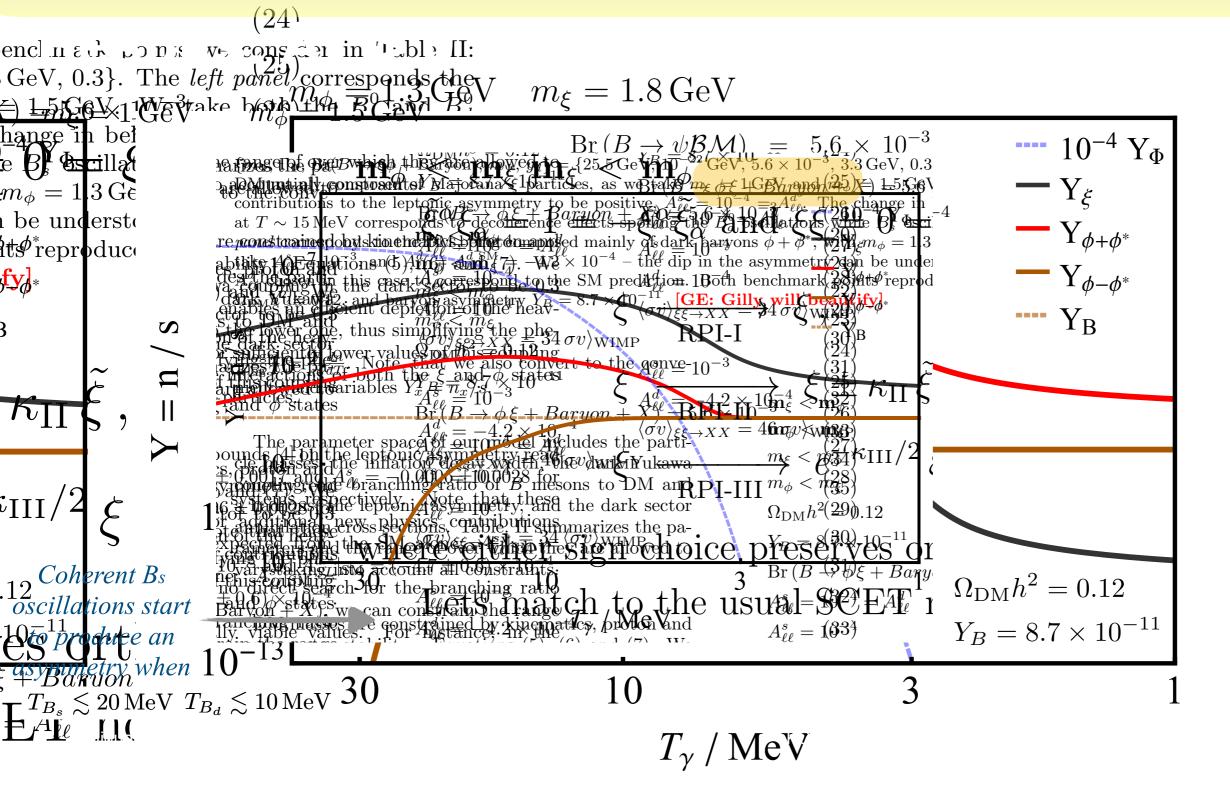


 $\langle \mathbf{a} \mathbf{a} \rangle$

1S

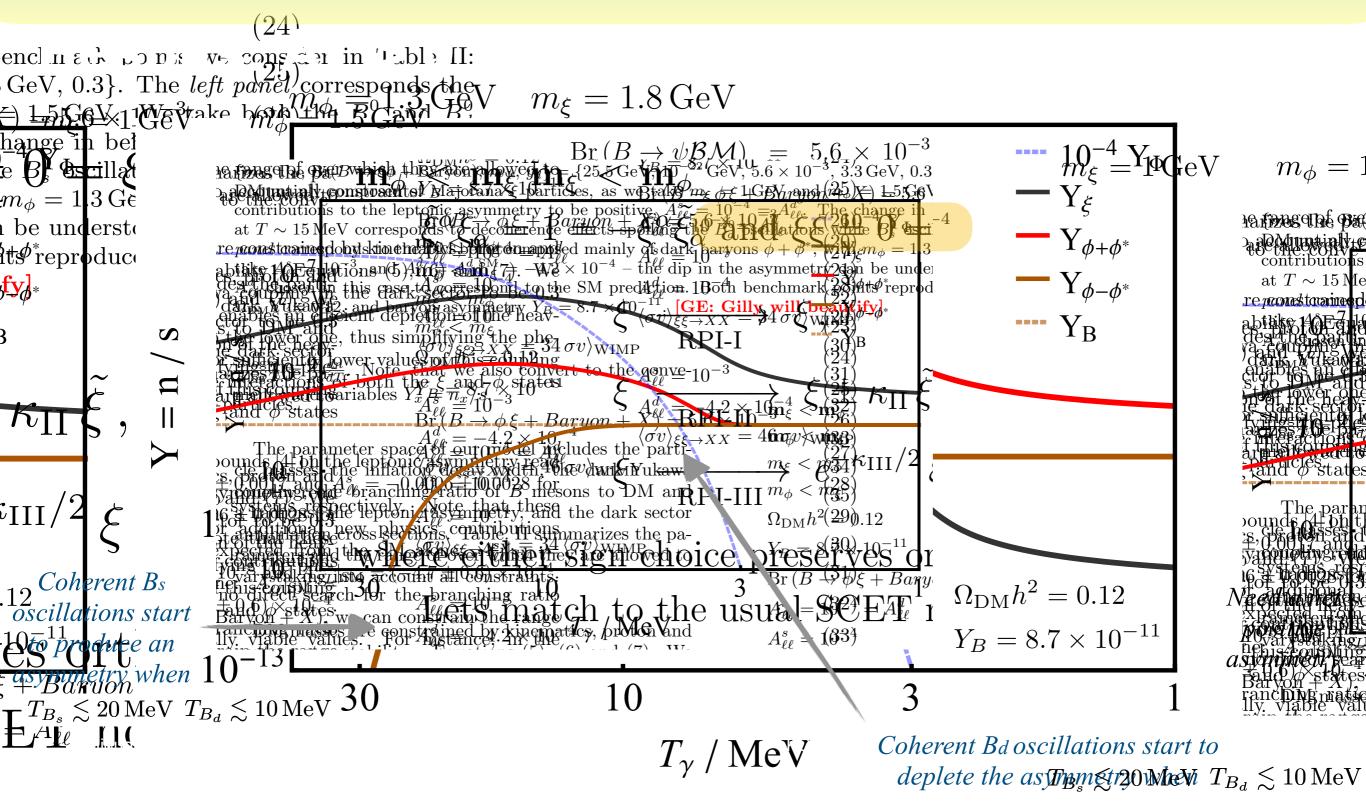
(20) Example Benchmark Point

1S



(1.1)(1.2)(1.3)

(20) Example Benchmark Point



 $\mathbf{1S}$

Numerical Result

$$Y_{\mathcal{B}} \simeq 8.7 \times 10^{-11} \frac{\operatorname{Br} \left(B \to \psi \mathcal{B} \mathcal{M} \right)}{10^{-2}} \sum_{q=s,d} \alpha_q \frac{A_{\mathrm{SL}}^q}{10^{-4}}$$

Successful *B*-Mesogenesis

$$A_{\mathrm{SL}}^{s,d} \times \mathrm{Br}\left(B^0 \to \psi \,\mathcal{B} \,\mathcal{M}\right) > 10^{-6}$$

Experimental Observables

Part II. A Roadmap for Discovering *B*-Mesogenesis



Based on:

GE with Gonzalo Alonso-Alvarez, Miguel Escudero, [arXiv:2101.02706]

Ongoing theoretical work: GE with Gonzalo Alonso-Alvarez, Jorge Martin Camalich, Miguel Escudero, Bartosz Fornal, Benjamin Grinstein (also see white paper coming out in August) Ongoing searches at Belle 1 and II

Possible searches at LHCb - see arXiv:2105.12668 and arXiv:2106.12870

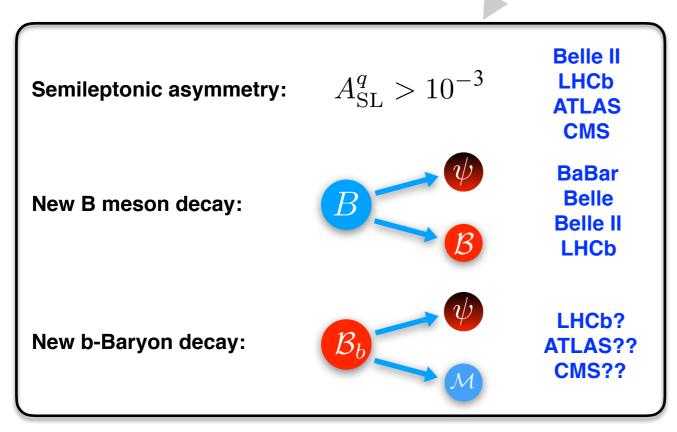
G. Elor

Signals of B-Mesogenesis

For successful baryogenesis: $A_{\rm SL}^{s,d} \times Br(B^0 \to \psi \mathcal{BM}) > 10^{-6}$

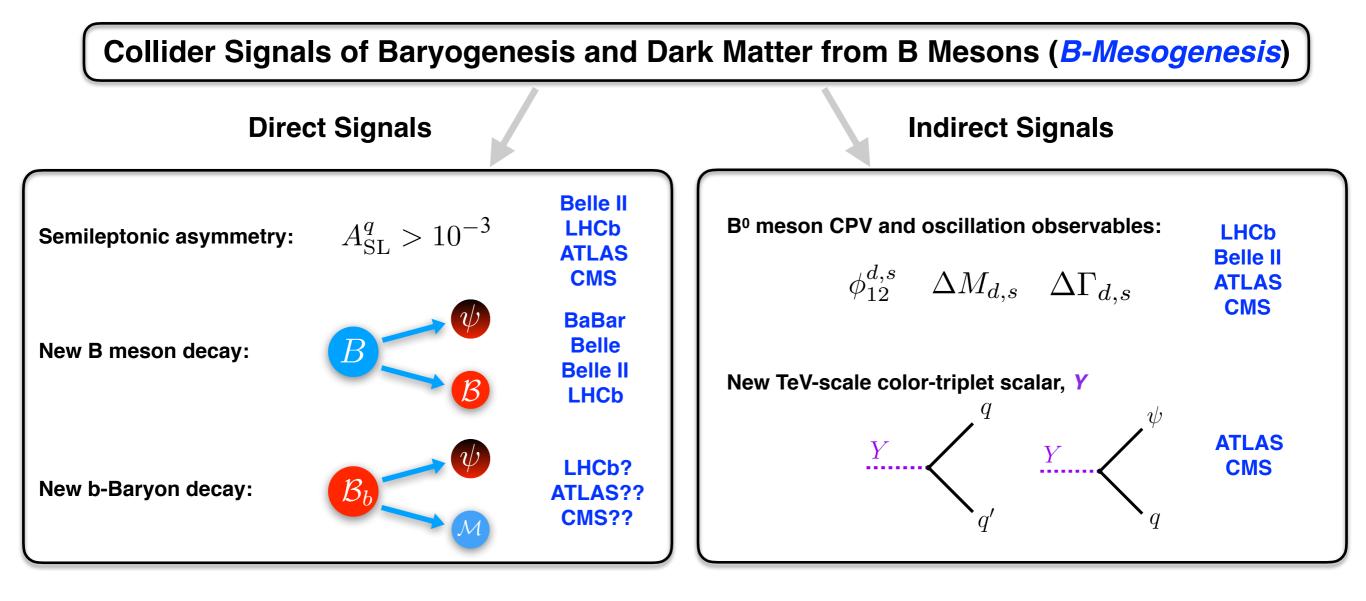
Collider Signals of Baryogenesis and Dark Matter from B Mesons (*B-Mesogenesis*)

Direct Signals



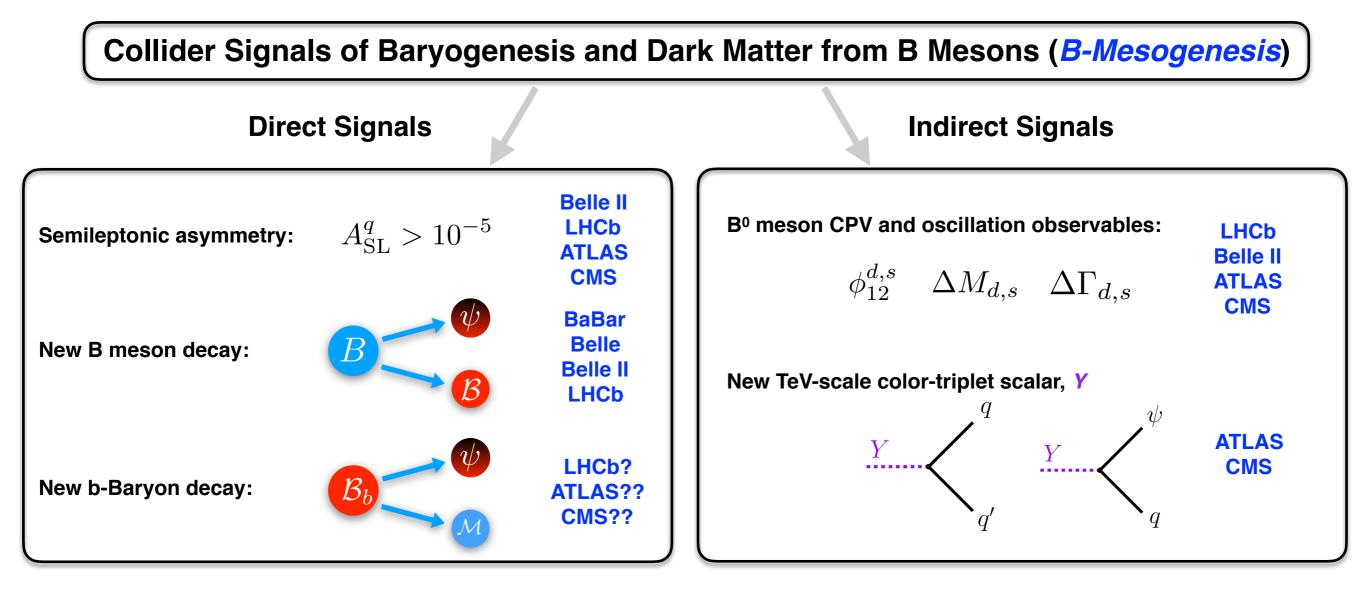
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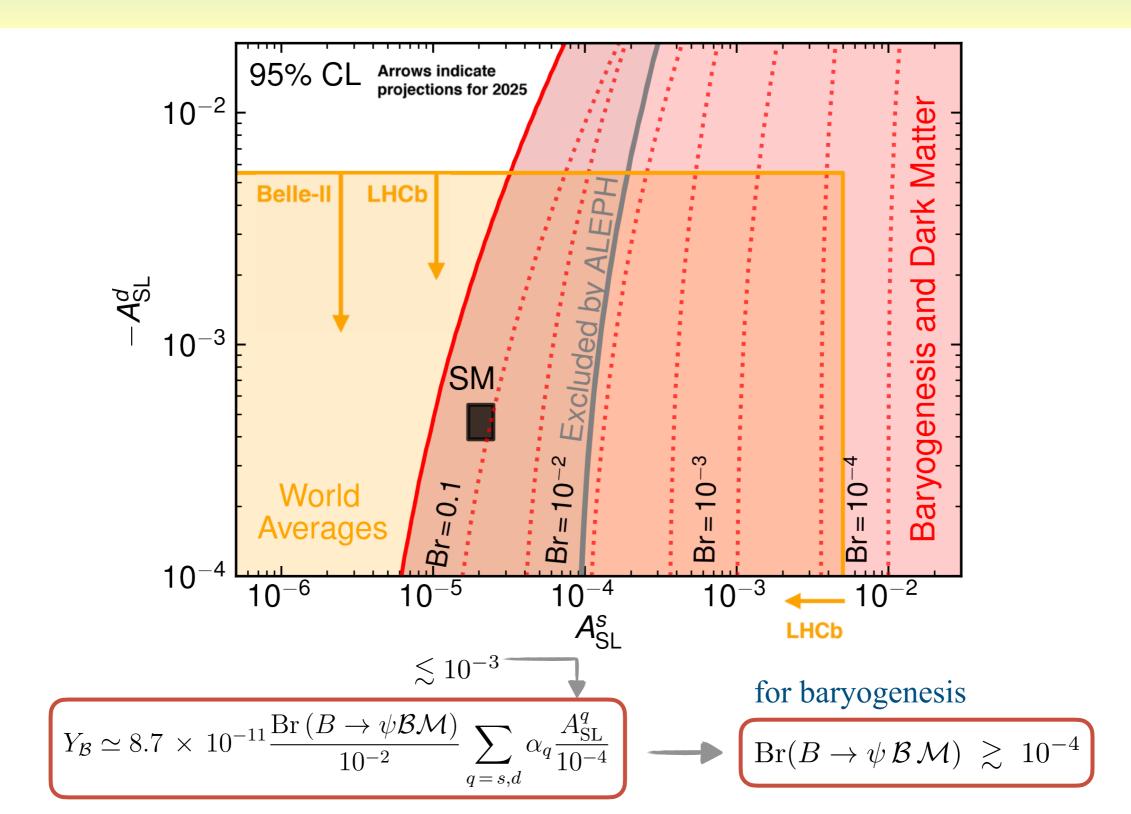
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Independent of UV model. Given a UV model there will be even more signals!

The Semi-Leptonic Asymmetry



G. Elor

Flavorful Variations

No a priori reason to expect a particular flavor structure.

Most general interactions:

$$\mathcal{L}_{-1/3} = -\sum_{i,j} y_{u_i d_j} Y^* \bar{u}_{iR} d_{jR}^c - \sum_k y_{\psi d_k} Y d_{kR}^c \bar{\psi} + \text{h.c.}$$

Possible operators:

$$egin{aligned} \mathcal{O}_{ud} &= \psi \, b \, u \, d \ \mathcal{O}_{us} &= \psi \, b \, u \, s \ \mathcal{O}_{cd} &= \psi \, b \, c \, d \ \mathcal{O}_{cs} &= \psi \, b \, c \, s \end{aligned}$$

B-Mesogenesis requires:

$$\operatorname{Br}(B \to \psi \,\mathcal{B} \,\mathcal{M}) \gtrsim 10^{-4}$$

Can be searched for at Belle, BaBar and LHCb

Operator/Decay	Initial State	Final state	
$\mathcal{O} = \psi b u d$ $\bar{b} \to \psi u d$	B_d B_s B^+ Λ_b	$\begin{split} \psi + n (udd) \\ \psi + \Lambda (uds) \\ \psi + p (duu) \\ \bar{\psi} + \pi^0 \end{split}$	
$\mathcal{O} = \psi b u s$ $\bar{b} \to \psi u s$	B_d B_s B^+ Λ_b	$\psi + \Lambda (usd)$ $\psi + \Xi^{0} (uss)$ $\psi + \Sigma^{+} (uus)$ $\overline{\psi} + K^{0}$	
$\mathcal{O} = \psi b c d$ $\overline{b} \to \psi c d$	B_d B_s B^+ Λ_b	$\psi + \Lambda_c + \pi^- (cdd)$ $\psi + \Xi_c^0 (cds)$ $\psi + \Lambda_c (dcu)$ $\overline{\psi} + \overline{D}^0$	
$\mathcal{O} = \psi b c s$ $\bar{b} \to \psi c s$	B_d B_s B^+ Λ_b	$\psi + \Xi_c^0 (csd)$ $\psi + \Omega_c (css)$ $\psi + \Xi_c^+ (csu)$ $\bar{\psi} + D^- + K^+$	

Can be searched for at Belle, BaBar and LHCb

Operator/Decay	Initial State	Final state	
$\mathcal{O} = \psi b u d$ $\bar{b} \to \psi u d$	$ \begin{array}{c} B_d \\ B_s \\ B^+ \\ \Lambda_b \end{array} $	$egin{aligned} \psi + n (udd) \ \psi + \Lambda (uds) \ \psi + p (duu) \ ar{\psi} + \pi^0 \end{aligned}$	 Directly related to baryon asymmetry Most stringent constraints actually comes from a 20 year
$\begin{aligned} \mathcal{O} &= \psi b u s \\ \bar{b} &\to \psi u s \end{aligned}$	$ \begin{array}{c} B_d \\ B_s \\ B^+ \\ \overline{\Lambda_b} \end{array} $	$\begin{split} \psi + \Lambda \left(usd \right) \\ \psi + \Xi^0 \left(uss \right) \\ \psi + \Sigma^+ \left(uus \right) \\ \bar{\psi} + K^0 \end{split}$	old search at LEP [hep-ex/0010022]
$\mathcal{O} = \psi b c d$ $\bar{b} \to \psi c d$	$ \begin{array}{c} B_d\\ B_s\\ B^+\\ \Lambda_b \end{array} $	$\psi + \Lambda_c + \pi^- (cdd)$ $\psi + \Xi_c^0 (cds)$ $\psi + \Lambda_c (dcu)$ $\bar{\psi} + \overline{D}^0$	
$\mathcal{O} = \psi b c s$ $\bar{b} \to \psi c s$	$ \begin{array}{c} B_d\\ B_s\\ B^+\\ \Lambda_b \end{array} $	$\psi + \Xi_c^0 (csd)$ $\psi + \Omega_c (css)$ $\psi + \Xi_c^+ (csu)$ $\bar{\psi} + D^- + K^+$	G. E

Can be searched for at Belle, BaBar and LHCb

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$\mathcal{O} = \psi b u d$ $\bar{b} \to \psi u d$	B_d B_s B^+ Λ_b	$\begin{split} \psi + n (udd) \\ \psi + \Lambda (uds) \\ \psi + p (duu) \\ \bar{\psi} + \pi^0 \end{split}$	Directly related to baryon asymmetry Most stringent constraints actually comes from a 20 year
$\mathcal{O} = \psi b u s$ $\bar{b} \to \psi u s$	$ \begin{array}{c} B_d \\ B_s \\ B^+ \\ \hline \Lambda_b \end{array} $	$\begin{split} \psi + \Lambda \left(usd \right) \\ \psi + \Xi^0 \left(uss \right) \\ \psi + \Sigma^+ \left(uus \right) \\ \bar{\psi} + K^0 \end{split}$	 old search at LEP [hep-ex/0010022] Indirectly constrains B-Mesogenesis. Charged track is an advantage for searches
$\mathcal{O} = \psi b c d$ $\bar{b} \to \psi c d$	$ \begin{array}{c} B_d \\ B_s \\ B^+ \\ \Lambda_b \end{array} $	$\psi + \Lambda_c + \pi^- (cdd)$ $\psi + \Xi_c^0 (cds)$ $\psi + \Lambda_c (dcu)$ $\overline{\psi} + \overline{D}^0$	$B^{+} \underbrace{u}_{V} \underbrace{u}$
$\mathcal{O} = \psi b c s$ $\bar{b} \to \psi c s$	B_d B_s B^+ Λ_b	$\psi + \Xi_c^0 (csd)$ $\psi + \Omega_c (css)$ $\psi + \Xi_c^+ (csu)$ $\bar{\psi} + D^- + K^+$	$B^{+} \underbrace{\begin{matrix} u \\ b \\ y \\ y \\ \psi \end{matrix}} \Lambda_{c} \qquad B^{+} \underbrace{\begin{matrix} u \\ b \\ y \\ y \\ \psi \end{matrix}} \Xi_{c}^{+}$ G. Elor

Can be searched for at Belle, BaBar and LHCb

Operator/Decay	Initial State	Final state	
$\mathcal{O} = \psi b u d$ $\overline{b} \to \psi u d$	$ \begin{array}{c} B_d \\ B_s \\ B^+ \\ \Lambda_b \end{array} $	$\begin{split} \psi + n (udd) \\ \psi + \Lambda (uds) \\ \psi + p (duu) \\ \bar{\psi} + \pi^0 \end{split}$	Directly related to baryon asymmetry Most stringent constraints actually comes from a 20 year
$\mathcal{O} = \psi b u s$ $\bar{b} \to \psi u s$	$ \begin{array}{c} B_d \\ B_s \\ B^+ \\ \hline \Lambda_b \end{array} $	$\begin{split} \psi + \Lambda \left(usd \right) \\ \psi + \Xi^0 \left(uss \right) \\ \psi + \Sigma^+ \left(uus \right) \\ \bar{\psi} + K^0 \end{split}$	 old search at LEP [hep-ex/0010022] Indirectly constrains B-Mesogenesis. Charged track is an advantage for searches
$\mathcal{O} = \psi b c d$ $\bar{b} \to \psi c d$	$ \begin{array}{c} B_d\\ B_s\\ B^+\\ \Lambda_b \end{array} $	$\psi + \Lambda_c + \pi^- (cdd)$ $\psi + \Xi_c^0 (cds)$ $\psi + \Lambda_c (dcu)$ $\overline{\psi} + \overline{D}^0$	b-flavored baryon decays can
$\mathcal{O} = \psi b c s$ $\bar{b} \to \psi c s$	$ \begin{array}{c} B_d \\ B_s \\ B^+ \\ \Lambda_b \end{array} $	$\psi + \Xi_c^0 (csd)$ $\psi + \Omega_c (css)$ $\psi + \Xi_c^+ (csu)$ $\bar{\psi} + D^- + K^+$	yield indirect constraints. expect: $Br(\mathcal{B}_b \to \bar{\psi} \mathcal{M}) > 10^{-4}$ G. Elo

Ongoing Belle-I and II Search

Flavorful variations:

Operator/Decay	Initial State	Final state	
$\mathcal{O} = \psi b u d$	B_d B_s	$\psi + n (udd) \ \psi + \Lambda (uds)$	
$\overline{b} \to \psi u d$	$B^+ \ \Lambda_b$	$\psi + p \left(duu ight) \ ar{\psi} + \pi^0$	
	B_d	$\psi + \Lambda \left(usd ight)$	
$\mathcal{O} = \psi b u s$	B_s	$\psi + \Xi^0 (uss)$	
$b \rightarrow \psi u s$	B^+	$\psi + \Sigma^+ (uus)$	
	Λ_b	$\bar{\psi} + K^0$	
	B_d	$\psi + \Lambda_c + \pi^- (cdd)$	
$\mathcal{O} = \psi b c d$ $\bar{b} \to \psi c d$	B_s	$\psi + \Xi_c^0 (cds)$	
	B^+	$\psi + \Lambda_c (dcu)$	
	Λ_b	$ar{\psi} + \overline{D}^0$	
	B_d	$\psi + \Xi_c^0 \ (csd)$	
$\mathcal{O} = \psi b c s$	B_s	$\psi + \Omega_c \left(css \right)$	
$b \rightarrow \psi c s$	B^+	$\psi + \Xi_c^+ (csu)$	
	Λ_b	$\bar{\psi} + D^- + K^+$	

Results to be made public at the next
 collaboration meeting. They plan to look for other decay modes as well

Possibilities at LHCb

[See our white paper on "Stealth Physics at LHCb" 2105.12668]

- No handle on initial energy of decaying *B* meson so measuring missing energy is non-trivial.
- But, LHCb has advantages: larger number of *B* mesons produced than at Belle, excellent vertex resolution, and good particle reconstruction efficiencies.
- Some possibilities for searches do exist. e.g. new paper just last week!

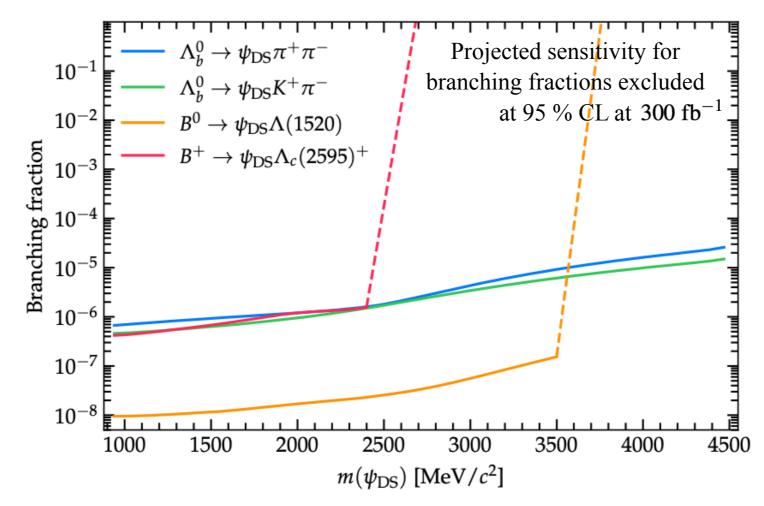
Prospects on searches for baryonic Dark Matter produced in *b***-hadron decays at LHCb** [2106.12870]

Alexandre Brea Rodríguez^{a,1}, Veronika Chobanova^{b,1}, Xabier Cid Vidal^{c,1}, Saúl López Soliño^{d,1}, Diego Martínez Santos^{e,1}, Titus Mombächer^{f,1}, Claire Prouvé^{g,1}, Emilio Xosé Rodríguez Fernández^{h,1}, Carlos Vázquez Sierra^{i,2}

¹Instituto Galego de Física de Altas Enerxías (IGFAE), Universidade de Santiago de Compostela, 15782 Santiago de Compostela, Spain ²European Organization for Nuclear Research (CERN), Geneva, Switzerland

Proposed Search at LHCb [2106.12870]

- Search for decays of *B* mesons and *b*-Flavored baryons into an excited baryon in the final state $B \to \psi B^*$
- The excited baryon promptly decay at the same decay point as original decay, allowing one to trigger on this decay.



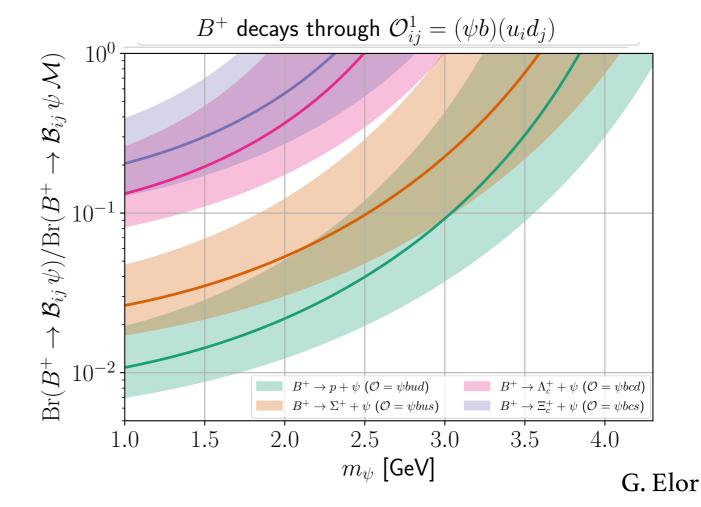
Caution: Inclusive vs. Exclusive Rates

• All decays (and their searches) discussed thus far have been *exclusive*. But, the observable controlling the baryon asymmetry is an *inclusive* rate.

$$Br(B \to \psi \mathcal{BM}) \gtrsim 10^{-4}$$

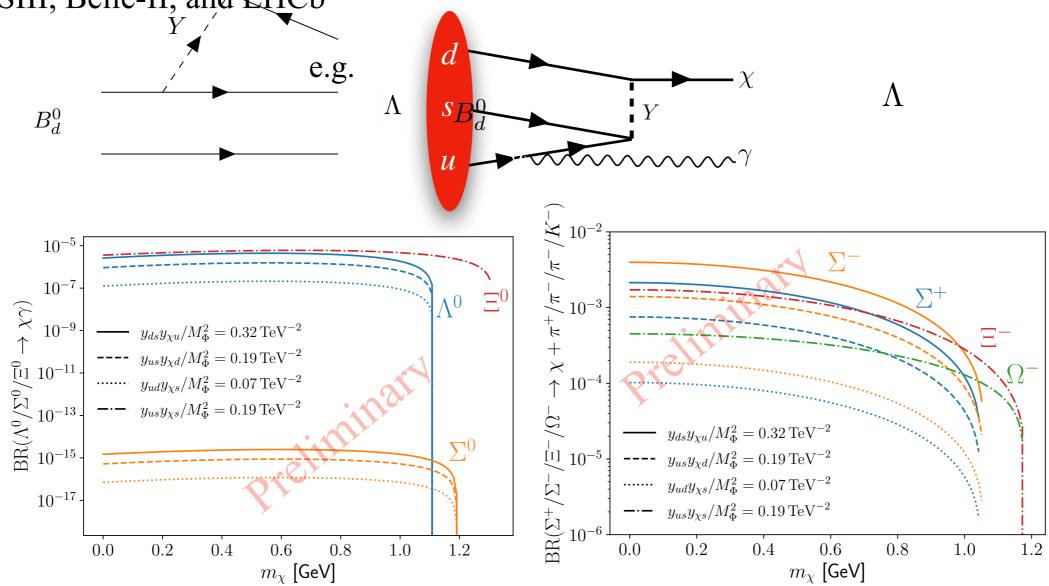
- Need a dedicated calculation using QCD sum rules or lattice techniques etc. to calculate form factors. Beyond my current expertise....
- Phase space method [Bigi, Phys.Lett.B 106, 510 (1981)]

$$\frac{\operatorname{Br}(B \to \psi \mathcal{B})}{\operatorname{Br}(B \to \psi \mathcal{B} \mathcal{M})} \gtrsim (1 - 10) \%.$$



New Hyperon Decays

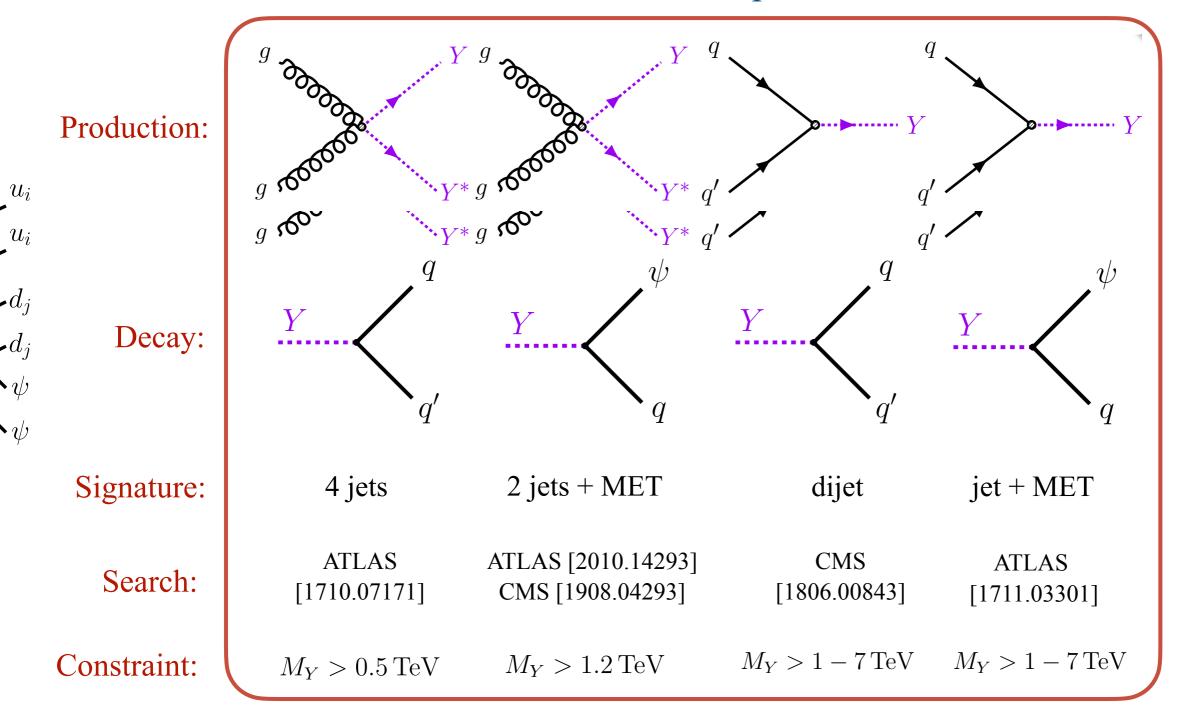
Light hadrons: we can compute form factors by matching onto chiral EFT. Hyperon decays are another indirect probe of B-Mesogenesis and are of interest at BESIII, Belle-II, and LHCb



[GE with Gonzalo Alonso-Alvarez, Jorge Martin Camalich, Miguel Escudero, Bartosz Fornal, Benjamin Grinstein]

Colored Triplet Scalar

Constraints from LHC squark searches

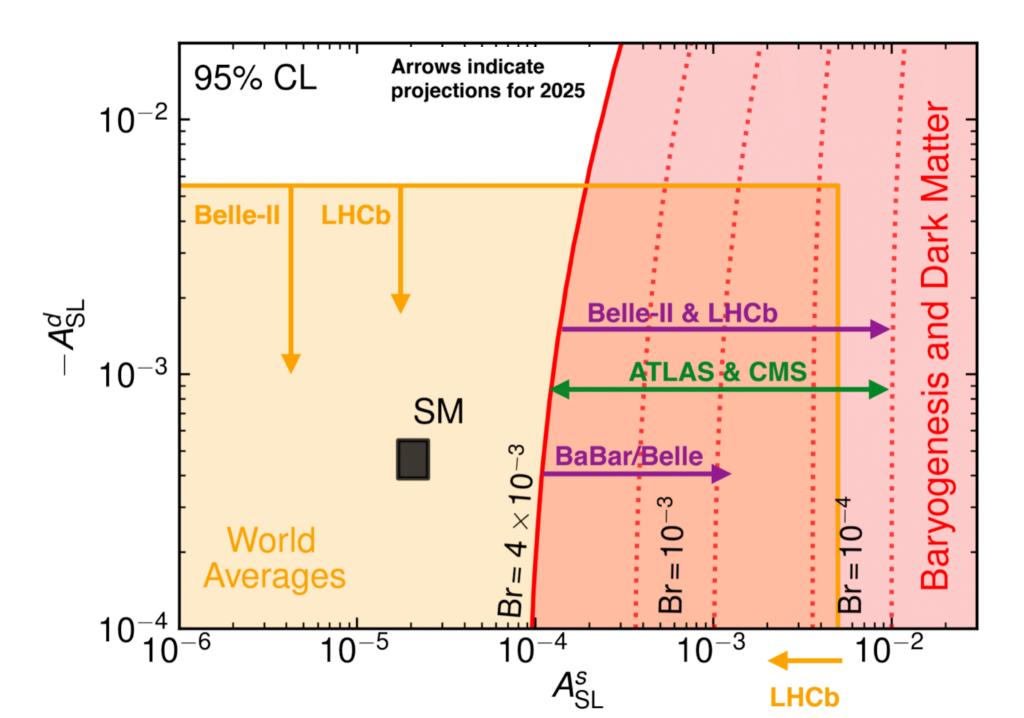


G. Elor

$$\frac{\overline{b} \rightarrow \psi u d}{Coolered Triplet Scalar}$$
Constraints from LHC squark searches
$$B-Mcsogenesis requires: Br(B \rightarrow \psi B M) \simeq 10^{-3} \left(\frac{\Delta m}{3 \text{ GeV}}\right)^4 \left(\frac{1.5 \text{ TeV}}{M_Y} \frac{\sqrt{y}u dy dy d}{0.53}\right)^4 \gtrsim 10^{-4}$$
Br > 10⁻⁴ $\Delta m = m_B - m_{\psi} - m_B - m_M$
Since collider bounds depend on the ratio $\frac{\sqrt{y}u_i d_j y d_{yd_k}}{M_Y}$ they will in turn constrain the branching fraction.
$$c.g. \int_{0}^{10^{-1}} \frac{\sqrt{y}u_i d_j y d_{yd_k}}{10^{-1}} \int_{0}^{10^{-2}} \frac{\sqrt{y}u_i d_j y d_{yd_k}}{10^{-1}} \int_{0}^{10^{-2}} \frac{\sqrt{y}u_i d_j y d_{yd_k}}{1.25 - 1.50 - 1.75} \int_{0}^{2.00} \frac{2.25 - 2.50}{2.50}$$

Discovering B-Mesogenesis Outlook

Could be fully tested in but a few years.



Discovering B-Mesogenesis Outlook

Predictions and Signals of B-Mesogenesis

• New *B* meson decay modes with

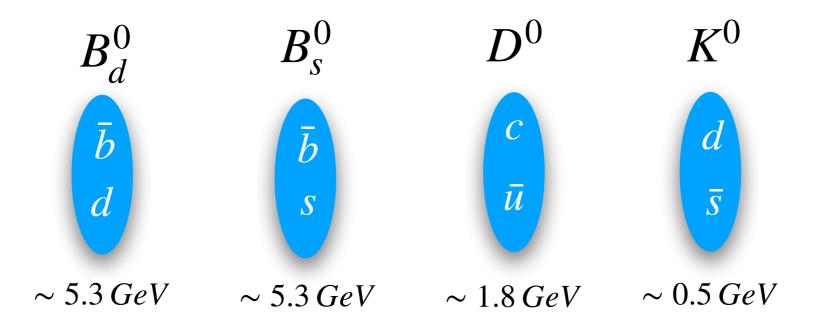
$$\operatorname{Br}(B \to \psi \,\mathcal{B} \,\mathcal{M}) \gtrsim 10^{-4}$$

• Positive semileptonic asymmetry

$$A_{\rm SL}^q > 10^{-4}$$

- The existence TeV scale colored triplet scalar
- Implications for flavor structure.
- Many more signals possible given a UV model see e.g. [1907.10612]

Why B Mesons?



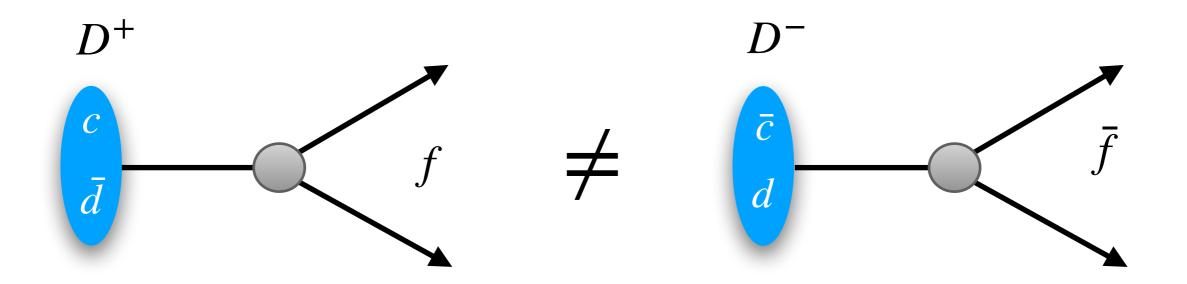
- Kinematics: Dark baryons must be GeV scale. Only *B* mesons are heavy enough to decay into GeV scale. Charge dark particle under lepton number instead, then it can be light.
- Neutral *D* Mesons don't have a lot of CP violation in their oscillations, but charged *D* Mesons have a lot of CP violation in their decays.

Part III. D-Mesogenesis

Based on:

GE with Robert McGehee, Phys. Rev. D [arXiv:2011.06115]

CPV in Charged D Decays



Observable:
$$A_{CP}^f = \frac{\Gamma(D^+ \to f) - \Gamma(D^- \to \bar{f})}{\Gamma(D^+ \to f) + \Gamma(D^- \to \bar{f})}$$

G. Elor

CPV in Charged D Decays

Example: Standard Model decays to an odd number of charged pions

D^+ decay mode	$A_{CP}^{f}/10^{-2}$	D^+ decay mode	$A_{CP}^{f}/10^{-2}$
$K_S^0 \pi^+$	-0.41 ± 0.09	$\pi^+\eta$	1.0 ± 1.5
$K^{-}\pi^{+}\pi^{+}$	-0.18 ± 0.16	$\pi^+\eta'(958)$	-0.6 ± 0.7
$K^{-}\pi^{+}\pi^{+}\pi^{0}$	$-0.3 \pm 0.6 \pm 0.4$	$K^+K^-\pi^+$	0.37 ± 0.29
$K_S^0 \pi^+ \pi^0$	$-0.1 \pm 0.7 \pm 0.2$	$\phi\pi^+$	0.01 ± 0.09
$K_S^0 \pi^+ \pi^+ \pi^-$	$0.0 \pm 1.2 \pm 0.3$	$a_0(1450)^0\pi^+$	$-19 \pm 12^{+8}_{-11}$
$\pi^+\pi^0$	2.4 ± 1.2	$\phi(1680)\pi^{+}$	$-9 \pm 22 \pm 14$
$\pi^+\eta$	1.0 ± 1.5	$\pi^+\pi^+\pi^-$	-1.7 ± 4.2

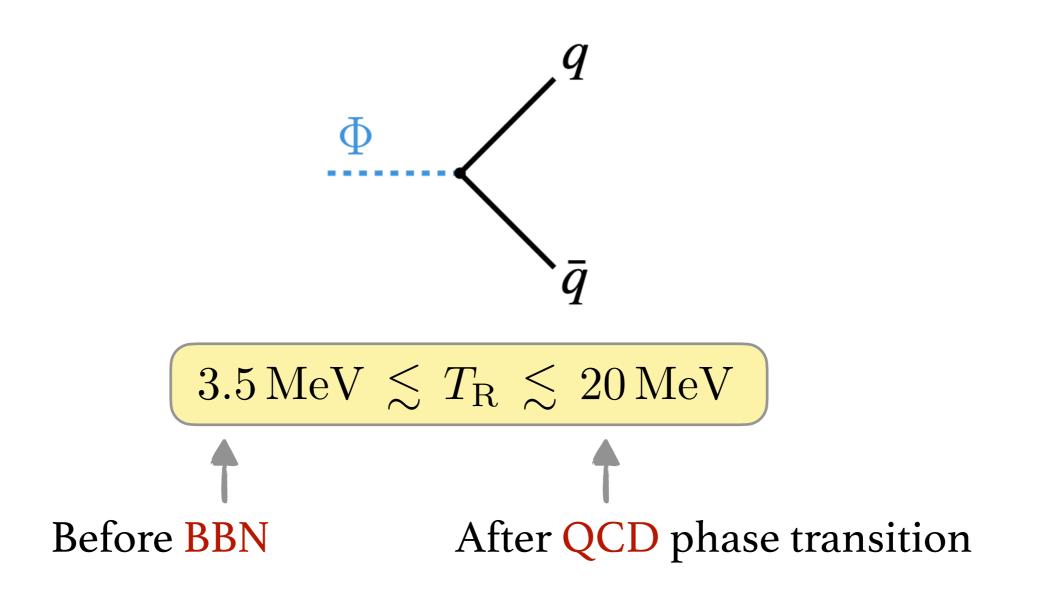
Not a small number if we want to explain

$$Y_B^{\text{obs}} = (8.718 \pm 0.004) \times 10^{-11}$$

Sakharov I. Out of Equilibrium

Late decay of an "inflaton-like" field

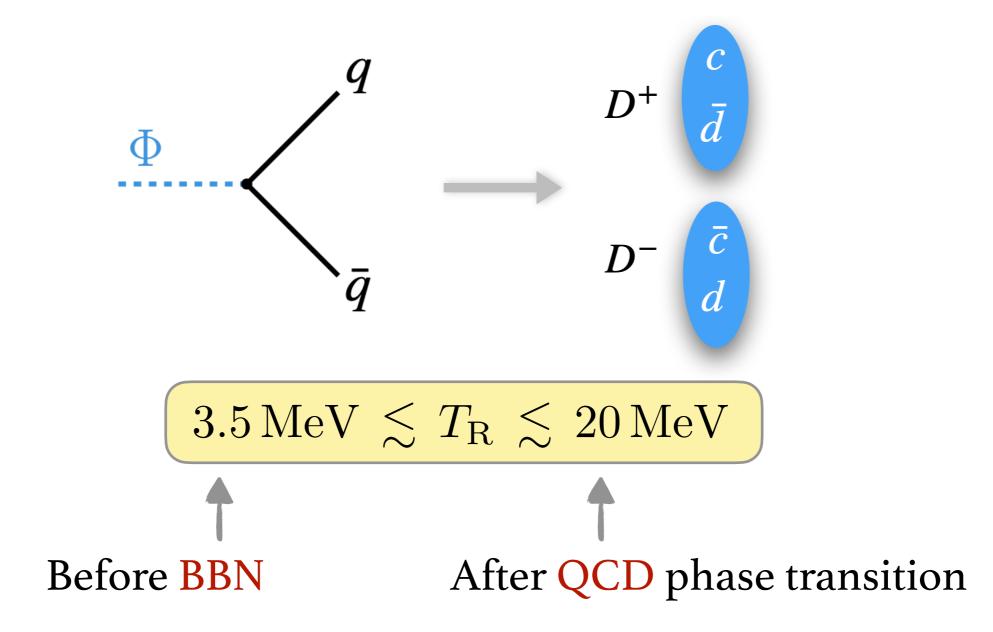
Decays at: $\Gamma_{\Phi} = 4H(T_R)$ to quarks $m_{\Phi} \in [5 \text{ GeV}, 100 \text{ GeV}]$



Sakharov I. Out of Equilibrium

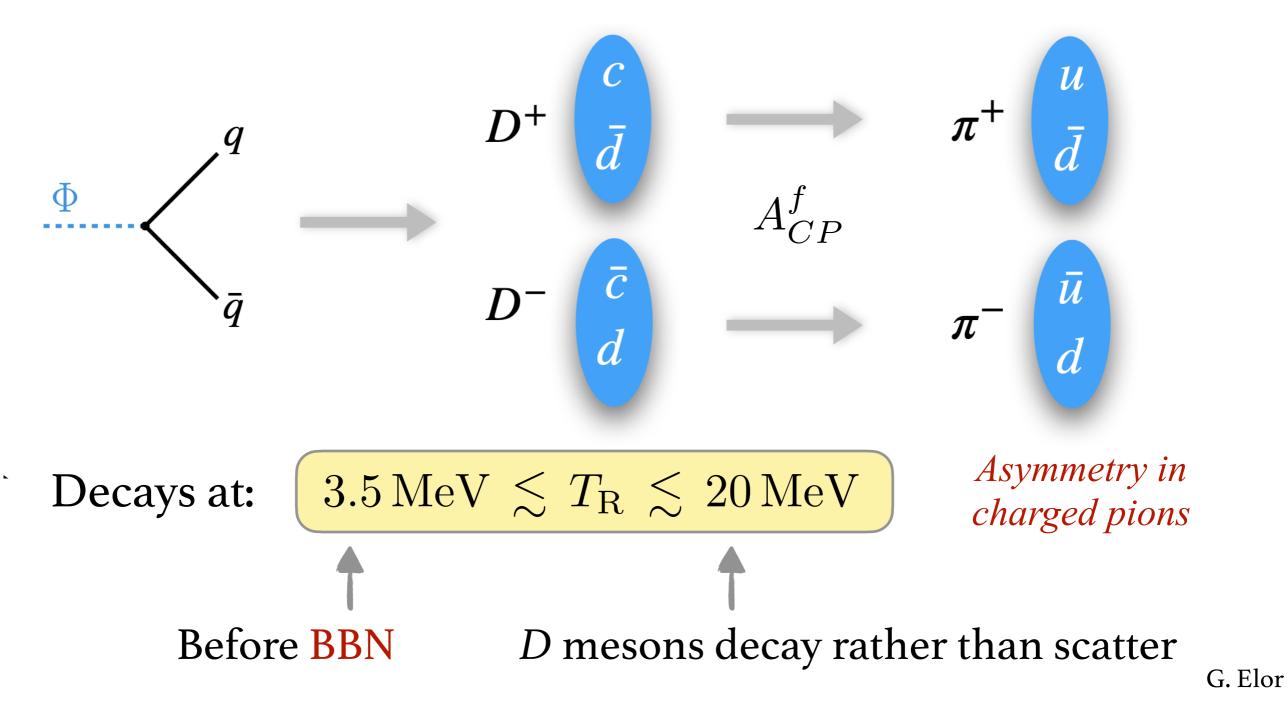
Late decay of an "inflaton-like" field

Decays at: $\Gamma_{\Phi} = 4H(T_R)$ to quarks $m_{\Phi} \in [5 \text{ GeV}, 100 \text{ GeV}]$



Sakharov II. CP Violation

D mesons quickly undergo Standard Model decays to pions



Sakharov III. B Violation?

Need a way to change baryon number



Hide baryon *and lepton* number in a dark sector without violating either.



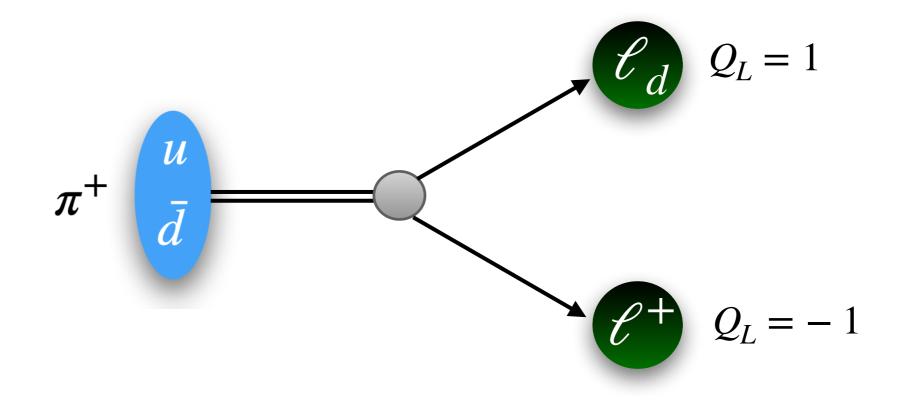
First generate a lepton asymmetry

Dark Sector Lepton

Portal Operator:

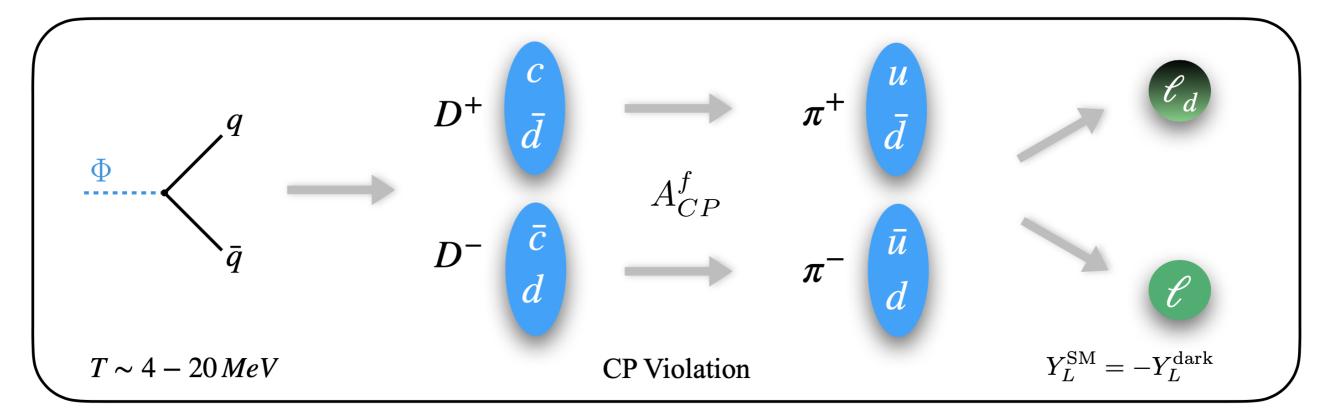
$$\mathcal{O} = \frac{1}{\Lambda^2} \Big[\bar{d} \Gamma^{\mu} u \Big] \Big[\bar{\ell}_d \Gamma_{\mu} \ell \Big] + \text{h.c.}$$

Pion Decays: $\pi^+ \to \ell_d + \ell^+$, $m_{\ell_d} < m_{\pi^+} - m_{\ell}$ Can be light



Generating a Lepton Asymmetry

Equal and opposite dark/visible sector lepton asymmetry



$$Y_L^{\text{dark}} \equiv \left(\frac{n_{\ell_d} - n_{\bar{\ell_d}}}{s}\right) \propto \text{Br}\left(\pi^+ \to \ell_d + \ell^+\right) \sum_f A_{\text{CP}}^f \times \text{Br}\left(D^+ \to f\right)$$

Generating a Baryon Asymmetry

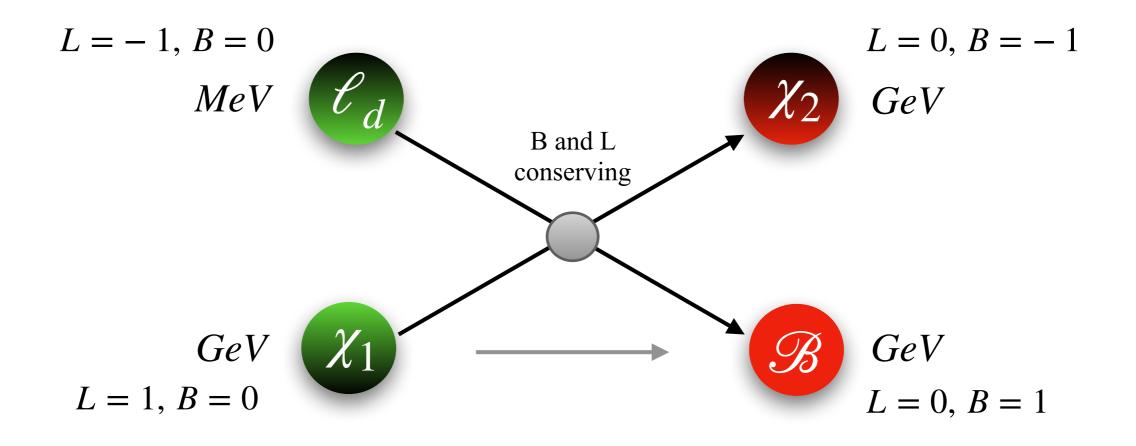
When you make the Universe at 20 MeV, you (of course) can not use Electroweak Sphalerons to transfer a lepton into a baryon asymmetry.

You also don't need them...

Freezing-In the Baryon Asymmetry

Scattering off dark sector states charged under SM L and/or B

$$\bar{\ell}_d + \chi_1 \to \chi_2 + \mathcal{B}$$



Equal and opposite dark and visible sector baryon asymmetry

Boltzmann Equations: Lepton Asymmetry

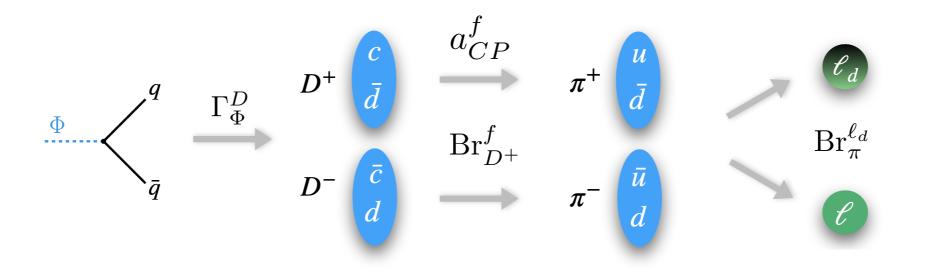
• Inflaton:

$$\frac{dn_{\Phi}}{dt} + 3Hn_{\Phi} = -\Gamma_{\Phi}n_{\Phi}$$

- Radiation: $\frac{d\rho_{\rm rad}}{dt} + 4H\rho_{\rm rad} = +\Gamma_{\Phi}m_{\Phi}n_{\Phi}$
- Hubble: $H^2 = \frac{8\pi}{3M_{\rm Pl}^2} \left(\rho_{\rm rad} + m_{\Phi} n_{\Phi} \right) \qquad \Gamma_{\Phi} = 4H \left(T_R \right)$
- The dark lepton asymmetry:

 $\Gamma_{\Phi}^{D} \equiv \Gamma_{\Phi} \mathrm{Br}(\Phi \to c) \mathrm{Br}(c \to D)$

$$\frac{d}{dt} \left(n_{\ell_d} - n_{\bar{\ell}_d} \right) + 3H \left(n_{\ell_d} - n_{\bar{\ell}_d} \right) = 2 \Gamma_{\Phi}^D n_{\Phi} \mathrm{Br}_{\pi}^{\ell_d} \sum_f N_{\pi}^f a_{CP}^f \mathrm{Br}_{D^+}^f$$



Boltzmann Equations: Lepton Asymmetry

• Inflaton: $\frac{dn_{\Phi}}{dt}$ +

$$\frac{dn_{\Phi}}{dt} + 3Hn_{\Phi} = -\Gamma_{\Phi}n_{\Phi}$$

- Radiation: $\frac{d\rho_{\rm rad}}{dt} + 4H\rho_{\rm rad} = +\Gamma_{\Phi}m_{\Phi}n_{\Phi}$
- Hubble: $H^2 = \frac{8\pi}{3M_{\rm Pl}^2} \left(\rho_{\rm rad} + m_{\Phi}n_{\Phi}\right) \qquad \Gamma_{\Phi} = 4H\left(T_R\right)$
- The dark lepton asymmetry: $\Gamma_{\Phi}^{D} \equiv \Gamma_{\Phi} Br(\Phi \to c) Br(c \to D)$

$$\frac{d}{dt} \left(n_{\ell_d} - n_{\bar{\ell}_d} \right) + 3H \left(n_{\ell_d} - n_{\bar{\ell}_d} \right) = 2 \Gamma_{\Phi}^D n_{\Phi} \mathrm{Br}_{\pi}^{\ell_d} \sum_f N_{\pi}^f a_{CP}^f \mathrm{Br}_{D^+}^f$$

Experimental Observables:

- SM charged D decays:
- Charged pion decays:

$$\begin{split} a_{CP}^{f} &\equiv A_{CP}^{f} / (1 + A_{CP}^{f}) \approx A_{CP}^{f} & LHCb, B \\ & factories \\ Br_{D^{+}}^{f} &\equiv Br \left(D^{+} \to f \right) \\ Br_{\pi}^{\ell_{d}} &\equiv Br \left(\pi^{+} \to \ell_{d} + \ell^{+} \right) & PIENU, PSI, etc. \\ & G. Elor \end{split}$$

Boltzmann Equations: Lepton Asymmetry

Numerically Solving for the Lepton asymmetry:

$$\frac{Y_L^{\text{dark}}}{Y_B^{\text{obs}}} \simeq \frac{\text{Br}_{\pi}^{\ell_d}}{10^{-3}} \frac{\sum_f N_{\pi}^f a_{CP}^f \text{Br}_{D^+}^f}{3 \times 10^{-5}} \frac{T_R}{20 \,\text{MeV}} \frac{10 \,\text{GeV}}{m_{\Phi}}$$

$$Y_B^{\text{obs}} = (8.718 \pm 0.004) \times 10^{-11}$$

Experimental Observables:

- SM charged D decays:
- Charged pion decays:

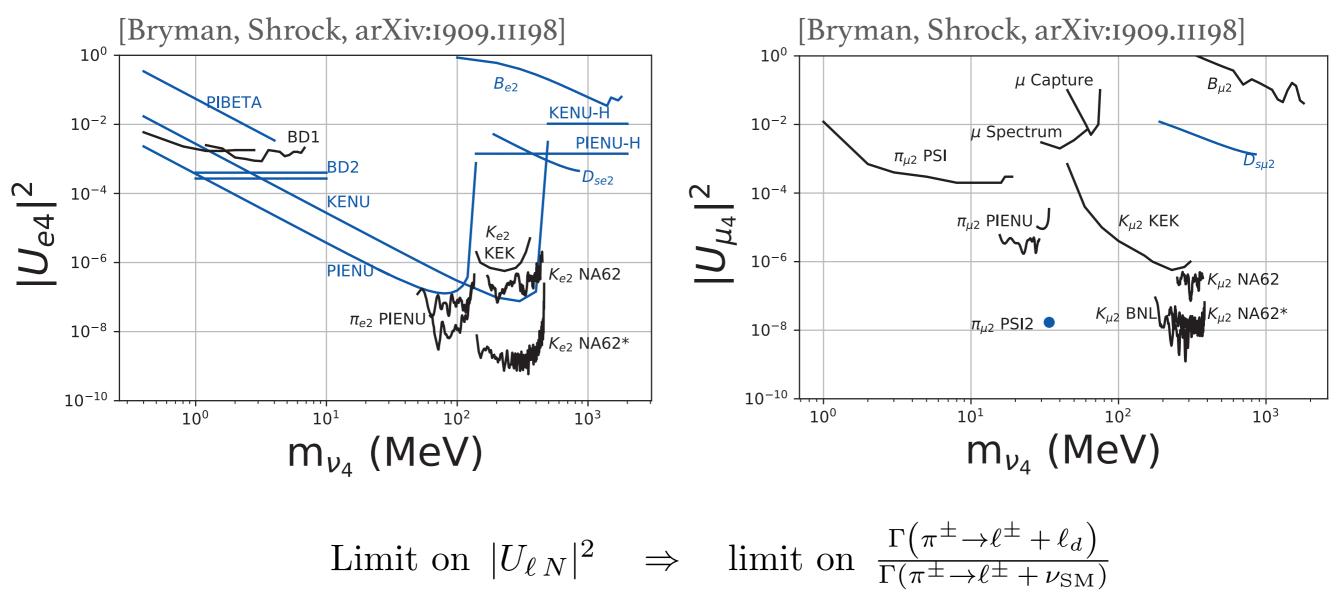
 $a_{CP}^{f} \equiv A_{CP}^{f} / (1 + A_{CP}^{f}) \approx A_{CP}^{f} \qquad LHCb, B$ factories $Br_{D^{+}}^{f} \equiv Br (D^{+} \to f)$ $Br_{\pi}^{\ell_{d}} \equiv Br (\pi^{+} \to \ell_{d} + \ell^{+}) \qquad PIENU, PSI, etc.$ G. Elor

Limits on D Decays

D^+ decay mode	$A_{CP}^{f}/10^{-2}$	${\rm Br}_{D^+}^f / 10^{-2}$
$K_S^0 \pi^+$	-0.41 ± 0.09	1.562 ± 0.031
$K^-\pi^+\pi^+$	-0.18 ± 0.16	9.38 ± 0.16
$K^-\pi^+\pi^+\pi^0$	$-0.3 \pm 0.6 \pm 0.4$	$5.98 \pm 0.08 \pm 0.16^{*}$
$K_S^0 \pi^+ \pi^0$	$-0.1 \pm 0.7 \pm 0.2$	$6.99 \pm 0.09 \pm 0.25^*$
•	•	•

$$\sum_{f} N_{\pi}^{f} a_{CP}^{f} \operatorname{Br}_{D^{+}}^{f} = \left(-9.3 \times 10^{-4}\right)_{-0.0039}^{+0.0031}$$

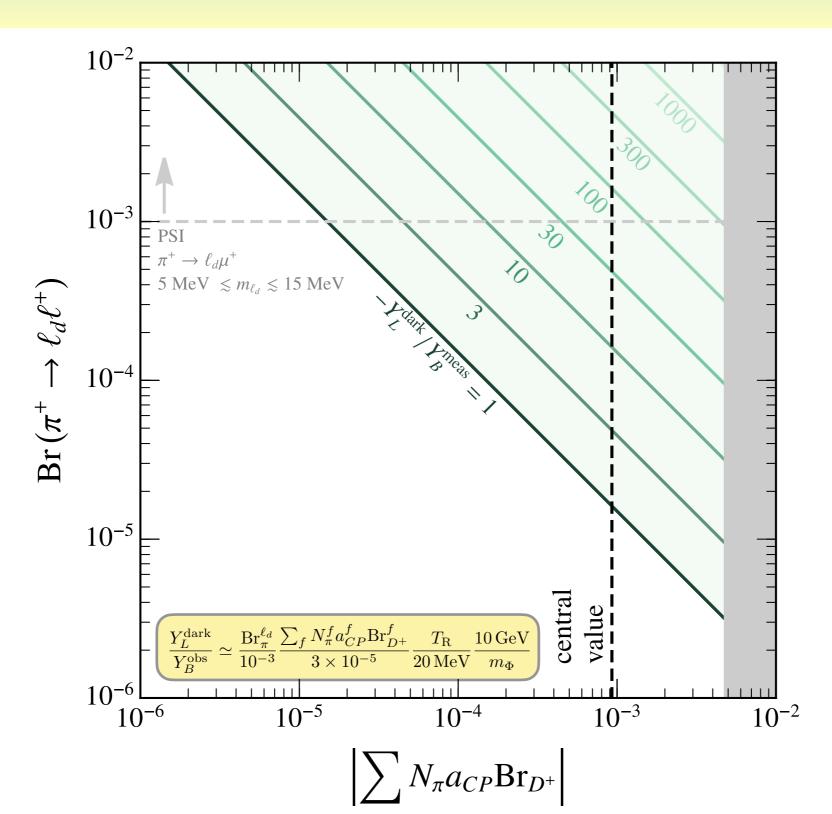
Limits on Pion Decays



[Shrock, Phys. Rev. D24, 1232 (1981)]

 $Br(\pi^{\pm} \to \mu^{\pm} + MET) \lesssim 10^{-3}$, for $5 \,MeV < m_{\ell_d} < 15 \,MeV$.

Generating a Lepton Asymmetry



Freezing-In a Baryon Asymmetry

Boltzmann Equations with scattering: $\bar{\ell}_d + \chi_1 \rightarrow \chi_2 + B$

- New dark lepton/lepto-baryon: $m_{\Phi} \gtrsim m_{\chi_1}$ $m_{\Phi} \gtrsim m_{\chi_2} + m_{\mathcal{B}}$ $\frac{dn_{\chi_1}}{dt} + 3Hn_{\chi_1} = \Gamma_{\Phi} n_{\Phi} \operatorname{Br} \left(\Phi \to \chi_1 \bar{\chi}_1\right) - \langle \sigma v \rangle n_{\bar{\ell}_d} n_{\chi_1}$
- Dark lepton:

$$\frac{d}{dt}\left(n_{\ell_d} - n_{\bar{\ell}_d}\right) + 3H\left(n_{\ell_d} - n_{\bar{\ell}_d}\right) = 2\Gamma_{\Phi}^D n_{\Phi} \mathrm{Br}_{\pi}^{\ell_d} \sum_f N_{\pi}^f a_{CP}^f \mathrm{Br}_{D^+}^f - \langle \sigma v \rangle n_{\chi_1} \left(n_{\ell_d} - n_{\bar{\ell}_d}\right)$$

• Baryon asymmetry:

$$\frac{d}{dt}\left(n_{\mathcal{B}} - n_{\overline{\mathcal{B}}}\right) + 3H\left(n_{\mathcal{B}} - n_{\overline{\mathcal{B}}}\right) = -\langle \sigma v \rangle n_{\chi_1}\left(n_{\ell_d} - n_{\overline{\ell}_d}\right)$$

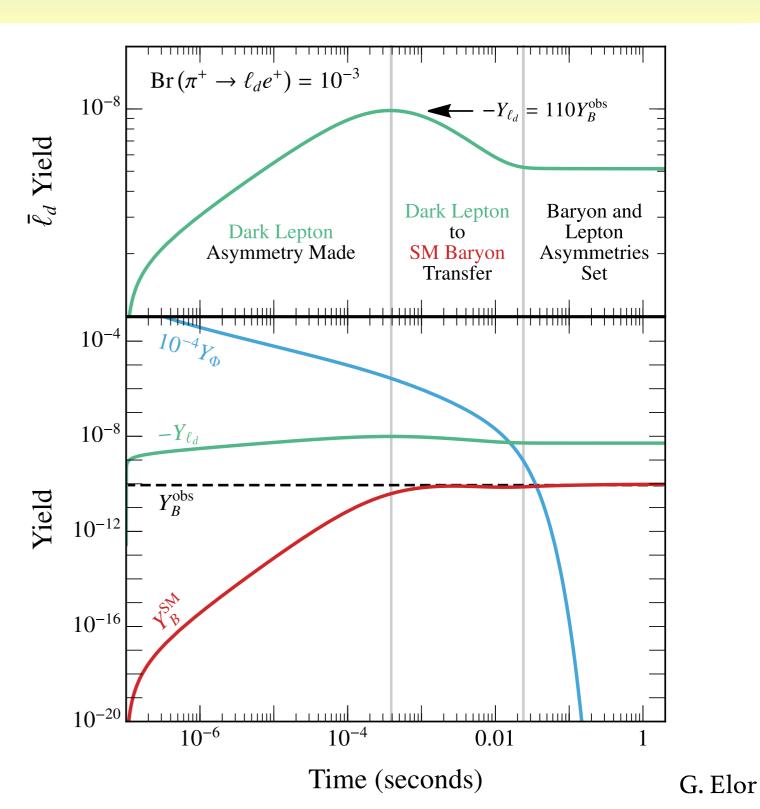
To efficiently transfer the asymmetry

$$\frac{n_{\chi_1} \langle \sigma v \rangle}{H(T)} \Big|_{T=T_R} \gtrsim \frac{Y_B^{\text{obs}}}{Y_L^{\text{dark}}}$$

Freezing-In a Baryon Asymmetry

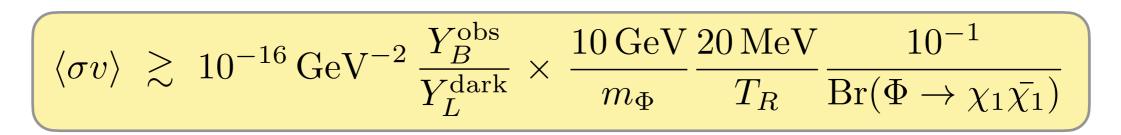
Example Benchmark point:

$$T_R = 10 \text{ MeV}, \ m_{\Phi} = 6 \text{ GeV}$$
$$\langle \sigma v \rangle = 1 \times 10^{-15} \text{ GeV}^{-2}$$
$$\operatorname{Br} \left(\Phi \to \chi_1 \bar{\chi}_1 \right) = 0.1$$
$$\sum_f N_{\pi}^f a_{CP}^f \operatorname{Br}_{D^+}^f = \left(-9.3 \times 10^{-4} \right)^{\frac{1}{2}}$$



Freezing-In a Baryon Asymmetry

Numerically:



How realistic is this?

Models

Proof of concept that what I have told you thus far is not (too) crazy.

• Some example models/dark sector charge assignments.

$$\bar{\ell}_d + \chi_1 \to \chi_2 + \mathcal{B}$$

• Estimation of the scattering cross section to confirm it can be large enough to transfer the asymmetry given current constraints.

$$\langle \sigma v \rangle \gtrsim 10^{-16} \,\mathrm{GeV}^{-2} \, \frac{Y_B^{\mathrm{obs}}}{Y_L^{\mathrm{dark}}} \times \frac{10 \,\mathrm{GeV}}{m_\Phi} \frac{20 \,\mathrm{MeV}}{T_R} \frac{10^{-1}}{\mathrm{Br}(\Phi \to \chi_1 \bar{\chi_1})}$$

Portal to the Dark Sector

Model Build for:

$$\bar{\ell}_d + \chi_1 \to \chi_2 + \mathcal{B}$$

New fields: (Same model as for *B*-Mesogenesis[arXiv:1810.00880])

	Field	Spin	L	В	\mathbb{Z}_2	Mass
Color triplet scalar mediator	Y	0	0	-2/3	+1	$\gtrsim 1{ m TeV}$
	ℓ_d	1/2	1	0	+1	$\mathcal{O}(10-140{ m MeV})$
Dark Baryon	ψ_B	1/2	0	-1	+1	$\gtrsim 1.2{ m GeV}$

 $\mathcal{L}_{\text{eff}} = \frac{y^2}{M_{-}^2} \bar{u}_i^c d_j d_k^c \psi_B$

Collider bounds (as just discussed)

Stability of matter, neutron star bounds

Allowed Interactions:

$$\mathcal{L} \supset y_{u_i d_j} Y^* \bar{u}_i d_j^c + y_{\psi d_k} Y \bar{\psi}_B d_k^c + h.c.$$

dark baryon-SM

baryon "mixing"

Dark Possibilities

$$\bar{\ell}_d + \chi_1 \to \chi_2 + \bar{\psi}_B$$

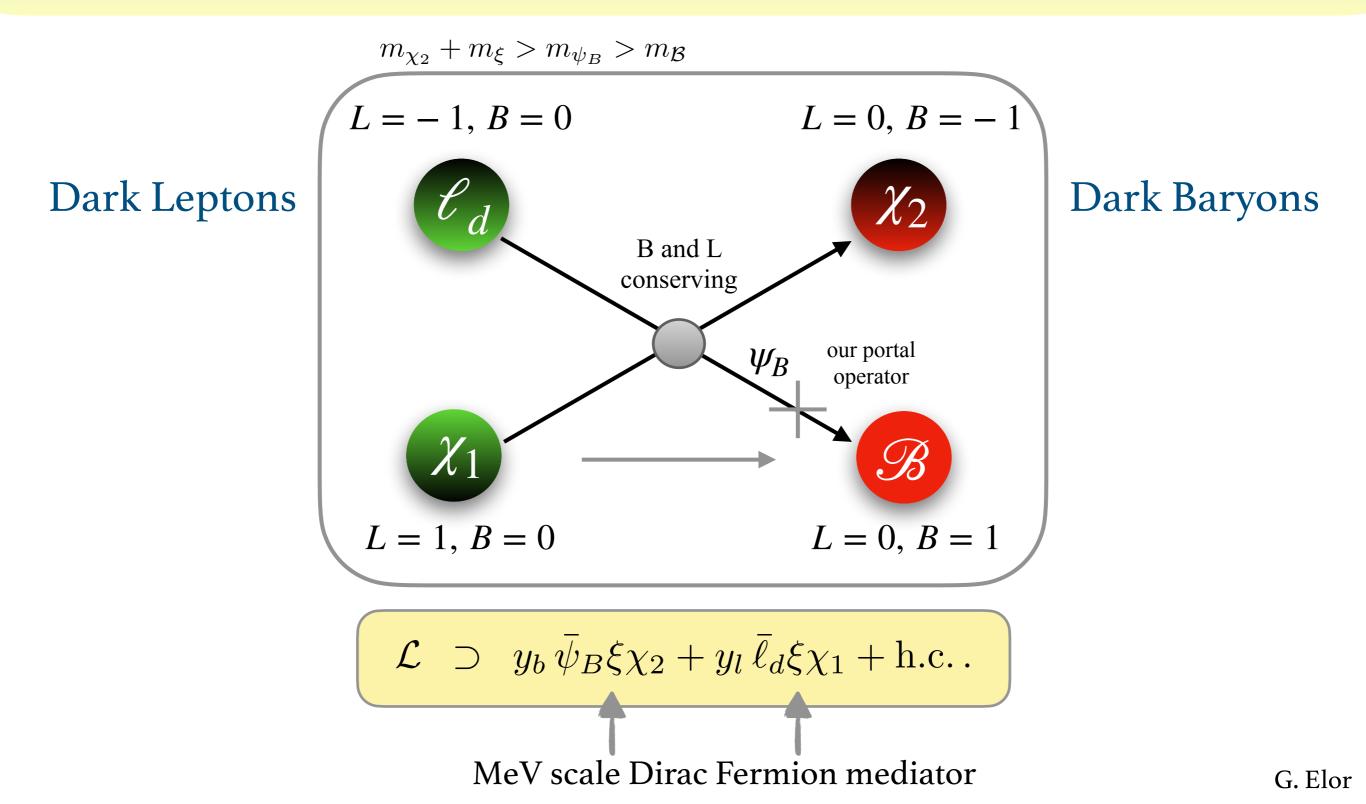
Field	L	В	Field	L	В
χ_1	1	0	χ_1	1	1
χ_2	0	-1	χ_2	0	0
χ_1	0	1	χ_1	0	0
χ_2	1	0	χ_2	-1	-1

Dark Possibilities

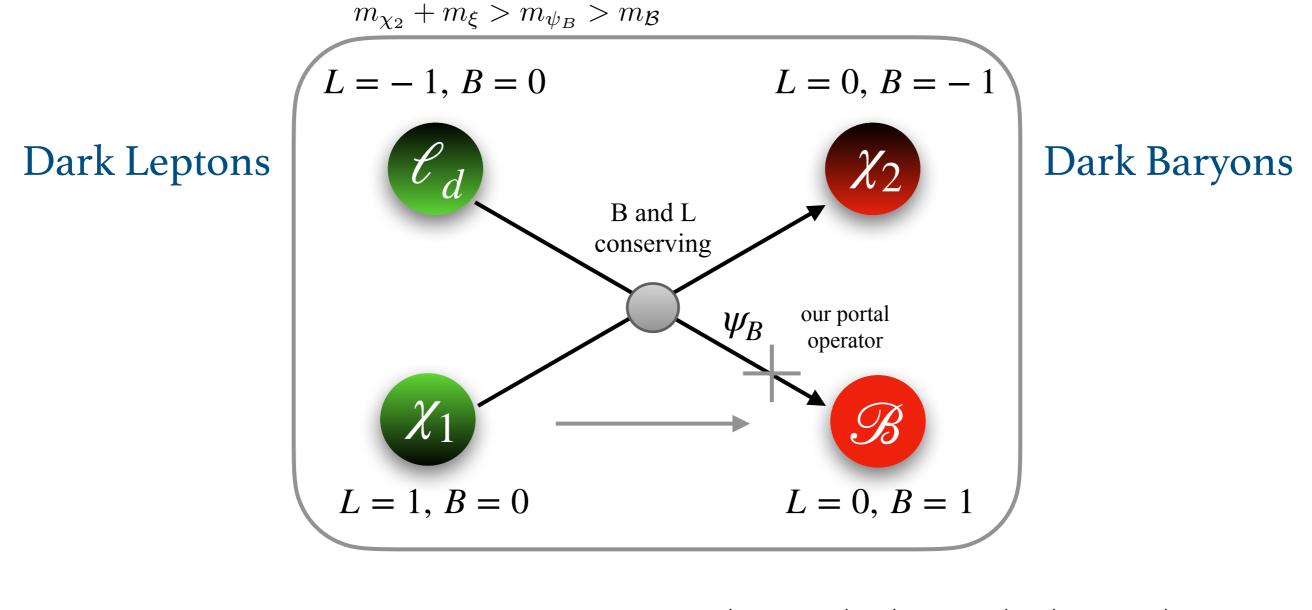
$$\bar{\ell}_d + \chi_1 \to \chi_2 + \bar{\psi}_B$$

Field	L	В	Field	L	В
χ_1	1	0	χ_1	1	1
χ_2	0	-1	χ_2	0	0
χ_1	0	1	χ_1	0	0
χ_2	1	0	χ_2	-1	-1

Example Charge Assignment



Example Charge Assignment



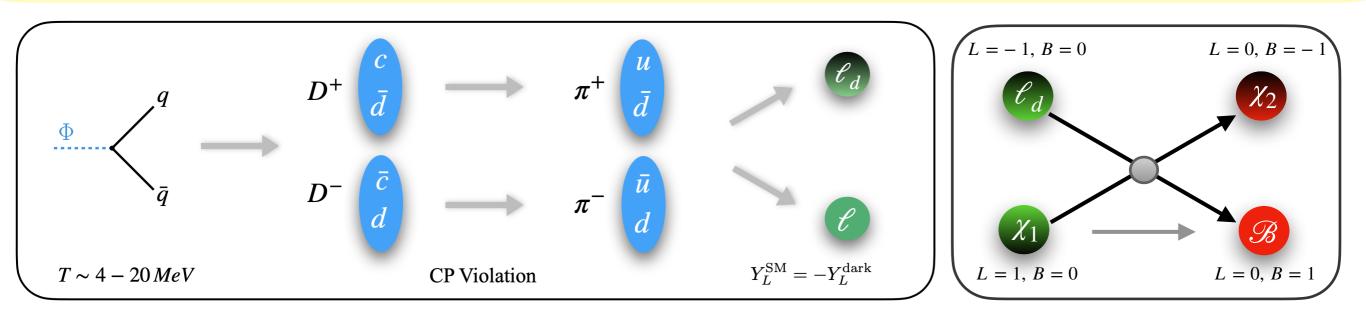
$$\langle \sigma v \rangle \simeq 10^{-15} \,\mathrm{GeV}^{-2} \, (y_l \, y_b)^2 \, \times \left(\frac{10 \,\mathrm{MeV}}{m_{\ell_d}}\right) \left(\frac{20 \,\mathrm{GeV}}{m_{\chi_1}}\right) \left(\frac{10 \,\mathrm{GeV}}{m_{\chi_2}}\right)$$

Dark Matter Generation

A few comments

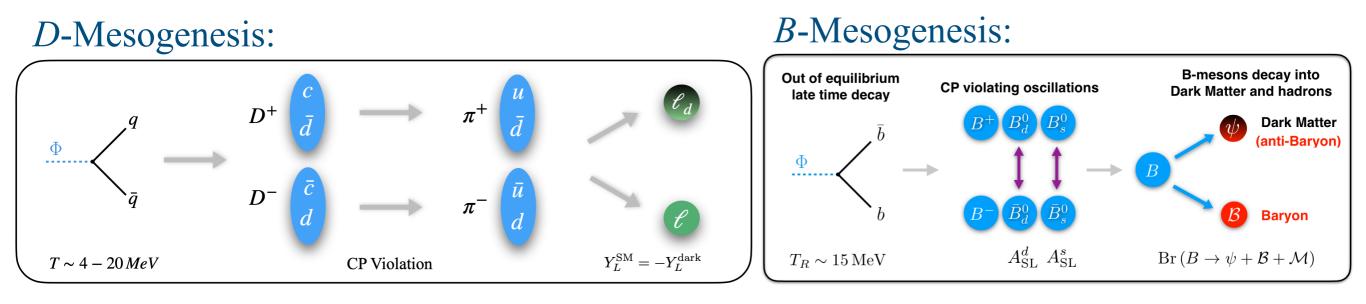
- Details depend on baryon and lepton number charge assignment.
- Symmetric component of any dark baryons will always contribute to some or all of the dark matter.
- Dark states will tend to be overproduced but additional dark sector dynamics can always be at play to achieve the observed dark matter relic abundance.

D-Mesogenesis



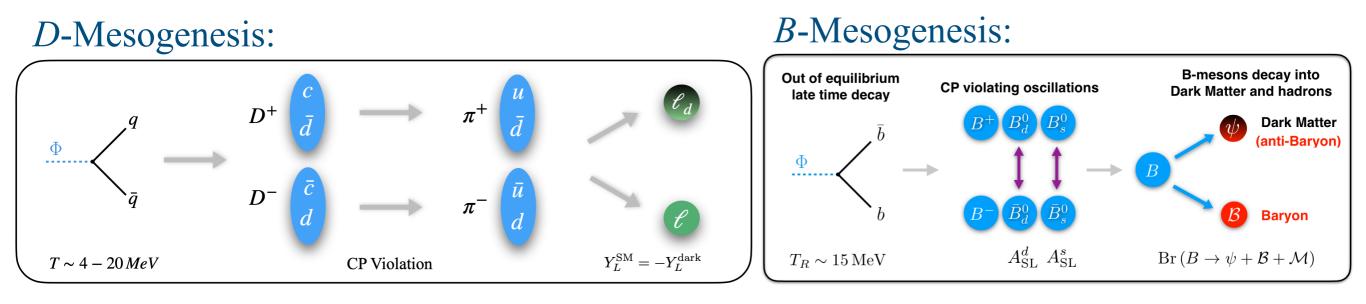
- First generates a lepton asymmetry and then freezes in a baryon asymmetry through dark sector scatterings.
- Baryogenesis and dark matter production are controlled by experimental observables of the charged *D* Mesons system.
- Upcoming experimental probes will better constrain or discover this mechanism.

Outlook



- Continued support of experimental efforts to probe Mesogenesis e.g. proper form factor calculation in the case of decays relevant for *B*-Mesogenesis.
- Mesogenesis in other meson systems?
- Explore UV embeddings (dark sector? mediators?) of both *B* and *D*-Mesogenesis and associated phenomenology.
- Theory of inflation and flavor.

Outlook



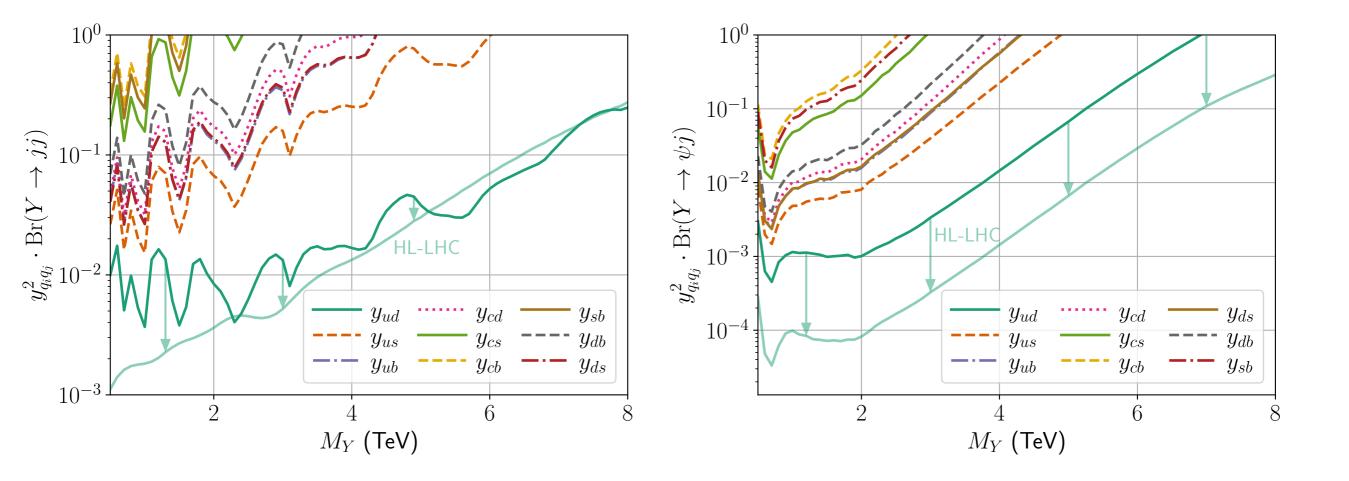
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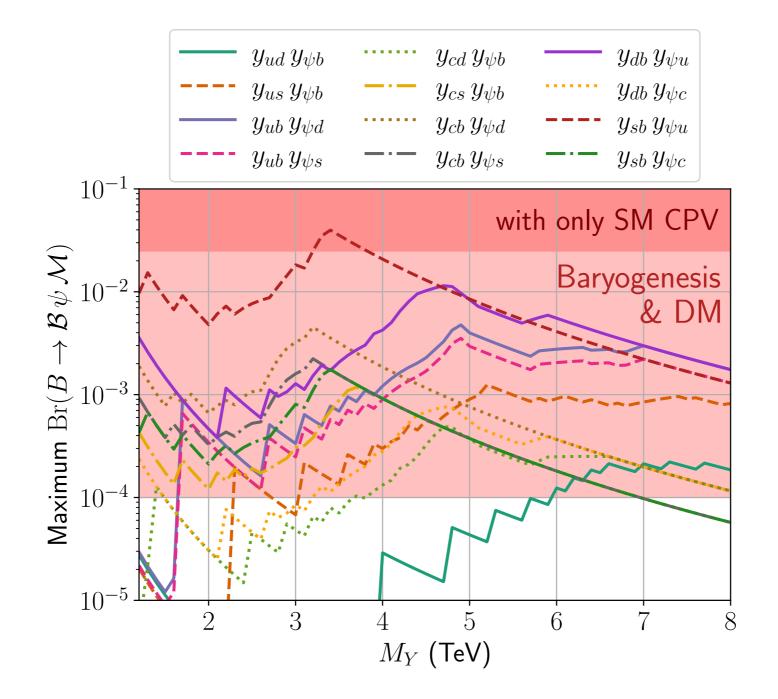
Colored Triplet Scalar

Constraints from LHC squark searches

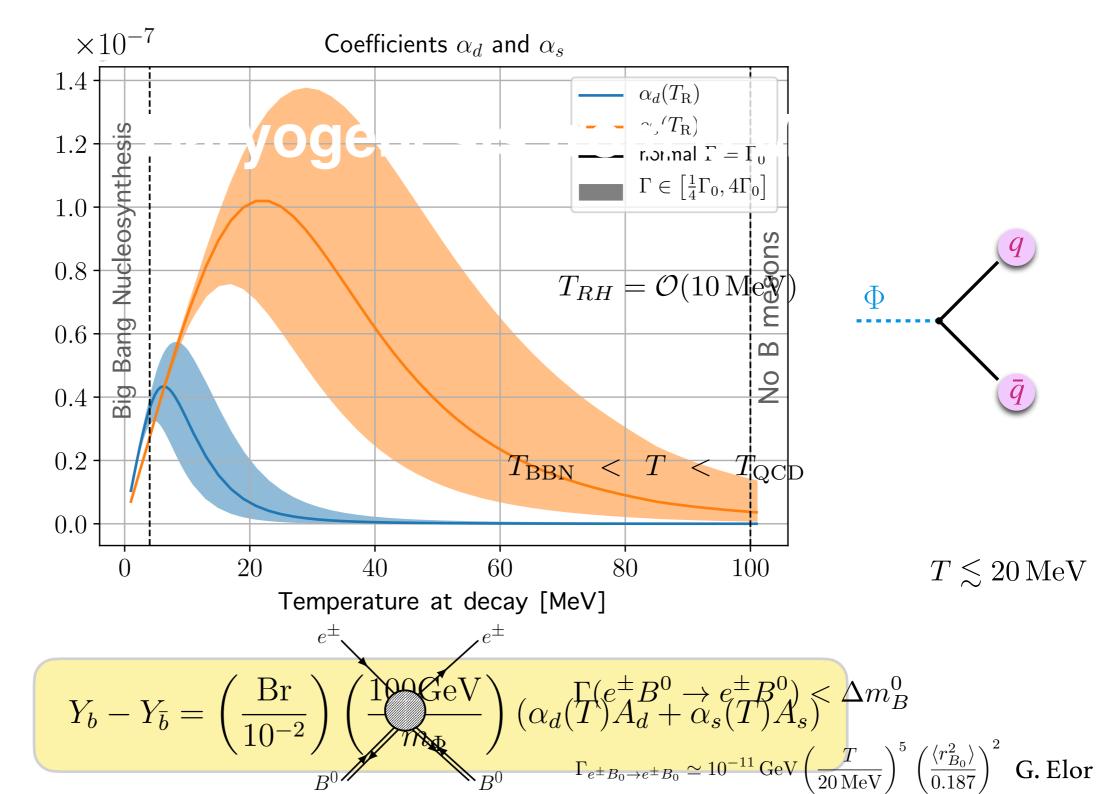


Colored Triplet Scalar

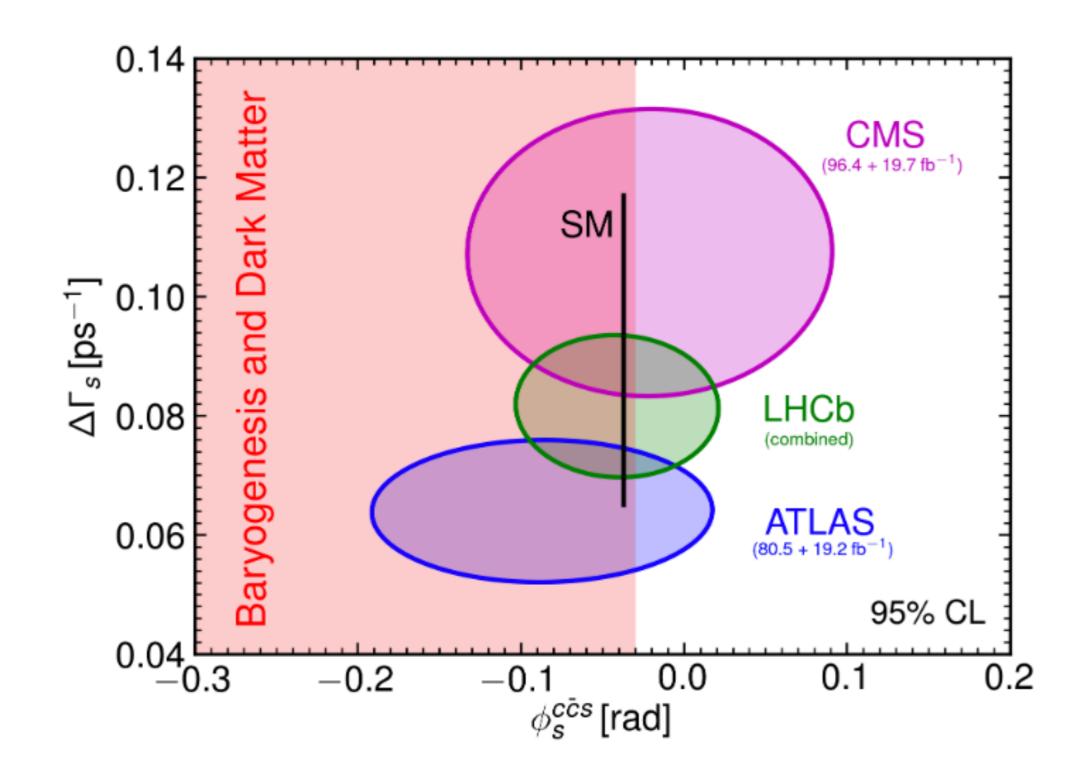
Constraints from LHC squark searches



Baryogenesis and Dark Matter from B Mesons



CP Observables



The Semi-Leptonic Asymmetry

$$A_{\rm SL}^q = \frac{\Gamma\left(\bar{B}_q^0 \to B_q^0 \to f\right) - \Gamma\left(B_q^0 \to \bar{B}_q^0 \to \bar{f}\right)}{\Gamma\left(\bar{B}_q^0 \to B_q^0 \to f\right) + \Gamma\left(B_q^0 \to \bar{B}_q^0 \to \bar{f}\right)} = -\left|\frac{\Gamma_{12}^q}{M_{12}^q}\right| \sin\left(\phi_{12}^q\right)$$

From SM box diagrams:

 $A_{\rm SL}^d|_{\rm SM} = (-4.7 \pm 0.4) \times 10^{-4}$ $A_{\rm SL}^s|_{\rm SM} = (2.1 \pm 0.2) \times 10^{-5}$

Lenz, Tetlalmatzi-Xolocotzi [1912.07621]

World average: As

$${}^{d}_{\rm SL} = (-2.1 \pm 1.7) \times 10^{-3}$$

 ${}^{s}_{\rm SL} = (-0.6 \pm 2.8) \times 10^{-3}$

HFLAG

Projected Sensitivities:

$$\begin{split} &\delta A_{\rm SL}^s = 10 \times 10^{-4} \; \left[{\rm LHCb} \left({33\,{\rm fb}^{-1}} \right) - 2025 \right] & [1812.07638, \\ &\delta A_{\rm SL}^s = 3 \times 10^{-4} \; \left[{\rm LHCb} \left({300\,{\rm fb}^{-1}} \right) - 2040 \right] & 1808.08865 \right] \\ &\delta A_{\rm SL}^d = 8 \times 10^{-4} \; \left[{\rm LHCb} \left({33\,{\rm fb}^{-1}} \right) - 2025 \right] \\ &\delta A_{\rm SL}^d = 2 \times 10^{-4} \; \left[{\rm LHCb} \left({300\,{\rm fb}^{-1}} \right) - 2040 \right] \\ &\delta A_{\rm SL}^d = 5 \times 10^{-4} \; \left[{\rm Belle} \; {\rm II} \left({50\,{\rm ab}^{-1}} \right) - 2025 \right] \end{split}$$

Searching for new b-Hadron Decays

Caution: Inclusive vs. Exclusive Rates

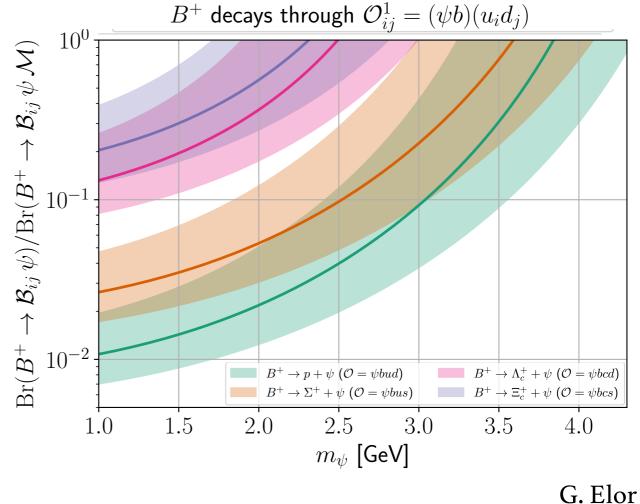
- All decays (and their searches) discussed thus far have been *exclusive*. But, the observable controlling the baryon asymmetry is an *inclusive* rate.
 - $Br(B \to \psi \,\mathcal{B} \,\mathcal{M}) \gtrsim 10^{-4}$
- Need a dedicated calculation using QCD sum rules or lattice techniques etc. to calculate form factors. Beyond my current expertise....
- Phase space method

$$\begin{bmatrix} \text{Bigi, Phys.Lett.B 106, 510 (1981)} \end{bmatrix}$$

$$\bar{b} \to u_i d_j \psi \qquad \gamma(\Lambda) \equiv \int_{(m_{u_i} + m_{d_j})^2}^{\Lambda^2} \frac{\partial \Gamma}{\partial M_{u_i d_j}^2} \, \mathrm{d}M_{u_i d_j}^2$$

$$\frac{\mathrm{Br}(B \to \mathcal{B}_{ij} + \psi)}{\mathrm{Br}(B \to \mathcal{B}_{ij} + \psi + \mathcal{M})} \simeq \frac{\gamma(m_{\mathcal{B}_{ij}})}{\gamma(m_b - m_{\psi})}$$

$$\frac{\mathrm{Br}(B \to \psi \mathcal{B})}{\mathrm{Br}(B \to \psi \mathcal{B} \mathcal{M})} \gtrsim (1 - 10) \%.$$



Flavorful Variations

No a priori reason to expect a particular flavor structure.

Most general interactions:

$$\mathcal{L}_{-1/3} = -\sum_{i,j} y_{u_i d_j} Y^* \bar{u}_{iR} d_{jR}^c - \sum_k y_{\psi d_k} Y d_{kR}^c \bar{\psi} + \text{h.c.}$$
$$\mathcal{L}_{2/3} = -\sum_{i,j} y_{d_i d_j} Y^* \bar{d}_{iR} d_{jR}^c - \sum_k y_{\psi u_k} Y u_{kR}^c \bar{\psi} + \text{h.c.}$$

Possible operators:

$$\mathcal{O}_{ud} = \psi \, b \, u \, d$$
$$\mathcal{O}_{us} = \psi \, b \, u \, s$$
$$\mathcal{O}_{cd} = \psi \, b \, c \, d$$
$$\mathcal{O}_{cs} = \psi \, b \, c \, s$$

New meson and b-flavored hadron decays:

Operator/Decay	Initial State	Final state	$\Delta M \ ({\rm MeV})$
$\mathcal{O} = \psi b u d$	B_d B_s	$egin{aligned} \psi + n (udd) \ \psi + \Lambda (uds) \end{aligned}$	$\begin{array}{c} 4340.07 \\ 4251.21 \end{array}$
$\overline{b} \to \psi u d$	B^+ Λ_b	$\psi + p \left(duu ight) \ ar{\psi} + \pi^0$	4341.05 5484.5
	B_d	$\psi + \Lambda(usd)$	4163.95
$\mathcal{O} = \psi b u s$	B_s B^+	$\psi + \Xi^0 (uss)$	4025.03
$b \rightarrow \psi u s$	$\frac{B}{\Lambda_b}$	$\frac{\psi + \Sigma^+ (uus)}{\bar{\psi} + K^0}$	4089.95 5121.9
	B_d	$\psi + \Lambda_c + \pi^- \left(c d d \right)$	2853.60
$\mathcal{O} = \psi b c d$	B_s	$\psi + \Xi_c^0 (cds)$	2895.02
$b \rightarrow \psi c d$	B^+	$\psi + \Lambda_c (dcu)$	2992.86
	Λ_b	$ar{\psi}+\overline{D}^0$	3754.7
	B_d	$\psi + \Xi_c^0 (csd)$	2807.76
$\mathcal{O} = \psi b c s$	B_s	$\psi + \Omega_c \ (css)$	2671.69
$\overline{b} o \psi c s$	B^+	$\psi + \Xi_c^+ (csu)$	2810.36
	Λ_b	$\bar{\psi} + D^- + K^+$	3256.2

Colored Triplet Scalar

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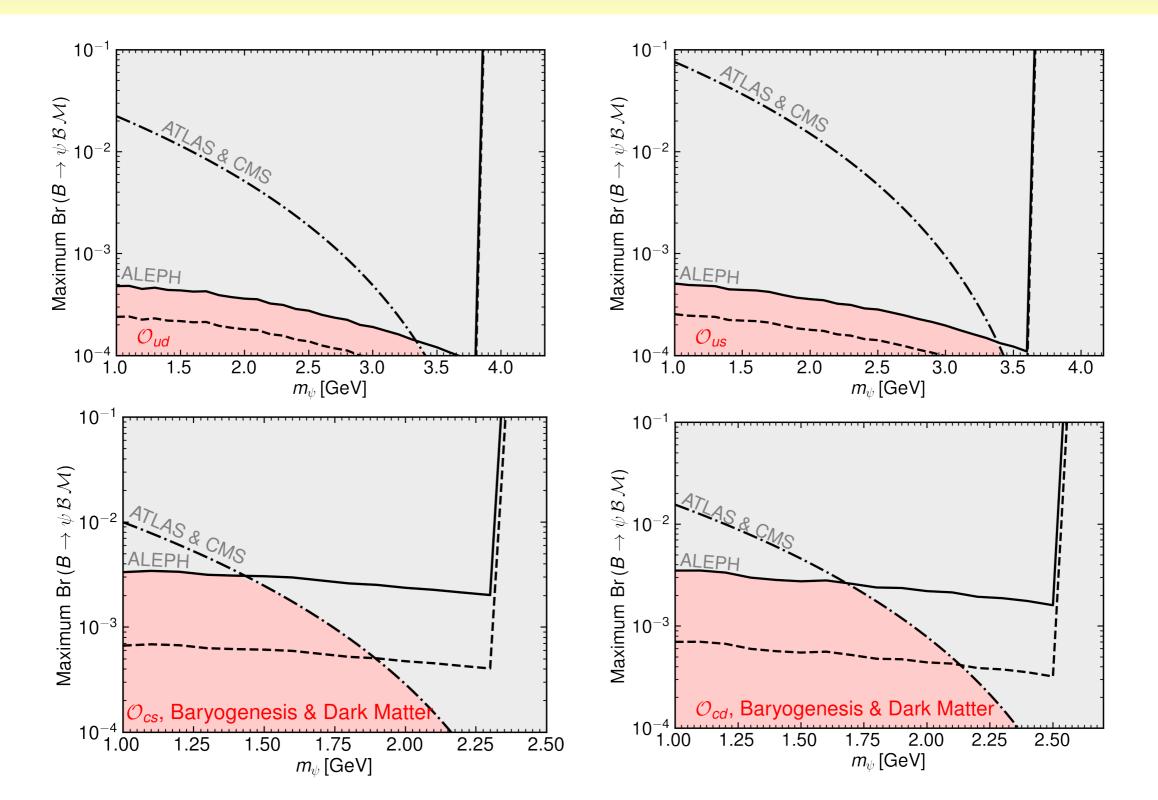
	Max:	Upper limit on m_{ψ}				
Operator	$\left[\sqrt{y^2} \underline{1.5 \mathrm{TeV}} \right]^4$	Inclusive $\operatorname{Br}(B \to \psi \mathcal{B} \mathcal{M})$				
	$\left\lfloor \overline{0.53} \ \overline{M_Y} \right\rfloor$	10^{-4}	10^{-3}	10^{-2}	2.5%	
$\mathcal{O}_{ud}^1 = (\psi b) (u d)$	0.3	2.0	_	-	-	
$\mathcal{O}_{ud}^2 = (\psi d) (u b)$	6.7	3.2	2.3	-	-	
$\mathcal{O}_{ud}^3 = (\psi u) (d b)$	16.2	3.4	2.8	1.6	-	
$\mathcal{O}_{us}^1 = (\psi b) (u s)$	2.4	2.7	1.6	-	-	
$\mathcal{O}_{us}^2 = \left(\psi s\right) \left(u b\right)$	6.7	3.2	2.3	-	-	
$\mathcal{O}_{us}^3 = (\psi u) (s b)$	75.8	3.5	3.0	2.2	1.8	
$\mathcal{O}_{cd}^1 = (\psi b) (c d)$	10.4	2.0	1.2	-	-	
$\mathcal{O}_{cd}^2 = (\psi d) (c b)$	96.6	2.4	1.9	1.2	-	
$\mathcal{O}_{cd}^{3} = \left(\psi c\right) \left(d b\right)$	16.2	2.1	1.4	-	-	
$\mathcal{O}_{cs}^1 = (\psi b) (c s)$	50.9	2.0	1.5	-	-	
$\mathcal{O}_{cs}^{2} = \left(\psi s\right) \left(c b\right)$	96.6	2.4	1.9	1.2	-	
$\mathcal{O}_{cs}^{3} = \left(\psi c\right) \left(s b\right)$	75.8	2.1	1.7	-	-	

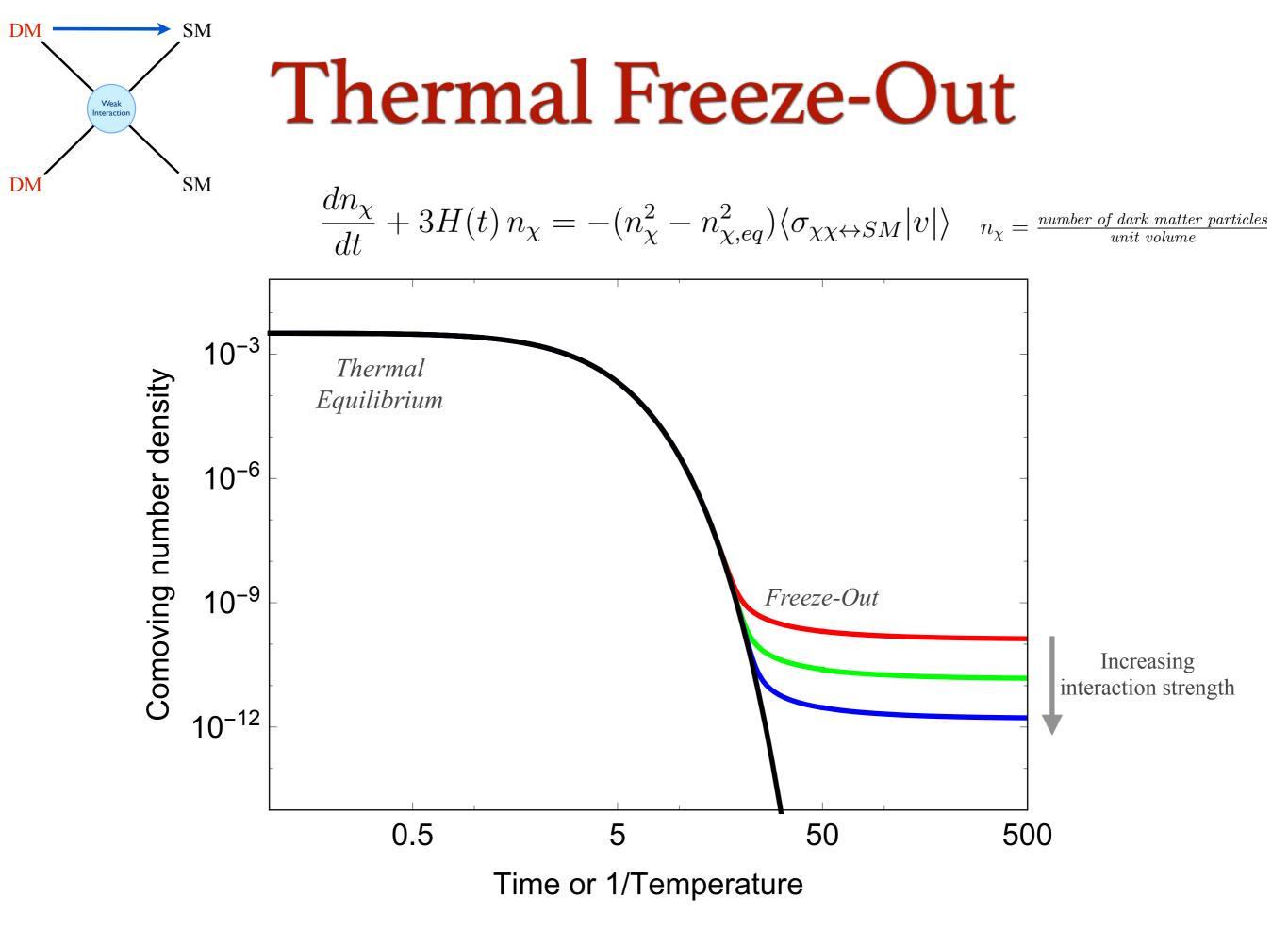
Bounds will depend on the product of coupling

$$\sqrt{y^2}/M_Y$$

Which in tern constrain branching fraction of relevance to the baryon asymmetry G. Elor

Colored Triplet Scalar





A Supersymmetric Theory

MSSM, R Symmetry, and Dirac Gauginos and Sterile Neutrios

Superfield	R-Charge	L no.
$\mathbf{U}^c, \mathbf{D}^c$	2/3	0
Q	4/3	0
$\mathbf{H}_{u},\mathbf{H}_{d}$	0	0
$\mathbf{R}_u, \mathbf{R}_d$	2	0
S	0	0
L	1	1
\mathbf{E}^{c}	1	-1
\mathbf{N}_{R}^{c}	1	-1

"RPV"
$$\mathbf{W} = y_u \mathbf{Q} \mathbf{H}_u \mathbf{U}^c - y_d \mathbf{Q} \mathbf{H}_d \mathbf{D}^c - y_e \mathbf{L} \mathbf{H}_d \mathbf{E}^c + \frac{1}{2} \lambda_{ijk}^{"} \mathbf{U}_i^c \mathbf{D}_j^c \mathbf{D}_k^c$$

 $+ \mu_u \mathbf{H}_u \mathbf{R}_d + \mu_d \mathbf{R}_u \mathbf{H}_d$
 $+ \lambda_u^t \mathbf{H}_u \mathbf{T} \mathbf{R}_d + \lambda_d^t \mathbf{R}_u \mathbf{T} \mathbf{H}_d + \lambda_d^s \mathbf{S} \mathbf{R}_u \mathbf{H}_d$.
 $\boldsymbol{\mathcal{L}} \stackrel{:}{=} \lambda_{113}^{"} \left(\tilde{d}_R^* u_R^\dagger b_R^\dagger + \tilde{u}_R^* d_R^\dagger b_R^\dagger + \tilde{b}_R^* u_R^\dagger d_R^\dagger \right) ,$
Gauge:
 $\mathcal{L}_{gauge} = -\sqrt{2}g(\phi T^a \psi^\dagger) \lambda^{a\dagger} + h.c.$

 $\Rightarrow -\sqrt{2}g(\tilde{d}_R^* d_R \tilde{B}^\dagger) - \sqrt{2}g(\tilde{d}_L d_L^\dagger \tilde{B}^\dagger) + \text{h.c.}$

Neutrio:

$$\mathbf{W} = \frac{\lambda_N}{4} \mathbf{S} \mathbf{N}_R^c \mathbf{N}_R^c + \mathbf{H}_u \mathbf{L}^i y_N^{ij} \mathbf{N}_R^{c,j} + \frac{1}{2} \mathbf{N}_R^c M_M \mathbf{N}_R^c + \text{h.c.},$$
$$\mathbf{W} = \frac{\lambda_N}{4} \mathbf{N}_R (\lambda_s \nu_R^\dagger \tilde{\nu}_R^* + \phi_s \nu_R^\dagger \nu_R^\dagger) + \text{h.c.}$$

Parameter space: "RPV" couplings and squark mass mixing

A Supersymmetric Theory

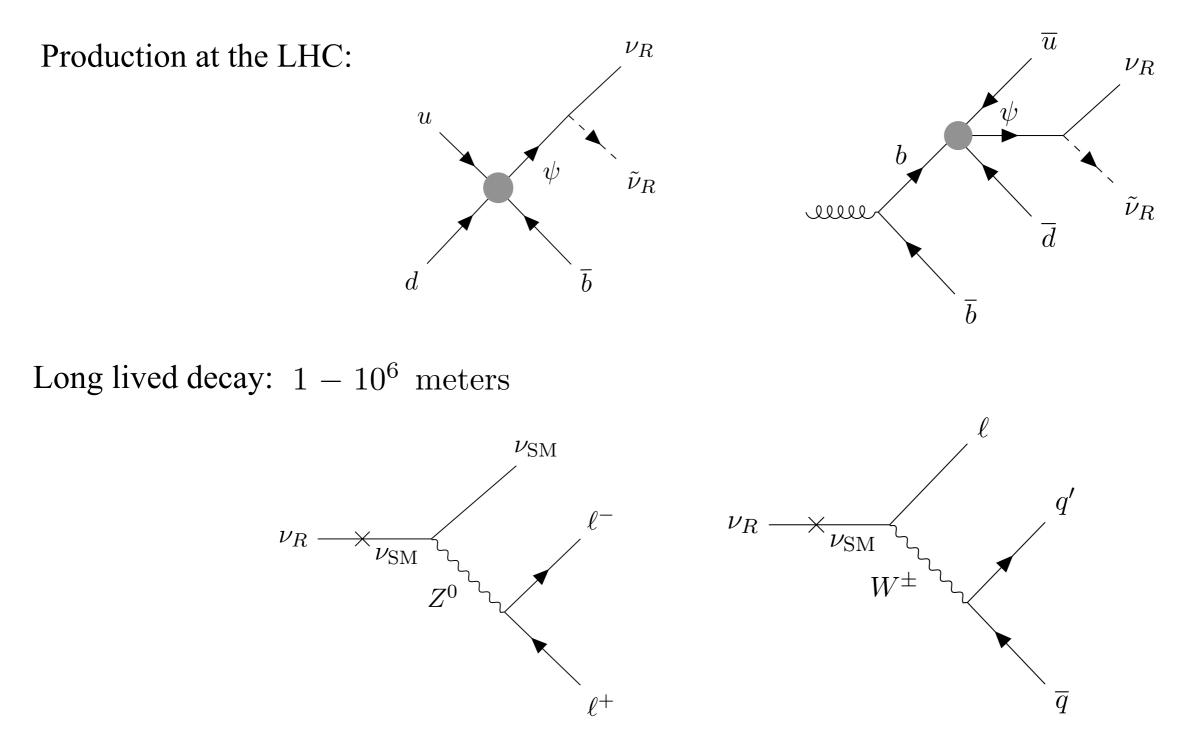
Superpartners and SM particles have different charge under an unbroken R-symmetry. We can identify this with Baryon number.

-----> Super

Superpartners as dark baryons.

	Field	Spin	Q_{EM}	Baryon no.	\mathbb{Z}_2	Mass
	Φ	0	0	0	+1	$11 - 100 \mathrm{GeV}$
MSSM Squark	\tilde{d}_R	0	-1/3	-2/3	+1	$\mathcal{O}({ m TeV})$
Dirac Bino	$\left[\begin{array}{c} \tilde{B} \\ \lambda_s^{\dagger} \end{array}\right]$	1/2	0	-1	+1	$\mathcal{O}({ m GeV})$
Right handed	$ u_R $	1/2	0	0	-1	$\mathcal{O}({ m GeV})$
neutrino multiplet	$\tilde{ u}_R$	0	0	-1	-1	$\mathcal{O}({ m GeV})$

Model Specific Signals



[G. Alonso-Alvarez, G. Elor, A. E. Nelson, H. Xiao JHEP [arXiv:1907.10612]] [G. Alonso-Alvarez, G. Elor, M. Escudero, D. McKeen [*in preparation*]]