# Mesogenesis

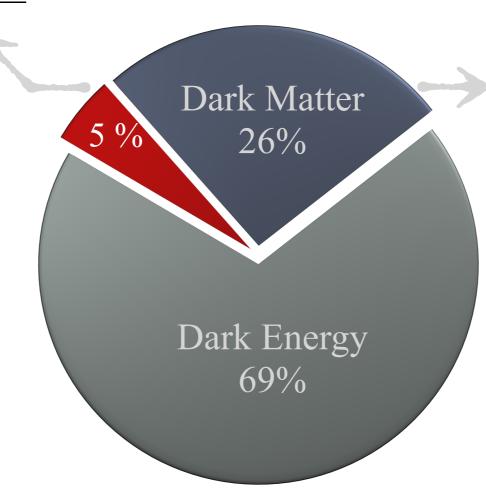
### Gilly Elor

MITP Particle Theory Seminar June 30 2021

### What is the Universe Made of?

#### From cosmological measurements we know:

The stuff we understand stars, planets, you (baryonic matter) Only 5 %

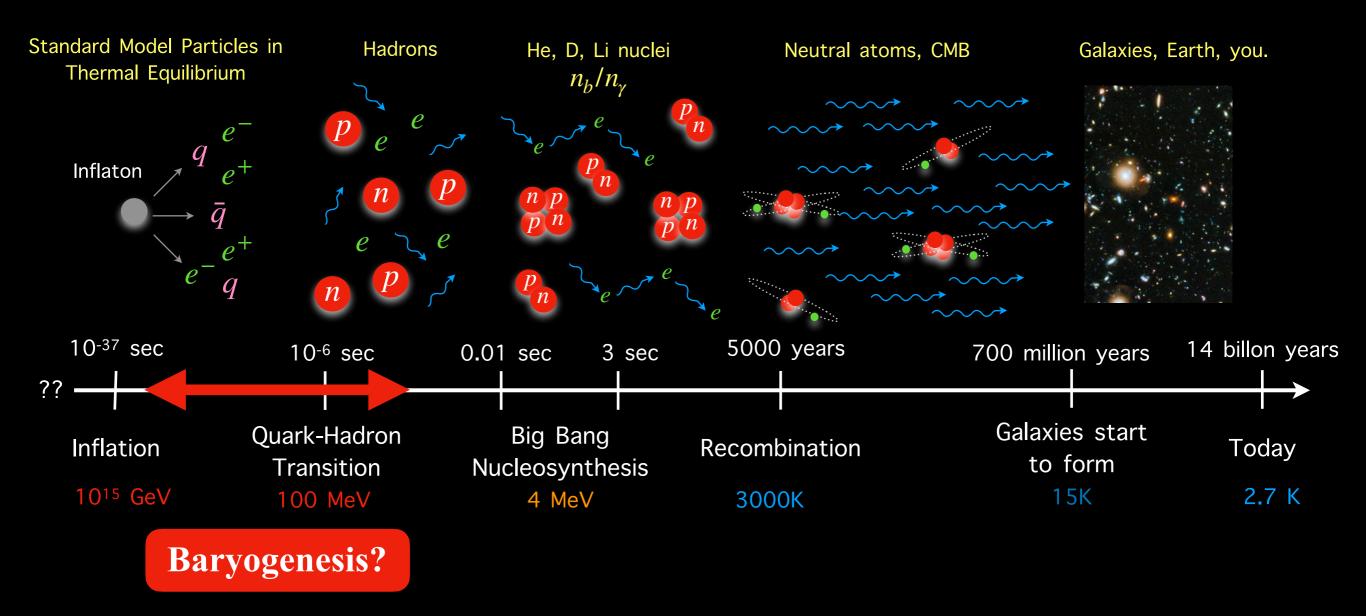


What is the nature and origin of dark matter?

Energy density today

We don't know where the 5 % of baryonic matter came from.

### The History of the Universe

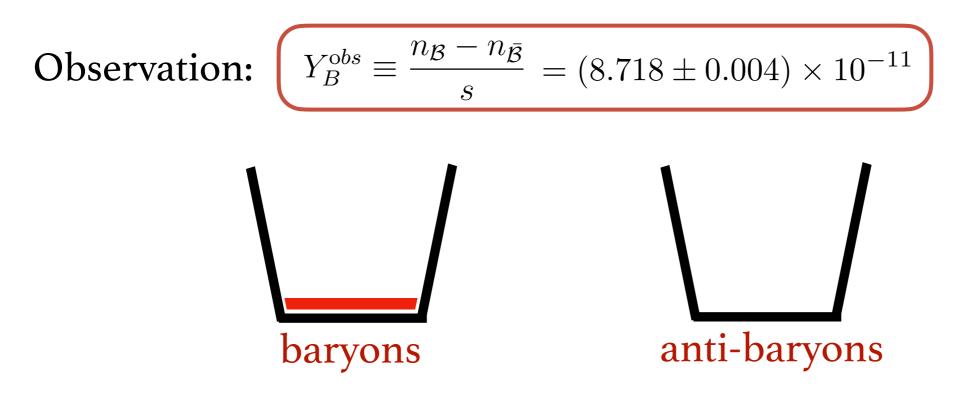


What mechanism generated the initial asymmetry? Observed to be (BBN, CMB):

$$Y_B^{\text{obs}} \equiv \frac{n_{\mathcal{B}} - n_{\bar{\mathcal{B}}}}{s} \sim 8 \times 10^{-11}$$

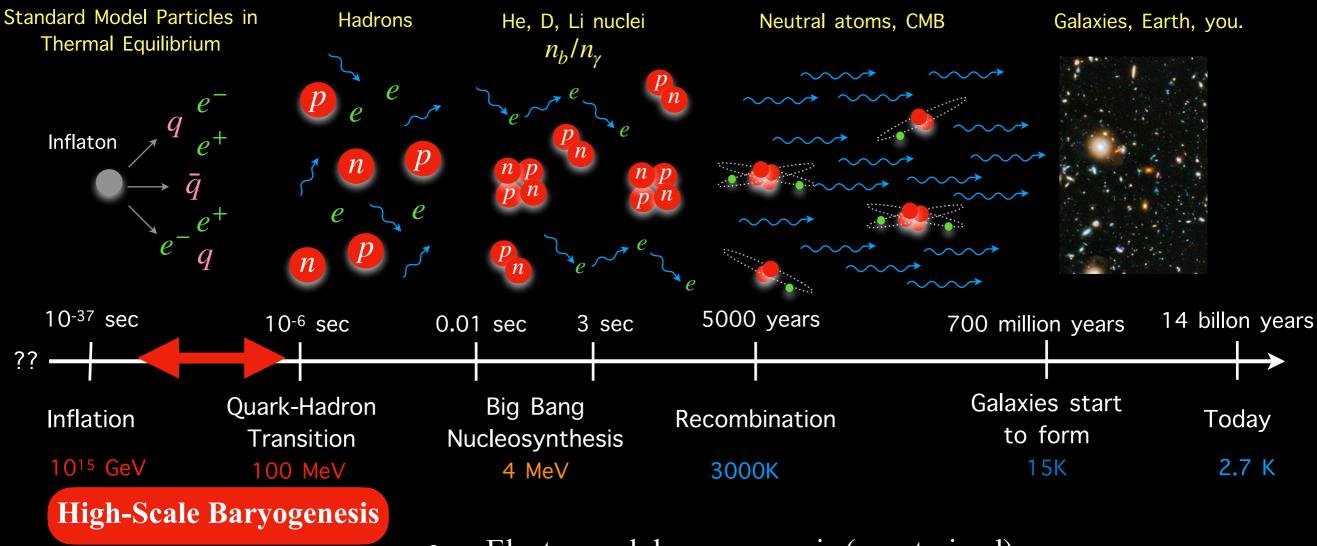
## Baryogenesis

Review: how to generate a baryon asymmetry?



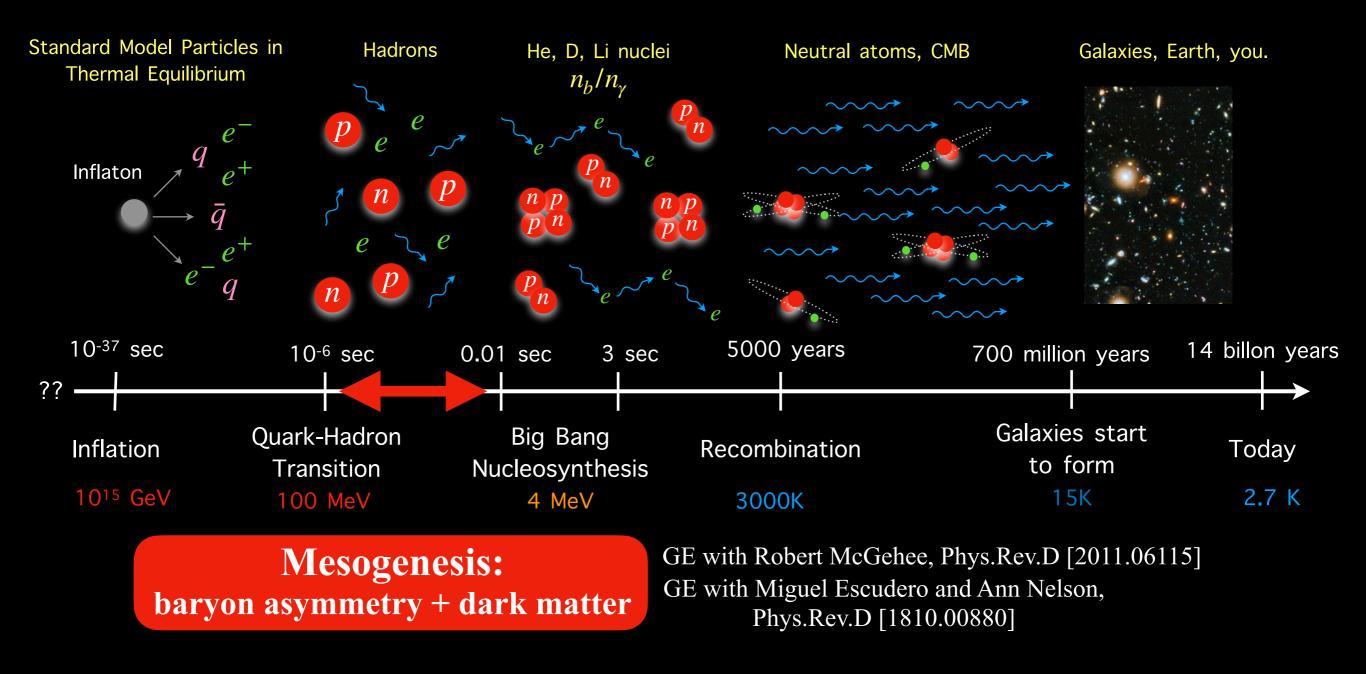
- Baryon number violation.
- Conjugate rates must be different.
- Out of thermal equilibrium.

### "Traditional" Baryogenesis



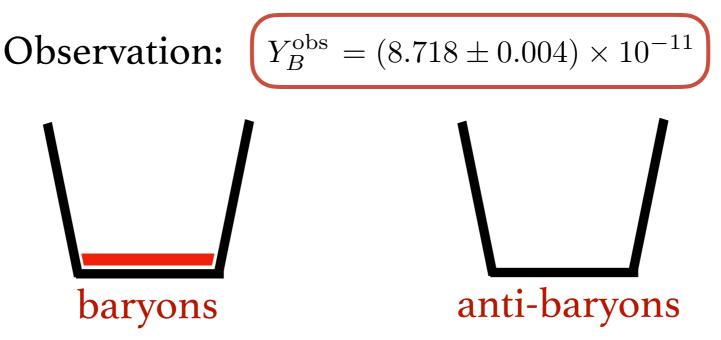
- Electroweak baryogengesis (constrained)
- Leptogenesis (hard to test)
- Affleck-Dine (very hard to test)
- • • • •

### Making the Universe at 20MeV



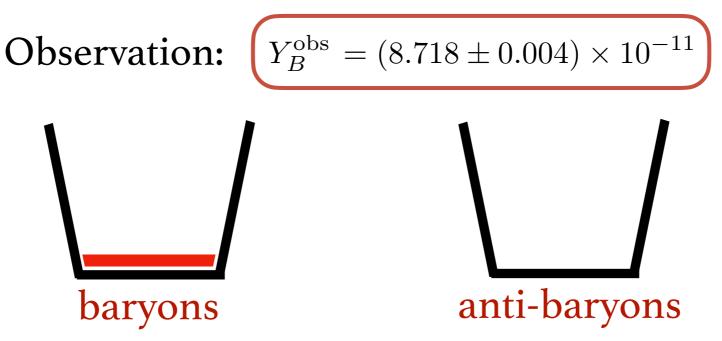
- Controlled by experimental observables. Signals!
- Theoretically appealing e.g. Relaxion and Nnaturalness require low scale baryogenesis.





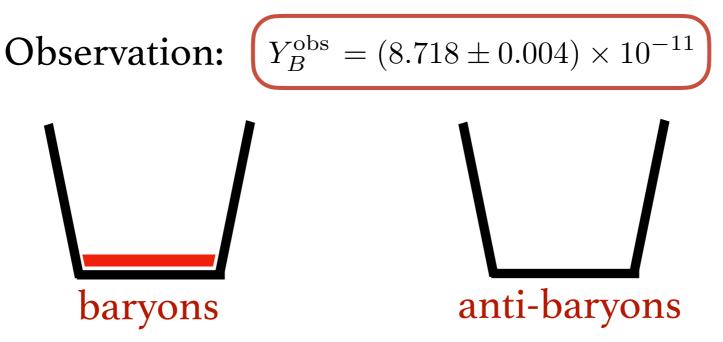
- Out of thermal equilibrium:
- CP Violation:
- "Baryon number violation":





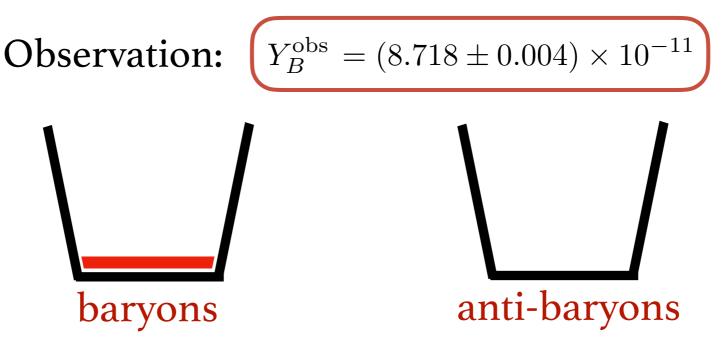
- Out of thermal equilibrium: Late decays of "inflaton" field to SM Mesons.
- CP Violation:
- "Baryon number violation":





- Out of thermal equilibrium: Late decays of "inflaton" field to SM Mesons.
- CP Violation: In SM Meson systems.
- "Baryon number violation":





- Out of thermal equilibrium: Late decays of "inflaton" field to SM Mesons.
- CP Violation: In SM Meson systems.
- "Baryon number violation": SM Meson decays to dark leptons or baryons.

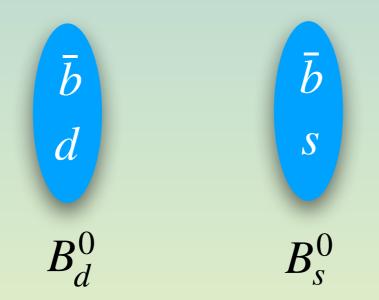
## Today's Talk

Baryogenesis and Dark Matter production using the CP Violation of Standard Model Meson Systems

**Observation:**  $Y_B^{\text{obs}} = (8.718 \pm 0.004) \times 10^{-11}$ 

- Part I. *B*-Mesogenesis: CP violation from neutral B meson oscillations [GE with Miguel Escudero and Ann Nelson arXiv:1810.00880].
- Part II. Discovering *B*-Mesogenesis: Details of signals and searches [GE with Gonzalo Alonso-Alvarez and Miguel Escudero arXiv:2101.02706, ongoing work, both theory and experimental].
- Part III. *D*-Mesogenesis: CP violation from charged *D* meson decays [GE with Robert McGehee arXiv:2011.06115].

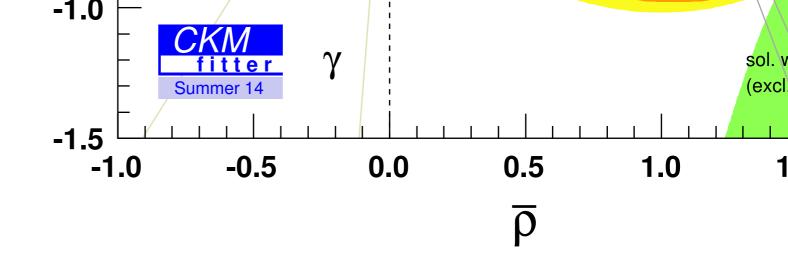
# Part I. B-Mesogenesis



Based on:

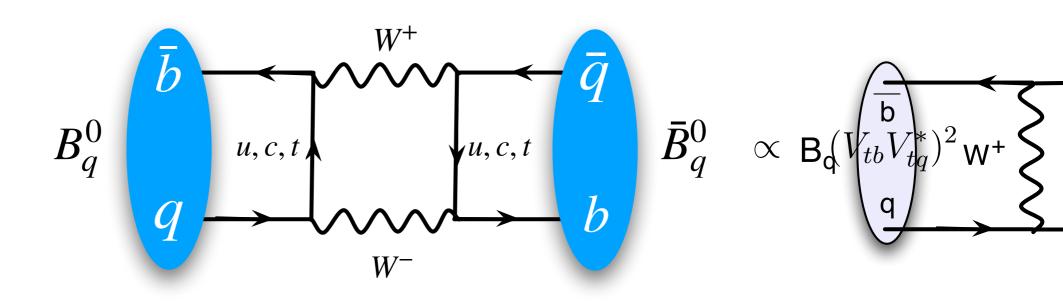
[GE with Miguel Escudero and Ann Nelson, Phys. Rev. D, arXiv:1810.00880] [GE with Gonzalo Alonso-Alvarez, Ann Nelson and Huangyu Xiao, JHEP, arXiv:1907.10612]

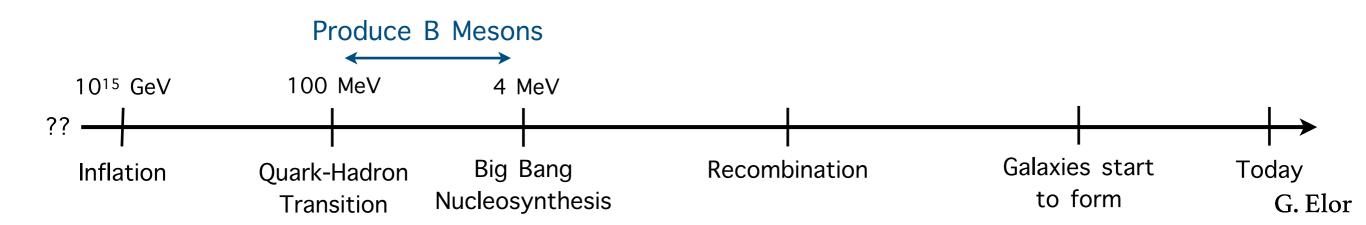
### Neutral B ]





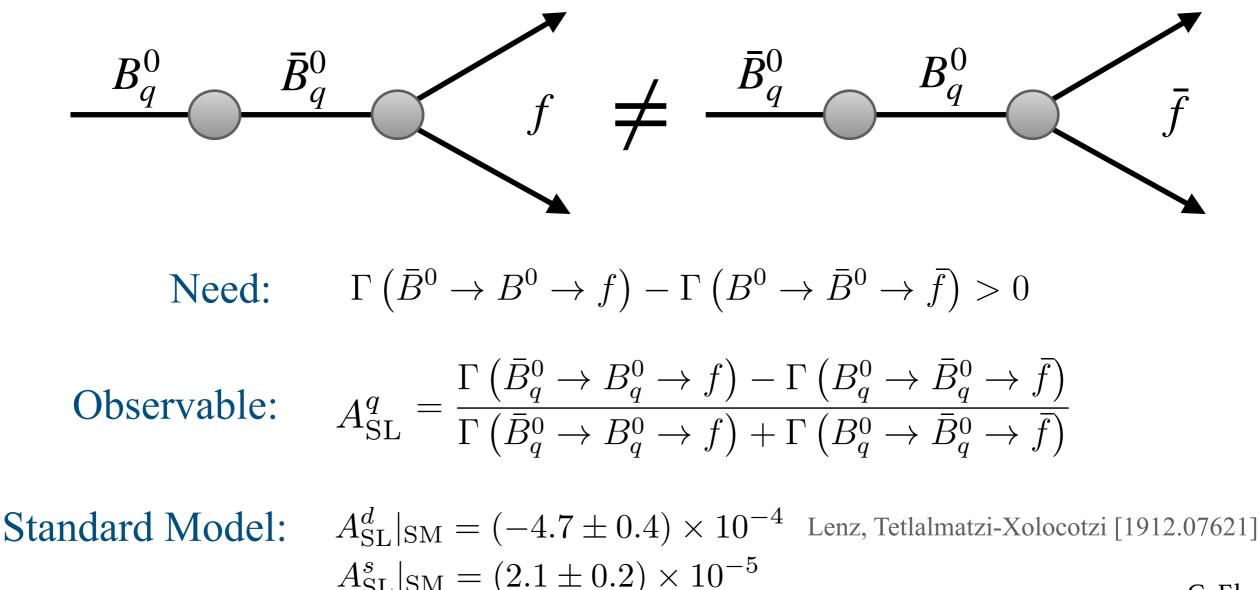
At low energies we can use CPV in *B* meson mixing e.g. from CKM phases in the case of the Standard Model (but new physics contributions are also not excluded)





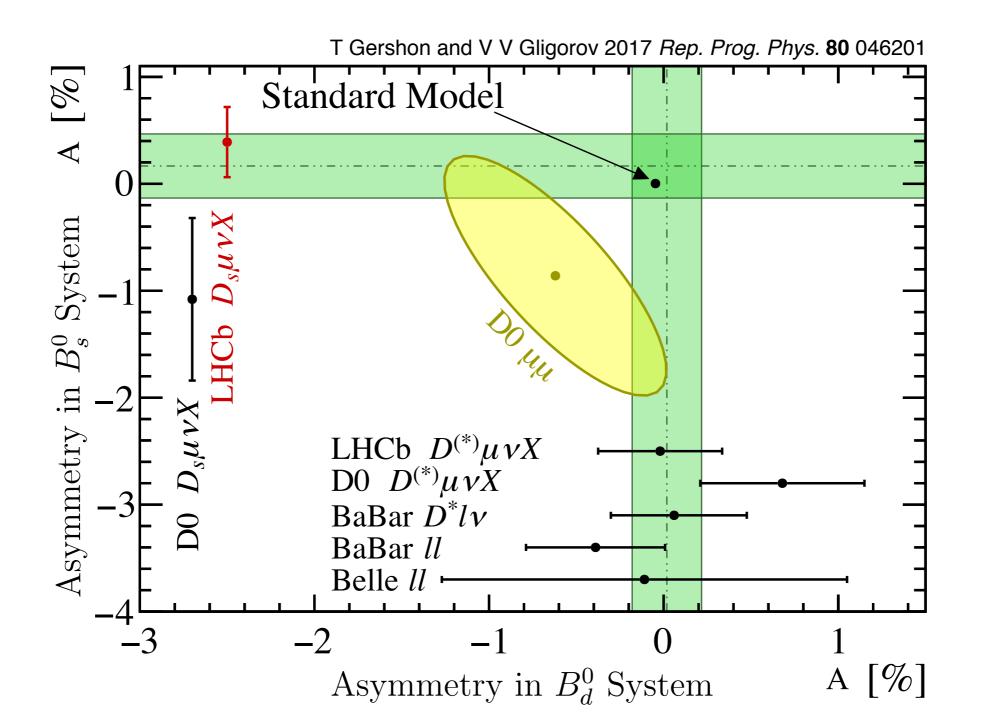
### **CP** Violation

B meson/anti-meson mixing has sizable CP violation



## Asymmetry in B Meson Mixing

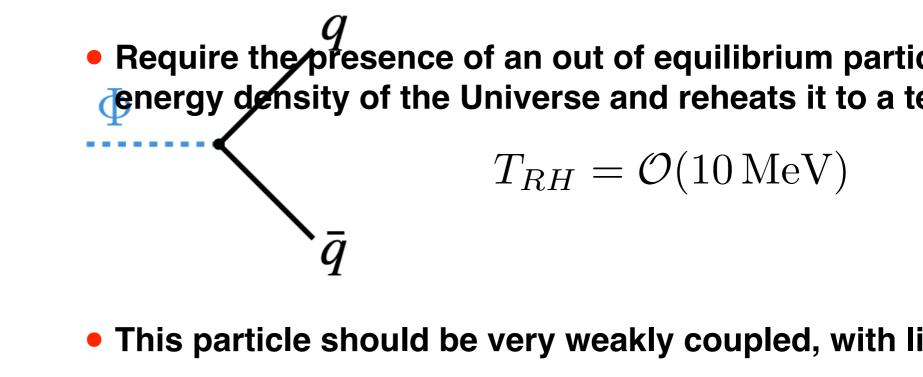
#### Can accommodate contributions from new physics



### Sakharov I. Out of Equilibrium

Late deczy of an "inflaton-like" field

Decays at:  $\Gamma_{\Phi} = 4H(T_R)$  to quarks  $m_{\Phi} \in [5 \text{ GeV}, 100 \text{ GeV}]$ 

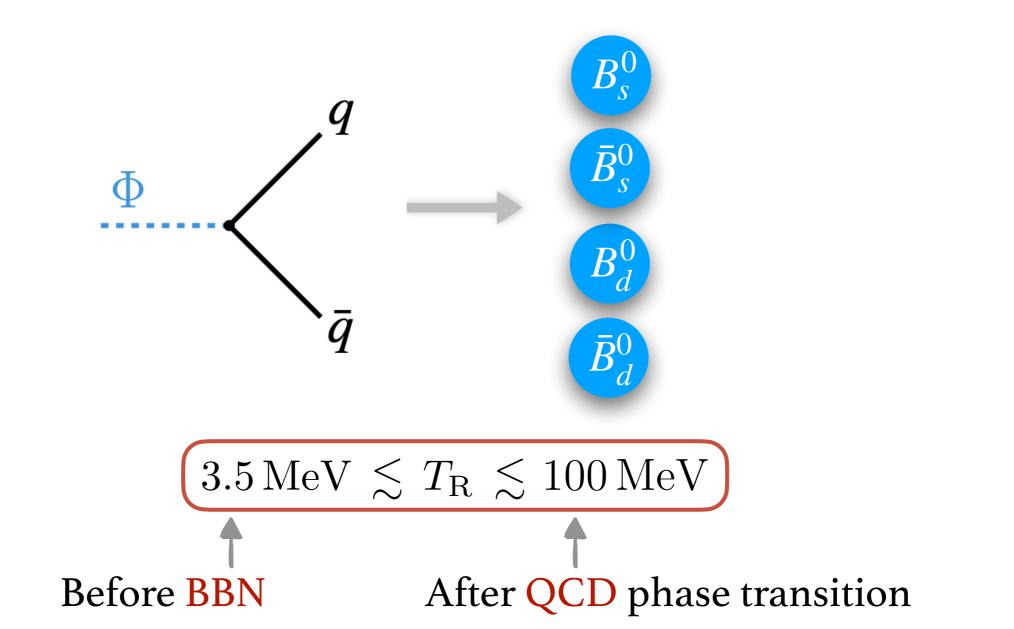


$$\begin{array}{ccc} 3.5\,\mathrm{MeV}\,\lesssim\,T_\mathrm{R}\,\lesssim\,100\,\mathrm{MeV} & \tau_\Phi=\mathcal{O}(10^{-3}\,\mathrm{s}) \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ &$$

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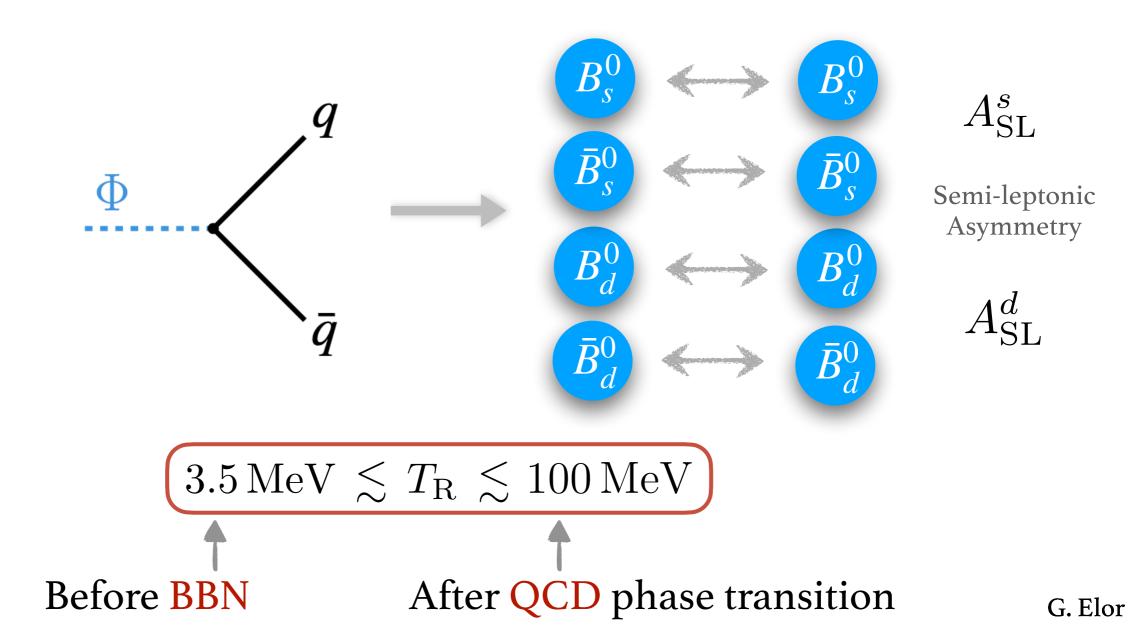
Decays at:  $\Gamma_{\Phi} = 4H(T_R)$  to quarks  $m_{\Phi} \in [5 \text{ GeV}, 100 \text{ GeV}]$ 



### Sakharov II. CP Violation

Late decay of an "inflaton-like" field

Decays at:  $\Gamma_{\Phi} = 4H(T_R)$  to quarks  $m_{\Phi} \in [5 \text{ GeV}, 100 \text{ GeV}]$ 



### Sakharov III. B Violation?

Need a way to change baryon number

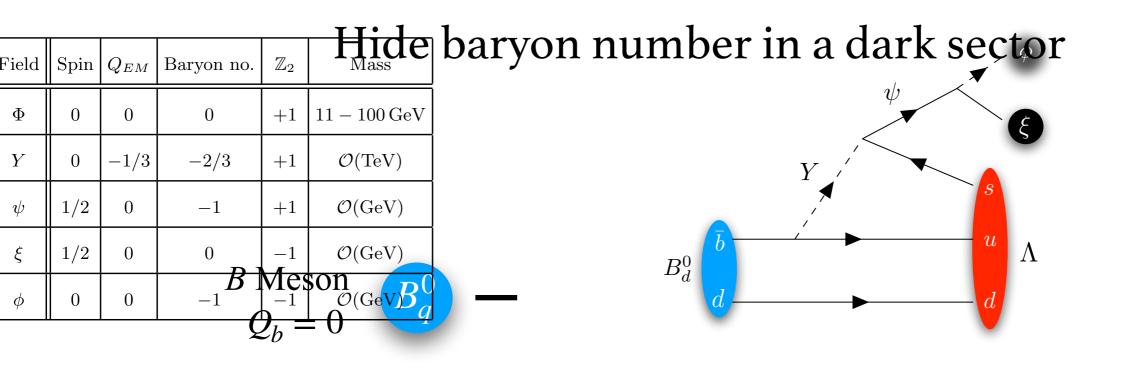


Hide baryon number in a dark sector rather than violate it



Will minimally need four new fields

### Dark Sector Baryon



Kinematics:  $m_{\psi} < m_B - m_{\text{Baryon}} < 4.3 \,\text{GeV}$ Proton stability:  $m_{\psi} > m_p - m_e \simeq 937.8 \,\text{MeV}$ 

Equal and opposite dark and visible baryon  $A^{\alpha}s\bar{y}n\bar{h}\bar{t}\bar{t}\bar{r}\bar{t}\bar{s}^{10}\bar{f}\bar{s}^{36}$  GeVed.  $Y_{\mathcal{B}} - Y_{\bar{\mathcal{B}}} = -(Y_{\psi} - Y_{\bar{\psi}})$ 

# Cchreat Appled Scale

	ſ	Field Sp	in $Q_{EM}$ Baryon	1 no. 2	$\mathbb{Z}_2$ Mass	
Field	Spin	$Q_{EM}$	Baryon no.	$\mathbb{Z}_2$	Mass	
Φ	0	0	0	+1	$11-100\mathrm{GeV}$	
Y	0	-1/3	-2/3	+1	$\mathcal{O}({ m TeV})$	
$\psi$	1/2	0	-1	+1	$\mathcal{O}({ m GeV})$	

Allowed by all the symmetries:

$$\mathcal{L} \supset -y_{ub} Y^* \, \bar{u} \, b^c - y_{\psi s} Y \bar{\psi}$$

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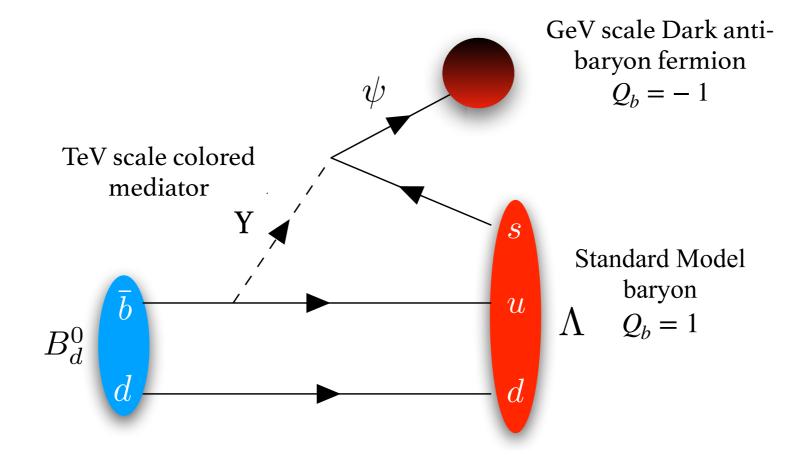
Colored triplet scalar must be TeV state to be consistent with dollider search so we will integrate it out:  $m_Y^{2b}y_{1/2}$  so  $m_Y^{2}$ 

$$\mathcal{H} = \bigcup_{Y}^{\kappa} b \, u \, s \, \psi$$

Note: this interaction does not change baryon nomber  $\frac{m_B}{2 \text{ GeV}}$ 

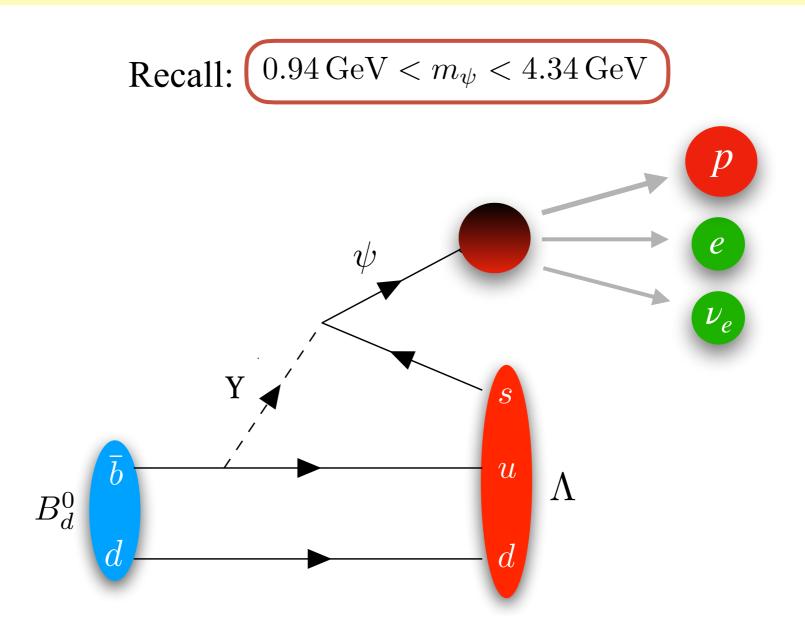
### New decay of the B Meson

For instance,  $\mathcal{O}_{us} = \psi b \, u \, s$  mediates the decay  $\bar{b} \to \psi + u + s$ 



Decay modes of a *B* meson to baryon and missing energy have historically not been searched for. Relatively unconstrained.

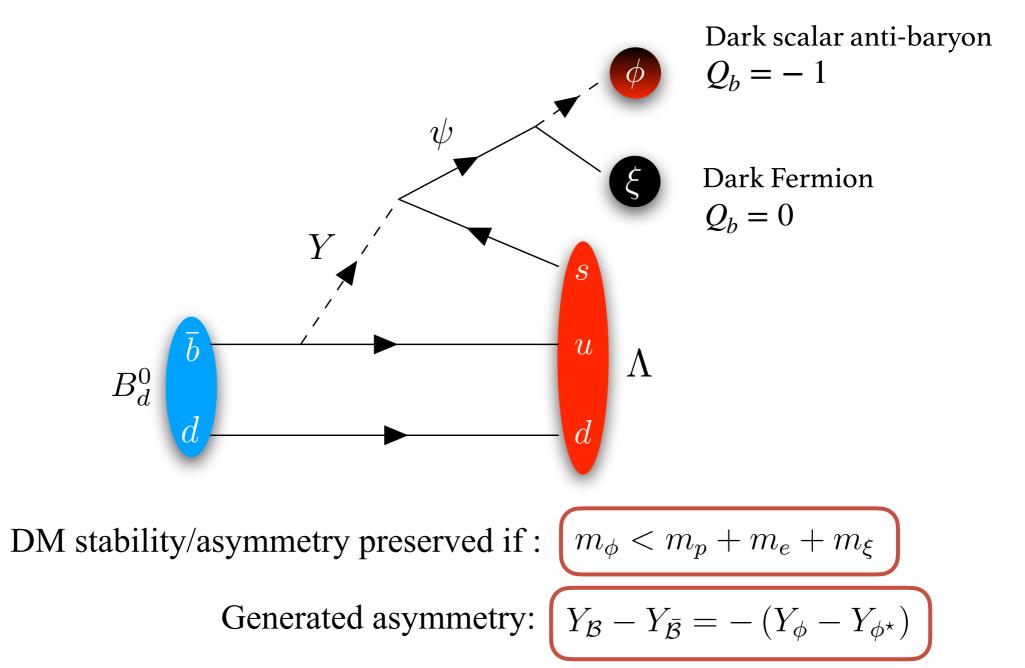
### Dark Matter



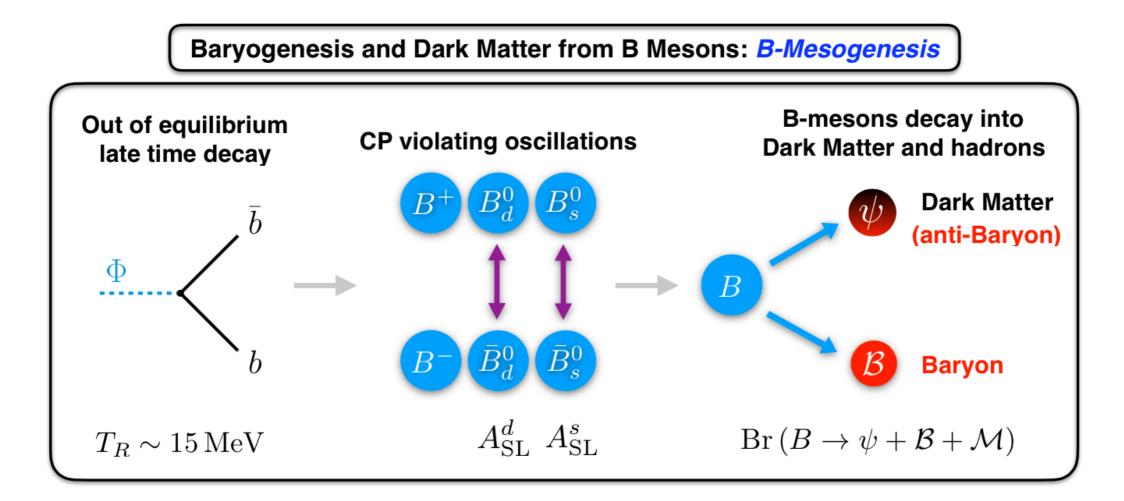
New dark baryon is unstable and will decay to baryonic matter, washing out the asymmetry in the process. It cannot be the dark matter.

### Dark Matter

Dark fermion must quickly decay within the dark sector into two new dark sector fields



### Summary: The Mechanism



#### Baryon asymmetry is related to experimental observables:

- Branching fraction for inclusive decay of B mesons into baryons, missing energy, and any number of final state mesons.
- Semi-leptonic asymmetries  $Y_B \propto \sum_{q=s,d} A_{SL}^q \times Br\left(B^0 \to \psi \mathcal{BM}\right)$  G. Elor

## Summary: Mirainal Mudal

#### *B* - Mesogenesis requires the minimal field content:

	Field	Spin	$Q_{EM}$	Baryon no.	$\mathbb{Z}_2$	Mass
Inflaton:	Φ	0	0	0	+1	$11 - 100 \mathrm{GeV}$
Colored Triplet Scalar:	Y	0	-1/3	-2/3	+1	$\mathcal{O}({ m TeV})$
Dark Dirac Baryon:	$\psi$	1/2	0	-1	+1	$\mathcal{O}({ m GeV})$
Dark Fermion:	ξ	1/2	0	0	-1	$\mathcal{O}({ m GeV})$
Dark Scalar Baryon:	$\phi$	0	0	-1	-1	$\mathcal{O}({ m GeV})$

## Sumary: Mainal Madel

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Aside: this can be embedded in e.g. a supersymmetric theory

 $Y \leftrightarrow \tilde{q}_R$  squark,  $\psi \leftrightarrow \begin{bmatrix} \tilde{B} \\ \lambda_s^{\dagger} \end{bmatrix}$  Dirac bino,  $\phi, \xi \leftrightarrow$  sterile neutrino multiplet

[GE with G. Alonso-Alvarez, A. E. Nelson, H. Xiao JHEP [arXiv:1907.10612]]

#### However, we do not need a UV model for numerics or general signals

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Colored Triplet Scalar:	Y	0	-1/3	-2/3	+1	$\mathcal{O}({ m TeV})$	Numerically study the evolution of the
Dark Dirac Baryon:	$\psi$	1/2	0	-1	+1	$\mathcal{O}({ m GeV})$	Boltzmann equations describing these fields
Dark Fermion:	ξ	1/2	0	0	-1	$\mathcal{O}({ m GeV})$	deserroting these fields
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#### However, we *do not need a UV model* for numerics or general signals

Late time decay of Inflaton

 $\Gamma_{\Phi} = 4H(T_R)$ 

• Inflaton:  $\frac{dn_{\Phi}}{dt} + 3Hn_{\Phi} = -\Gamma_{\Phi}n_{\Phi}$ • Radiation:  $\frac{d\rho_{\text{rad}}}{dt} + 4H\rho_{\text{rad}} = +\Gamma_{\Phi}m_{\Phi}n_{\Phi}$ • Hubble:  $H^{2} = \frac{8\pi}{3M_{\text{Pl}}^{2}}\left(\rho_{\text{rad}} + m_{\Phi}n_{\Phi}\right)$ 

#### Dark Matter

• Symmetric component of the dark scalar baryon

$$\frac{dn_{\phi+\phi^*}}{dt} + 3Hn_{\phi+\phi^*} = 2\Gamma_{\Phi}^B n_{\Phi} - 2\langle\sigma v\rangle_{\phi} \left(n_{\phi+\phi^*}^2 - n_{\text{eq},\phi+\phi^*}^2\right)$$

• The dark Majorana fermion

$$\frac{dn_{\xi}}{dt} + 3Hn_{\xi} = 2\Gamma_{\Phi}^{B}n_{\Phi} - \langle \sigma v \rangle_{\xi} \left( n_{\xi}^{2} - n_{\text{eq},\xi}^{2} \right)$$

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$$\Gamma_{\Phi}^{B} \equiv \Gamma_{\Phi} \times \operatorname{Br} \left( B \to \psi \mathcal{B} \mathcal{M} \right)$$

Simplification: For the (low) temperature range of interest we can check that the *B* mesons decay more quickly than they annihilate

#### Dark Matter

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$$\frac{dn_{\phi+\phi^*}}{dt} + 3Hn_{\phi+\phi^*} = 2\Gamma_{\Phi}^B n_{\Phi} - 2\langle\sigma v\rangle_{\phi} \left(n_{\phi+\phi^*}^2 - n_{\text{eq},\phi+\phi^*}^2\right)$$

• The dark Majorana fermion

Overproduced particle must be depleted by additional dark interactions.

$$\frac{dn_{\xi}}{dt} + 3Hn_{\xi} = 2\Gamma_{\Phi}^{B}n_{\Phi} - \langle \sigma v \rangle_{\xi} \left(n_{\xi}^{2} - n_{eq,\xi}^{2}\right)$$
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Simplification: For the (low) temperature range of interest we can check that the *B* mesons decay more quickly than they annihilate

### Numerics

#### The Baryon Asymmetry

• Anti-symmetric dark sector baryon makes up the baryon asymmetry

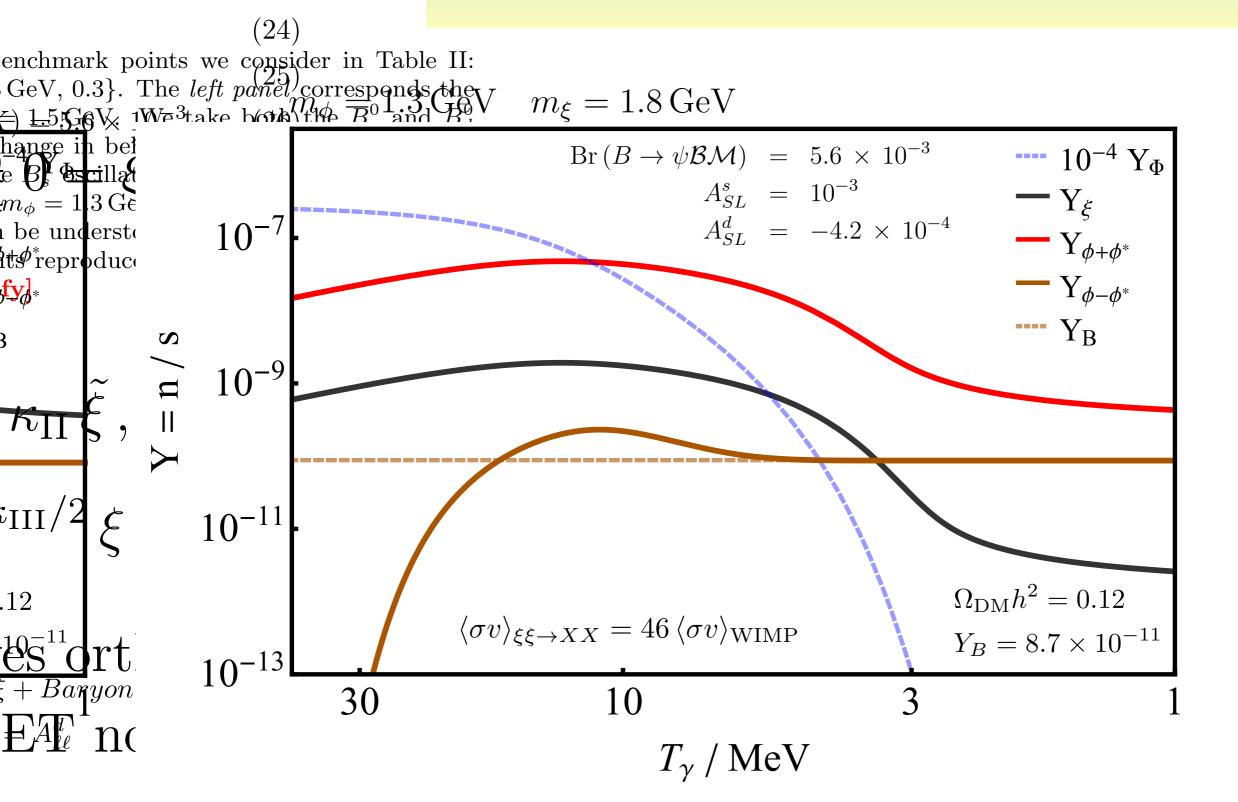
eV

$$\underbrace{\frac{\partial n_{\phi-\phi^*}}{\partial t} + 3Hn_{\phi-\phi^*}}_{b} = 2\Gamma^B_{\Phi}\sum_q \operatorname{Br}\left(\bar{b} \to B^0_q\right)A^q_{\mathrm{SL}}f^q_{\mathrm{deco}}n_{\Phi}$$

 $200 \,\mathrm{MeV}$ Coherent *B* meson oscillations maintained for 20 MeV scales and below

 $e^{\pm} \qquad \Gamma\left(e^{\pm}B_{q}^{0} \to e^{\pm}B_{q}^{0}\right) = 10^{-11} \text{GeV}\left(\frac{T}{20 \text{MeV}}\right)^{5}$   $f_{\text{deco}}^{q} = e^{-\Gamma\left(e^{\pm}B_{q}^{0} \to e^{\pm}B_{q}^{0}\right)/\Delta m_{B_{q}}}$   $B^{0} \qquad T_{B_{s}} \leq 20 \text{ MeV and } T_{B_{d}} \leq 10 \text{ MeV}$ 

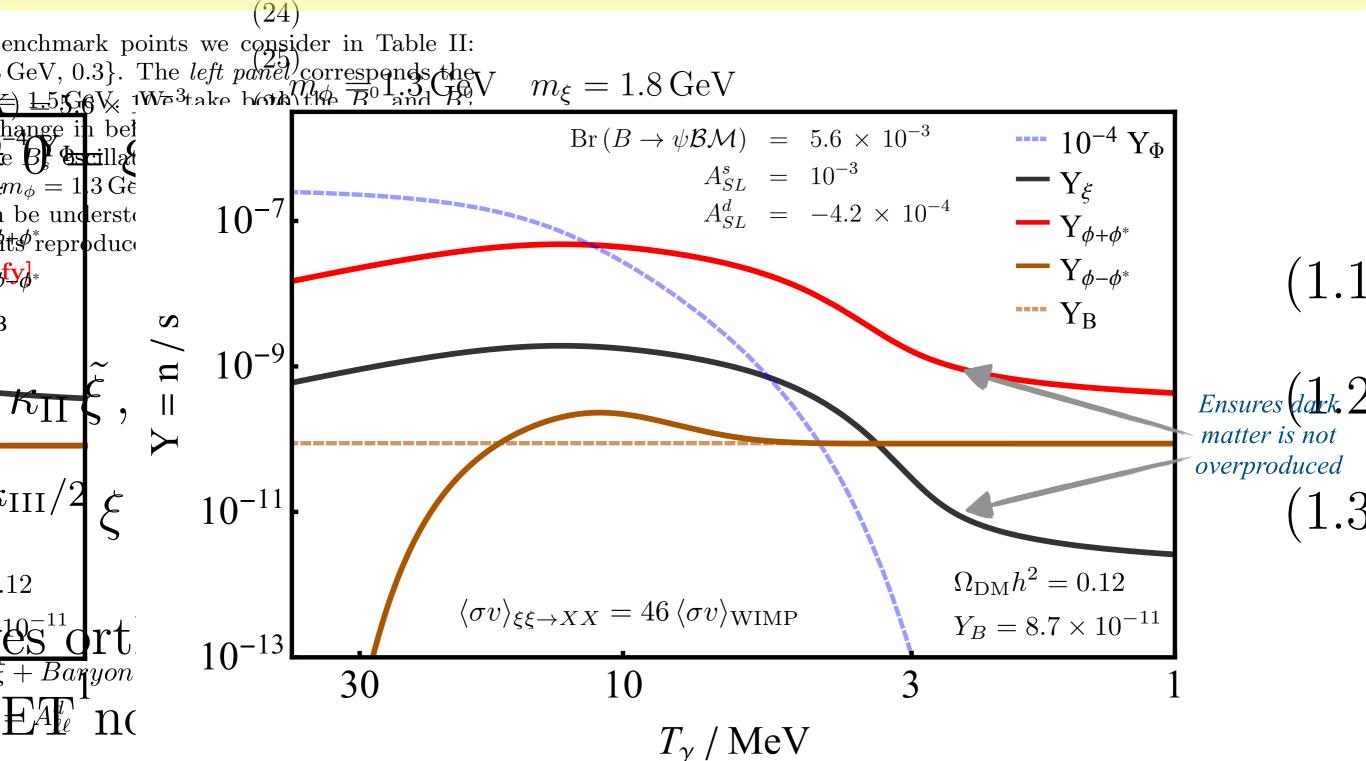
#### (20) **Example Benchmark Point** (23)



(1.1)(1.2)(1.3)

lS

### (20) Example Benchmark Point

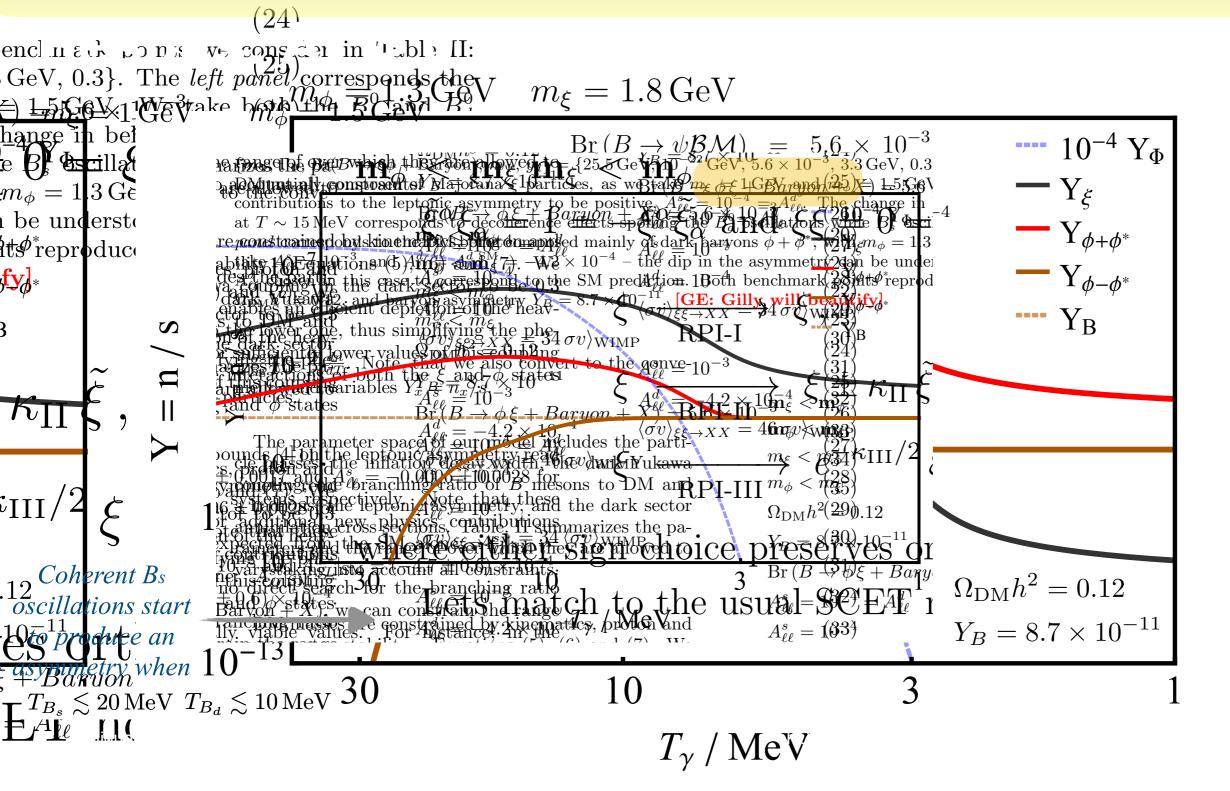


 $\langle \mathbf{a} \mathbf{a} \rangle$ 

**1S** 

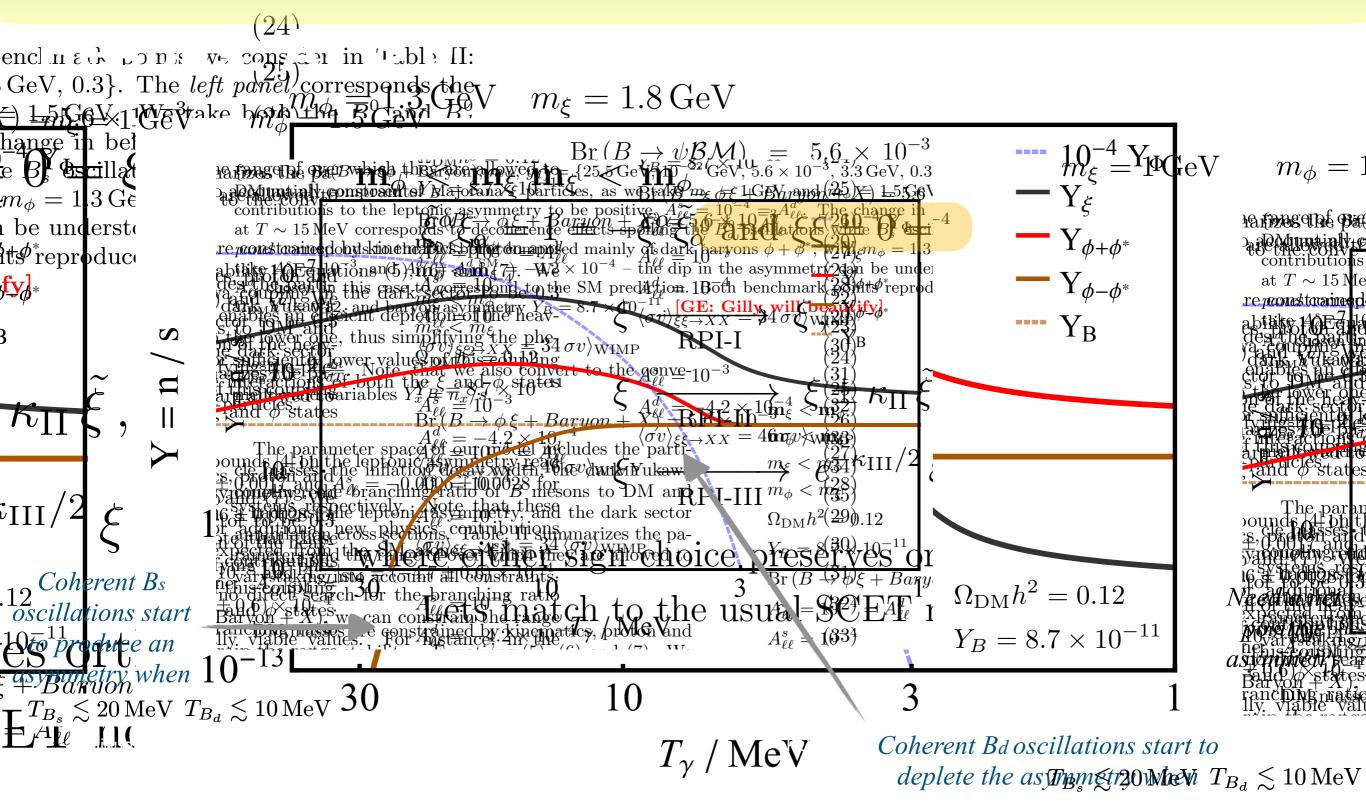
### (20) Example Benchmark Point

**1S** 



(1.1)(1.2)(1.3)

#### (20) Example Benchmark Point



 $\mathbf{1S}$ 

#### Numerical Result

$$Y_{\mathcal{B}} \simeq 8.7 \times 10^{-11} \frac{\operatorname{Br} \left( B \to \psi \mathcal{B} \mathcal{M} \right)}{10^{-2}} \sum_{q=s,d} \alpha_q \frac{A_{\mathrm{SL}}^q}{10^{-4}}$$

Successful *B*-Mesogenesis

$$A_{\mathrm{SL}}^{s,d} \times \mathrm{Br}\left(B^0 \to \psi \,\mathcal{B} \,\mathcal{M}\right) > 10^{-6}$$

**Experimental Observables** 

# Part II. A Roadmap for Discovering *B*-Mesogenesis



Based on:

GE with Gonzalo Alonso-Alvarez, Miguel Escudero, [arXiv:2101.02706]

Ongoing theoretical work: GE with Gonzalo Alonso-Alvarez, Jorge Martin Camalich, Miguel Escudero, Bartosz Fornal, Benjamin Grinstein (also see white paper coming out in August) Ongoing searches at Belle 1 and II

Possible searches at LHCb - see arXiv:2105.12668 and arXiv:2106.12870

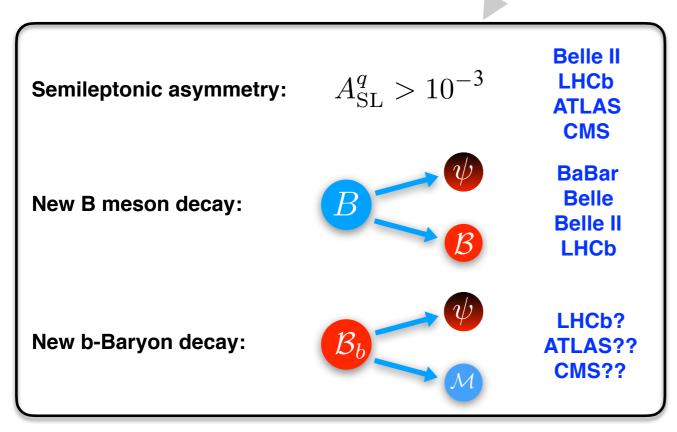
G. Elor

# Signals of B-Mesogenesis

For successful baryogenesis:  $A_{\rm SL}^{s,d} \times Br(B^0 \to \psi \mathcal{BM}) > 10^{-6}$ 

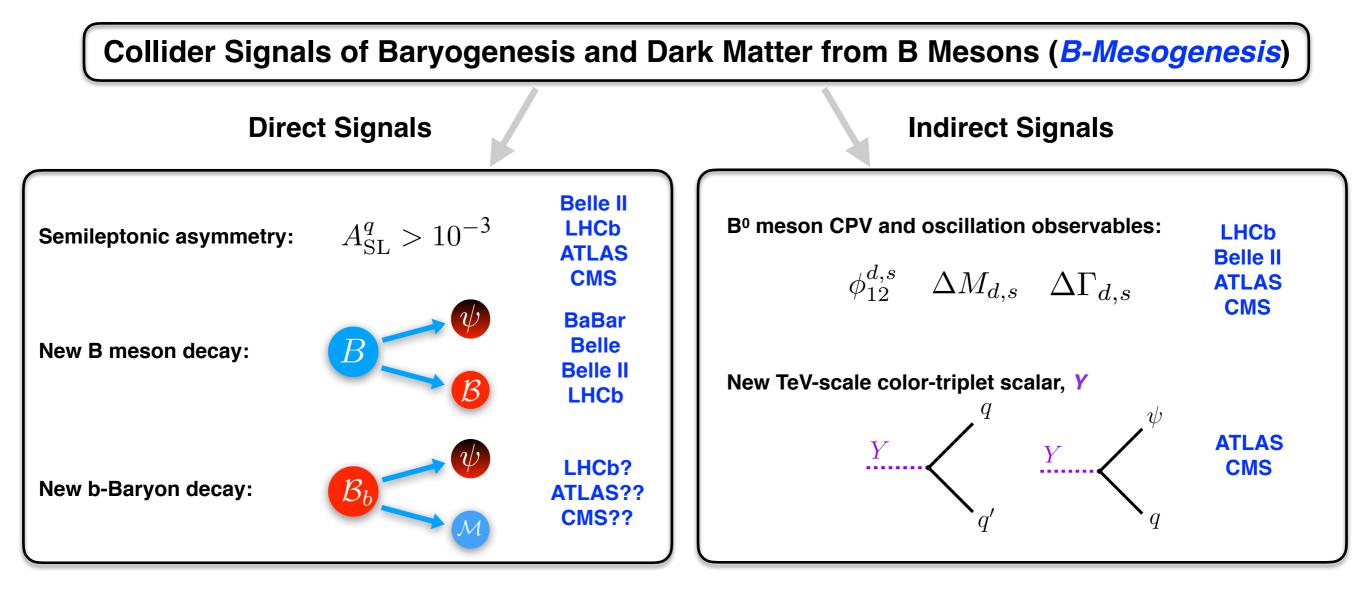
Collider Signals of Baryogenesis and Dark Matter from B Mesons (*B-Mesogenesis*)

**Direct Signals** 



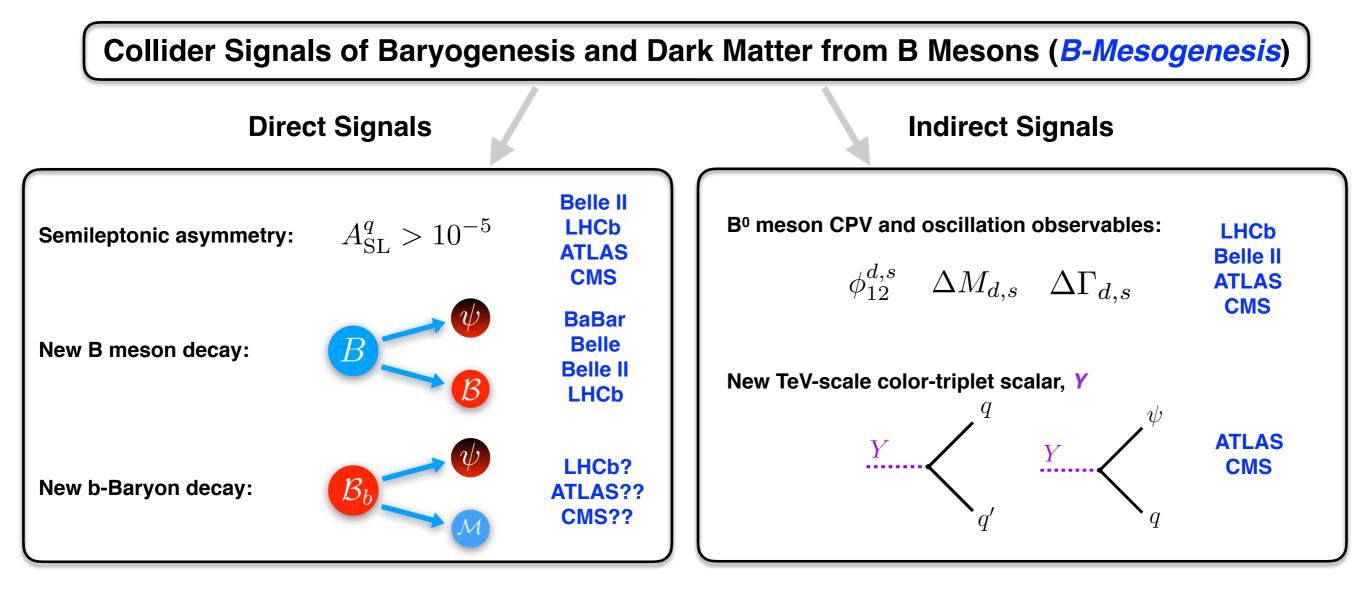
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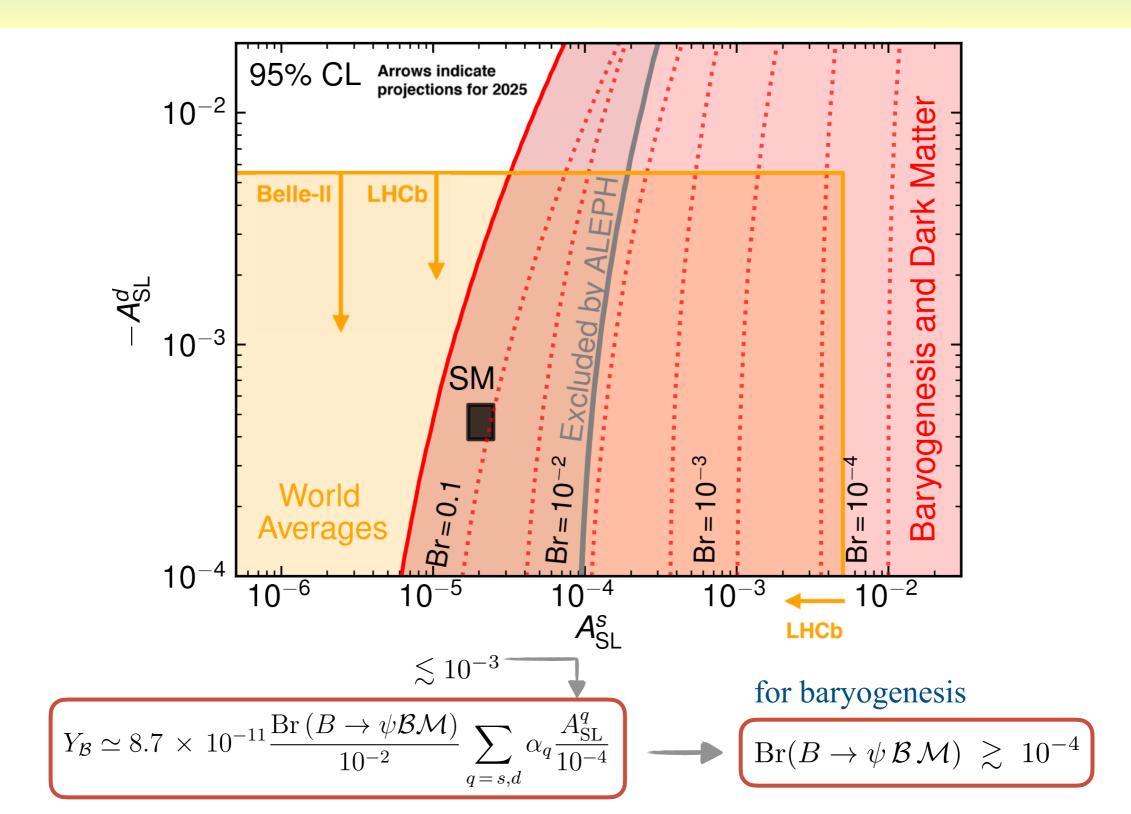
# Signals of B-Mesogenesis

For successful baryogenesis:  $A_{\rm SL}^{s,d} \times {\rm Br} \left( B^0 \to \psi \, \mathcal{B} \, \mathcal{M} \right) > 10^{-6}$ 



Independent of UV model. Given a UV model there will be even more signals!

### The Semi-Leptonic Asymmetry



G. Elor

## Flavorful Variations

No a priori reason to expect a particular flavor structure.

Most general interactions:

$$\mathcal{L}_{-1/3} = -\sum_{i,j} y_{u_i d_j} Y^* \bar{u}_{iR} d_{jR}^c - \sum_k y_{\psi d_k} Y d_{kR}^c \bar{\psi} + \text{h.c.}$$

Possible operators:

$$egin{aligned} \mathcal{O}_{ud} &= \psi \, b \, u \, d \ \mathcal{O}_{us} &= \psi \, b \, u \, s \ \mathcal{O}_{cd} &= \psi \, b \, c \, d \ \mathcal{O}_{cs} &= \psi \, b \, c \, s \end{aligned}$$

*B*-Mesogenesis requires:

$$\operatorname{Br}(B \to \psi \,\mathcal{B} \,\mathcal{M}) \gtrsim 10^{-4}$$

#### Can be searched for at Belle, BaBar and LHCb

Operator/Decay	Initial State	Final state	
$\mathcal{O} = \psi  b  u  d$ $\bar{b} \to \psi  u  d$	$B_d$ $B_s$ $B^+$ $\Lambda_b$	$\begin{split} \psi + n  (udd) \\ \psi + \Lambda  (uds) \\ \psi + p  (duu) \\ \bar{\psi} + \pi^0 \end{split}$	
$\mathcal{O} = \psi  b  u  s$ $\bar{b} \to \psi  u  s$	$B_d$ $B_s$ $B^+$ $\Lambda_b$	$\psi + \Lambda (usd)$ $\psi + \Xi^{0} (uss)$ $\psi + \Sigma^{+} (uus)$ $\overline{\psi} + K^{0}$	
$\mathcal{O} = \psi  b  c  d$ $\overline{b} \to \psi  c  d$	$B_d$ $B_s$ $B^+$ $\Lambda_b$	$\psi + \Lambda_c + \pi^- (cdd)$ $\psi + \Xi_c^0 (cds)$ $\psi + \Lambda_c (dcu)$ $\overline{\psi} + \overline{D}^0$	
$\mathcal{O} = \psi  b  c  s$ $\bar{b} \to \psi  c  s$	$B_d$ $B_s$ $B^+$ $\Lambda_b$	$\psi + \Xi_c^0 (csd)$ $\psi + \Omega_c (css)$ $\psi + \Xi_c^+ (csu)$ $\bar{\psi} + D^- + K^+$	

#### Can be searched for at Belle, BaBar and LHCb

Operator/Decay	Initial State	Final state	
$\mathcal{O} = \psi  b  u  d$ $\bar{b} \to \psi  u  d$	$ \begin{array}{c} B_d \\ B_s \\ B^+ \\ \Lambda_b \end{array} $	$egin{aligned} \psi + n  (udd) \ \psi + \Lambda  (uds) \ \psi + p  (duu) \ ar{\psi} + \pi^0 \end{aligned}$	 Directly related to baryon asymmetry Most stringent constraints actually comes from a 20 year
$\begin{aligned} \mathcal{O} &= \psi  b  u  s \\ \bar{b} &\to \psi  u  s \end{aligned}$	$ \begin{array}{c} B_d \\ B_s \\ B^+ \\ \overline{\Lambda_b} \end{array} $	$\begin{split} \psi + \Lambda \left( usd \right) \\ \psi + \Xi^0 \left( uss \right) \\ \psi + \Sigma^+ \left( uus \right) \\ \bar{\psi} + K^0 \end{split}$	old search at LEP [hep-ex/0010022]
$\mathcal{O} = \psi  b  c  d$ $\bar{b} \to \psi  c  d$	$ \begin{array}{c} B_d\\ B_s\\ B^+\\ \Lambda_b \end{array} $	$\psi + \Lambda_c + \pi^- (cdd)$ $\psi + \Xi_c^0 (cds)$ $\psi + \Lambda_c (dcu)$ $\bar{\psi} + \overline{D}^0$	
$\mathcal{O} = \psi  b  c  s$ $\bar{b} \to \psi  c  s$	$ \begin{array}{c} B_d\\ B_s\\ B^+\\ \Lambda_b \end{array} $	$\psi + \Xi_c^0 (csd)$ $\psi + \Omega_c (css)$ $\psi + \Xi_c^+ (csu)$ $\bar{\psi} + D^- + K^+$	G. E

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Operator/Decay	Initial State	Final state	
$\mathcal{O} = \psi  b  u  d$ $\bar{b} \to \psi  u  d$	$B_d$ $B_s$ $B^+$ $\Lambda_b$	$\begin{split} \psi + n  (udd) \\ \psi + \Lambda  (uds) \\ \psi + p  (duu) \\ \bar{\psi} + \pi^0 \end{split}$	Directly related to baryon asymmetry Most stringent constraints actually comes from a 20 year
$\mathcal{O} = \psi  b  u  s$ $\bar{b} \to \psi  u  s$	$ \begin{array}{c} B_d \\ B_s \\ B^+ \\ \hline \Lambda_b \end{array} $	$\begin{split} \psi + \Lambda \left( usd \right) \\ \psi + \Xi^0 \left( uss \right) \\ \psi + \Sigma^+ \left( uus \right) \\ \bar{\psi} + K^0 \end{split}$	<ul> <li>old search at LEP         <ul> <li>[hep-ex/0010022]</li> <li>Indirectly constrains B-Mesogenesis.</li> <li>Charged track is an advantage for searches</li> </ul> </li> </ul>
$\mathcal{O} = \psi  b  c  d$ $\bar{b} \to \psi  c  d$	$ \begin{array}{c} B_d \\ B_s \\ B^+ \\ \Lambda_b \end{array} $	$\psi + \Lambda_c + \pi^- (cdd)$ $\psi + \Xi_c^0 (cds)$ $\psi + \Lambda_c (dcu)$ $\overline{\psi} + \overline{D}^0$	$B^{+} \underbrace{u}_{V} \underbrace{u}$
$\mathcal{O} = \psi  b  c  s$ $\bar{b} \to \psi  c  s$	$B_d$ $B_s$ $B^+$ $\Lambda_b$	$\psi + \Xi_c^0 (csd)$ $\psi + \Omega_c (css)$ $\psi + \Xi_c^+ (csu)$ $\bar{\psi} + D^- + K^+$	$B^{+} \underbrace{\begin{matrix} u \\ b \\ y \\ y \\ \psi \end{matrix}} \Lambda_{c} \qquad B^{+} \underbrace{\begin{matrix} u \\ b \\ y \\ y \\ \psi \end{matrix}} \Xi_{c}^{+}$ G. Elor

#### Can be searched for at Belle, BaBar and LHCb

Operator/Decay	Initial State	Final state	
$\mathcal{O} = \psi  b  u  d$ $\overline{b} \to \psi  u  d$	$ \begin{array}{c} B_d \\ B_s \\ B^+ \\ \Lambda_b \end{array} $	$\begin{split} \psi + n  (udd) \\ \psi + \Lambda  (uds) \\ \psi + p  (duu) \\ \bar{\psi} + \pi^0 \end{split}$	Directly related to baryon asymmetry Most stringent constraints actually comes from a 20 year
$\mathcal{O} = \psi  b  u  s$ $\bar{b} \to \psi  u  s$	$ \begin{array}{c} B_d \\ B_s \\ B^+ \\ \hline \Lambda_b \end{array} $	$\begin{split} \psi + \Lambda \left( usd \right) \\ \psi + \Xi^0 \left( uss \right) \\ \psi + \Sigma^+ \left( uus \right) \\ \bar{\psi} + K^0 \end{split}$	<ul> <li>old search at LEP         <ul> <li>[hep-ex/0010022]</li> <li>Indirectly constrains B-Mesogenesis.</li> <li>Charged track is an advantage for searches</li> </ul> </li> </ul>
$\mathcal{O} = \psi  b  c  d$ $\bar{b} \to \psi  c  d$	$ \begin{array}{c} B_d\\ B_s\\ B^+\\ \Lambda_b \end{array} $	$\psi + \Lambda_c + \pi^- (cdd)$ $\psi + \Xi_c^0 (cds)$ $\psi + \Lambda_c (dcu)$ $\overline{\psi} + \overline{D}^0$	b-flavored baryon decays can
$\mathcal{O} = \psi  b  c  s$ $\bar{b} \to \psi  c  s$	$ \begin{array}{c} B_d \\ B_s \\ B^+ \\ \Lambda_b \end{array} $	$\psi + \Xi_c^0 (csd)$ $\psi + \Omega_c (css)$ $\psi + \Xi_c^+ (csu)$ $\bar{\psi} + D^- + K^+$	yield indirect constraints. expect: $Br(\mathcal{B}_b \to \bar{\psi} \mathcal{M}) > 10^{-4}$ G. Elo

#### Ongoing Belle-I and II Search

Flavorful variations:

Operator/Decay	Initial State	Final state	
$\mathcal{O} = \psi  b  u  d$	$B_d$ $B_s$	$\psi + n  (udd) \ \psi + \Lambda  (uds)$	
$\overline{b} \to \psi  u  d$	$B^+ \ \Lambda_b$	$\psi + p \left( duu  ight) \ ar{\psi} + \pi^0$	
	$B_d$	$\psi + \Lambda \left( usd  ight)$	
$\mathcal{O} = \psi  b  u  s$	$B_s$	$\psi + \Xi^0 (uss)$	
$b \rightarrow \psi  u  s$	$B^+$	$\psi + \Sigma^+ (uus)$	
	$\Lambda_b$	$\bar{\psi} + K^0$	
	$B_d$	$\psi + \Lambda_c + \pi^- (cdd)$	
$\mathcal{O} = \psi  b  c  d$ $\bar{b} \to \psi  c  d$	$B_s$	$\psi + \Xi_c^0  (cds)$	
	$B^+$	$\psi + \Lambda_c (dcu)$	
	$\Lambda_b$	$ar{\psi} + \overline{D}^0$	
	$B_d$	$\psi + \Xi_c^0 \ (csd)$	
$\mathcal{O} = \psi  b  c  s$	$B_s$	$\psi + \Omega_c \left( css \right)$	
$b \rightarrow \psi  c  s$	$B^+$	$\psi + \Xi_c^+ (csu)$	
	$\Lambda_b$	$\bar{\psi} + D^- + K^+$	

Results to be made public at the next
 collaboration meeting. They plan to look for other decay modes as well

#### Possibilities at LHCb

[See our white paper on "Stealth Physics at LHCb" 2105.12668]

- No handle on initial energy of decaying *B* meson so measuring missing energy is non-trivial.
- But, LHCb has advantages: larger number of *B* mesons produced than at Belle, excellent vertex resolution, and good particle reconstruction efficiencies.
- Some possibilities for searches do exist. e.g. new paper just last week!

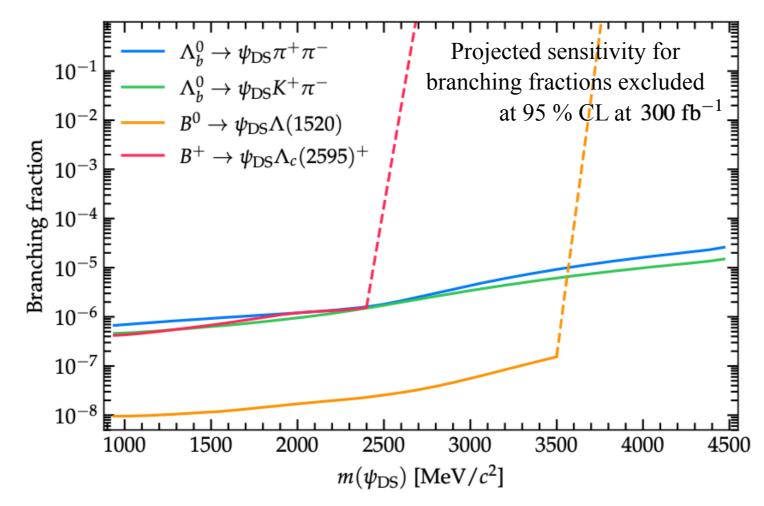
**Prospects on searches for baryonic Dark Matter produced in** *b***-hadron decays at LHCb** [2106.12870]

Alexandre Brea Rodríguez<sup>a,1</sup>, Veronika Chobanova<sup>b,1</sup>, Xabier Cid Vidal<sup>c,1</sup>, Saúl López Soliño<sup>d,1</sup>, Diego Martínez Santos<sup>e,1</sup>, Titus Mombächer<sup>f,1</sup>, Claire Prouvé<sup>g,1</sup>, Emilio Xosé Rodríguez Fernández<sup>h,1</sup>, Carlos Vázquez Sierra<sup>i,2</sup>

<sup>1</sup>Instituto Galego de Física de Altas Enerxías (IGFAE), Universidade de Santiago de Compostela, 15782 Santiago de Compostela, Spain <sup>2</sup>European Organization for Nuclear Research (CERN), Geneva, Switzerland

Proposed Search at LHCb [2106.12870]

- Search for decays of *B* mesons and *b*-Flavored baryons into an excited baryon in the final state  $B \to \psi B^*$
- The excited baryon promptly decay at the same decay point as original decay, allowing one to trigger on this decay.



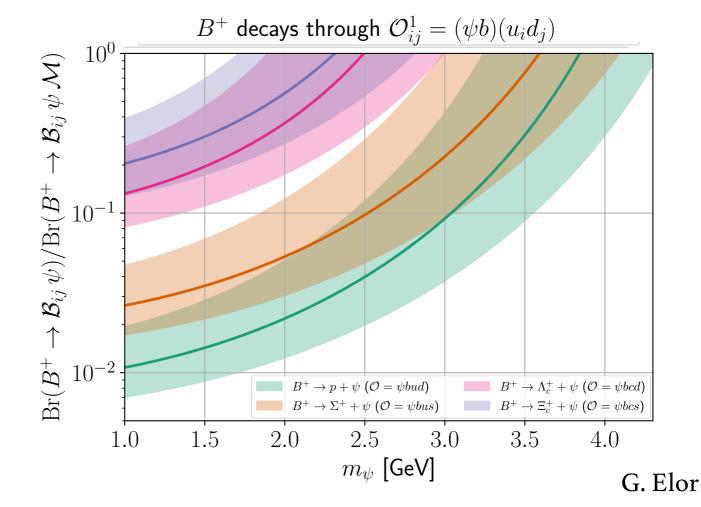
#### Caution: Inclusive vs. Exclusive Rates

• All decays (and their searches) discussed thus far have been *exclusive*. But, the observable controlling the baryon asymmetry is an *inclusive* rate.

$$Br(B \to \psi \mathcal{BM}) \gtrsim 10^{-4}$$

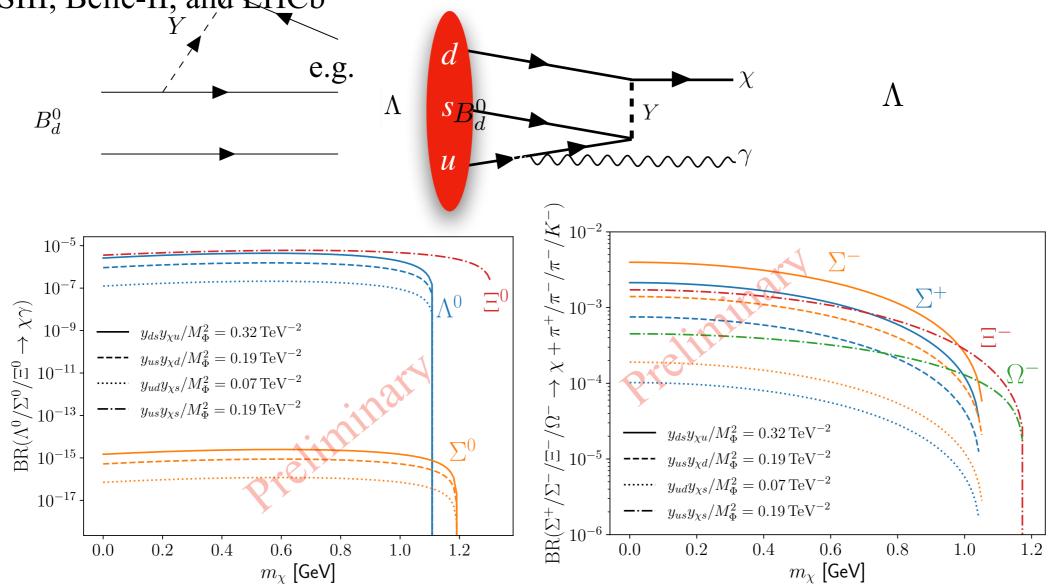
- Need a dedicated calculation using QCD sum rules or lattice techniques etc. to calculate form factors. Beyond my current expertise....
- Phase space method [Bigi, Phys.Lett.B 106, 510 (1981)]

$$\frac{\operatorname{Br}(B \to \psi \mathcal{B})}{\operatorname{Br}(B \to \psi \mathcal{B} \mathcal{M})} \gtrsim (1 - 10) \%.$$



## New Hyperon Decays

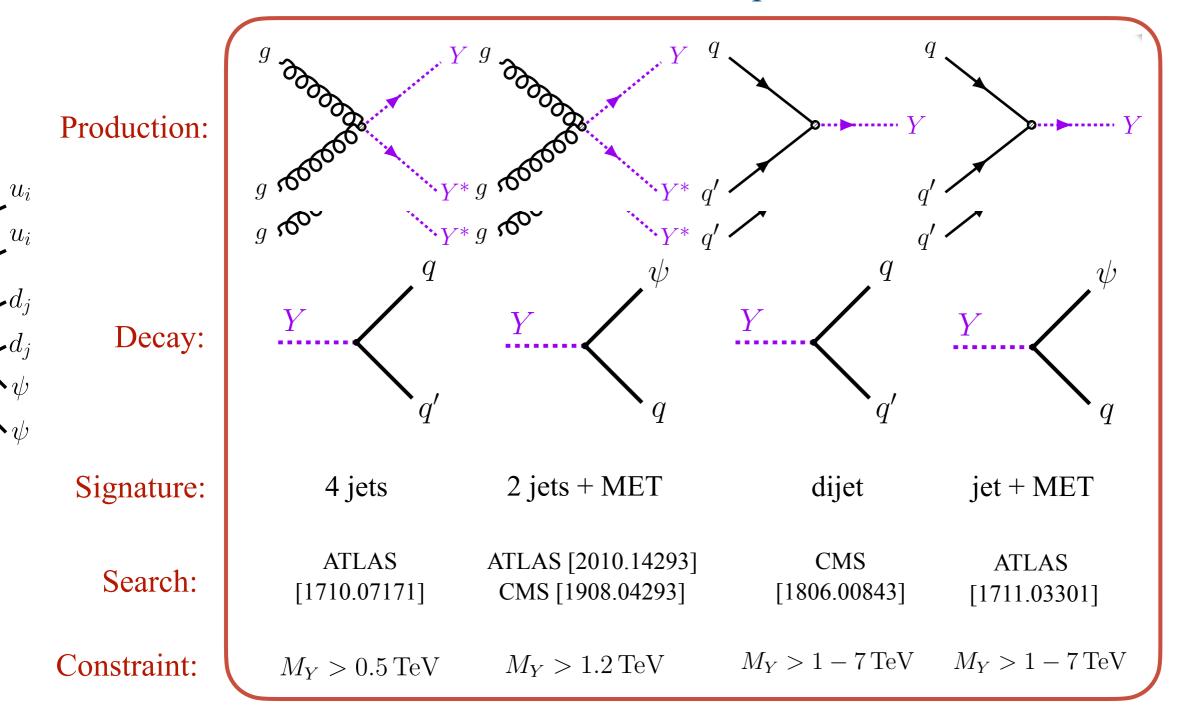
Light hadrons: we can compute form factors by matching onto chiral EFT. Hyperon decays are another indirect probe of B-Mesogenesis and are of interest at BESIII, Belle-II, and LHCb



[GE with Gonzalo Alonso-Alvarez, Jorge Martin Camalich, Miguel Escudero, Bartosz Fornal, Benjamin Grinstein]

# **Colored Triplet Scalar**

Constraints from LHC squark searches

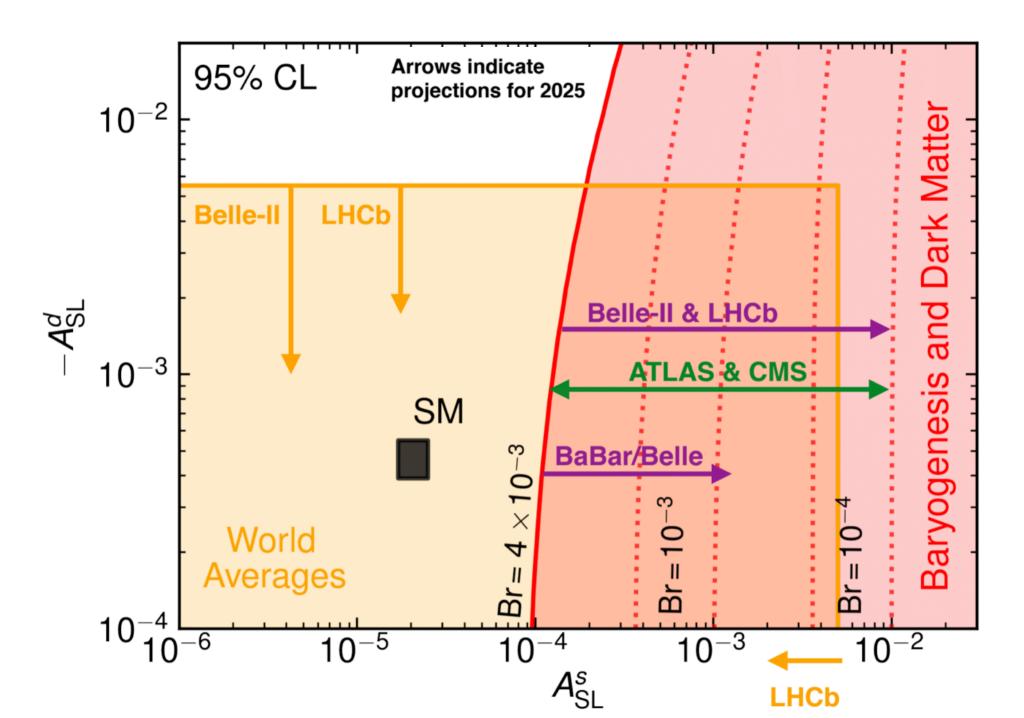


G. Elor

$$\frac{\overline{b} \rightarrow \psi u d}{Coolered Triplet Scalar}$$
Constraints from LHC squark searches
$$B-Mcsogenesis requires: Br(B \rightarrow \psi B M) \simeq 10^{-3} \left(\frac{\Delta m}{3 \text{ GeV}}\right)^4 \left(\frac{1.5 \text{ TeV}}{M_Y} \frac{\sqrt{y}u dy dy d}{0.53}\right)^4 \gtrsim 10^{-4}$$
Br > 10<sup>-4</sup>  $\Delta m = m_B - m_{\psi} - m_B - m_M$ 
Since collider bounds depend on the ratio  $\frac{\sqrt{y}u_i d_j y d_{yd_k}}{M_Y}$  they will in turn constrain the branching fraction.
$$c.g. \int_{0}^{10^{-1}} \frac{\sqrt{y}u_i d_j y d_{yd_k}}{10^{-1}} \int_{0}^{10^{-2}} \frac{\sqrt{y}u_i d_j y d_{yd_k}}{10^{-1}} \int_{0}^{10^{-2}} \frac{\sqrt{y}u_i d_j y d_{yd_k}}{1.25 - 1.50 - 1.75} \int_{0}^{2.00} \frac{2.25 - 2.50}{2.50}$$

#### Discovering B-Mesogenesis Outlook

Could be fully tested in but a few years.



## Discovering B-Mesogenesis Outlook

#### Predictions and Signals of B-Mesogenesis

• New *B* meson decay modes with

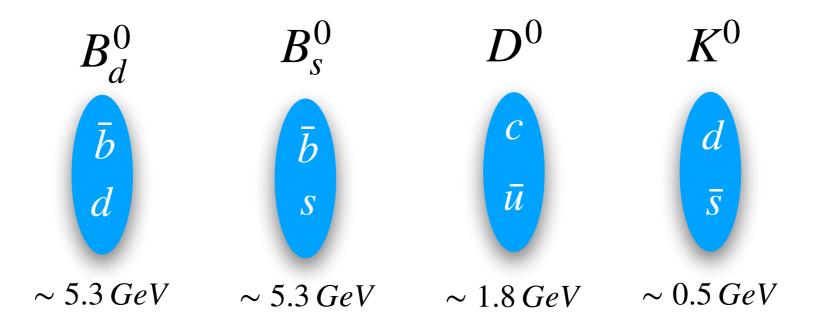
$$\operatorname{Br}(B \to \psi \,\mathcal{B} \,\mathcal{M}) \gtrsim 10^{-4}$$

• Positive semileptonic asymmetry

$$A_{\rm SL}^q > 10^{-4}$$

- The existence TeV scale colored triplet scalar
- Implications for flavor structure.
- Many more signals possible given a UV model see e.g. [1907.10612]

# Why B Mesons?



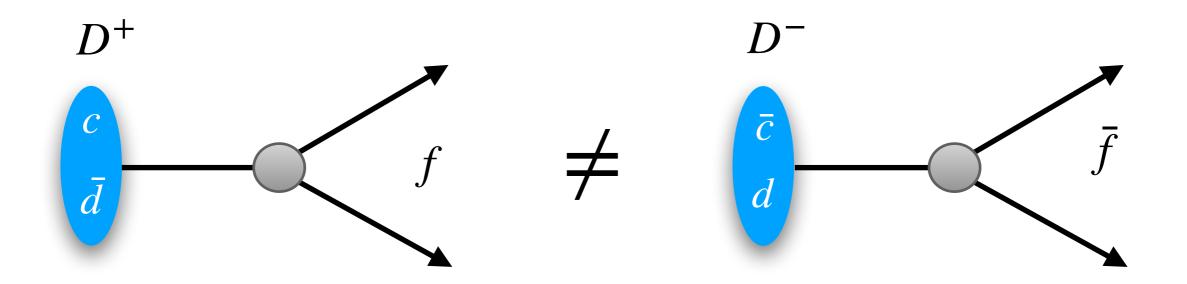
- Kinematics: Dark baryons must be GeV scale. Only *B* mesons are heavy enough to decay into GeV scale. Charge dark particle under lepton number instead, then it can be light.
- Neutral *D* Mesons don't have a lot of CP violation in their oscillations, but charged *D* Mesons have a lot of CP violation in their decays.

# Part III. D-Mesogenesis

Based on:

GE with Robert McGehee, Phys. Rev. D [arXiv:2011.06115]

## CPV in Charged D Decays



**Observable:** 
$$A_{CP}^f = \frac{\Gamma(D^+ \to f) - \Gamma(D^- \to \bar{f})}{\Gamma(D^+ \to f) + \Gamma(D^- \to \bar{f})}$$

G. Elor

# CPV in Charged D Decays

Example: Standard Model decays to an odd number of charged pions

$D^+$ decay mode	$A_{CP}^{f}/10^{-2}$	$D^+$ decay mode	$A_{CP}^{f}/10^{-2}$
$K_S^0 \pi^+$	$-0.41 \pm 0.09$	$\pi^+\eta$	$1.0 \pm 1.5$
$K^{-}\pi^{+}\pi^{+}$	$-0.18 \pm 0.16$	$\pi^+\eta'(958)$	$-0.6\pm0.7$
$K^{-}\pi^{+}\pi^{+}\pi^{0}$	$-0.3 \pm 0.6 \pm 0.4$	$K^+K^-\pi^+$	$0.37 \pm 0.29$
$K_S^0 \pi^+ \pi^0$	$-0.1 \pm 0.7 \pm 0.2$	$\phi\pi^+$	$0.01 \pm 0.09$
$K_S^0 \pi^+ \pi^+ \pi^-$	$0.0 \pm 1.2 \pm 0.3$	$a_0(1450)^0\pi^+$	$-19 \pm 12^{+8}_{-11}$
$\pi^+\pi^0$	$2.4 \pm 1.2$	$\phi(1680)\pi^{+}$	$-9 \pm 22 \pm 14$
$\pi^+\eta$	$1.0 \pm 1.5$	$\pi^+\pi^+\pi^-$	$-1.7 \pm 4.2$

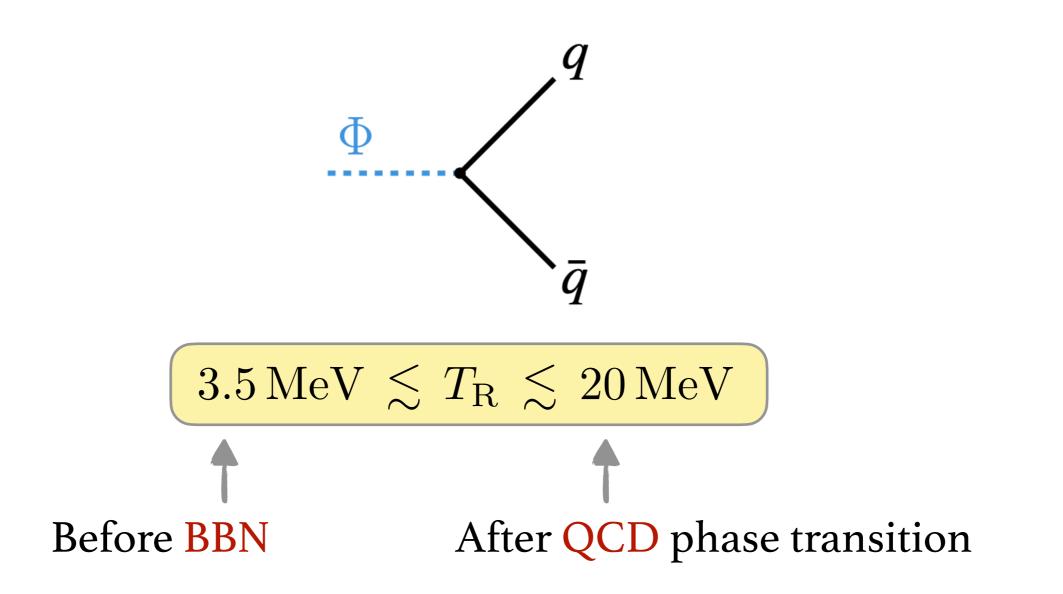
Not a small number if we want to explain

$$Y_B^{\text{obs}} = (8.718 \pm 0.004) \times 10^{-11}$$

### Sakharov I. Out of Equilibrium

Late decay of an "inflaton-like" field

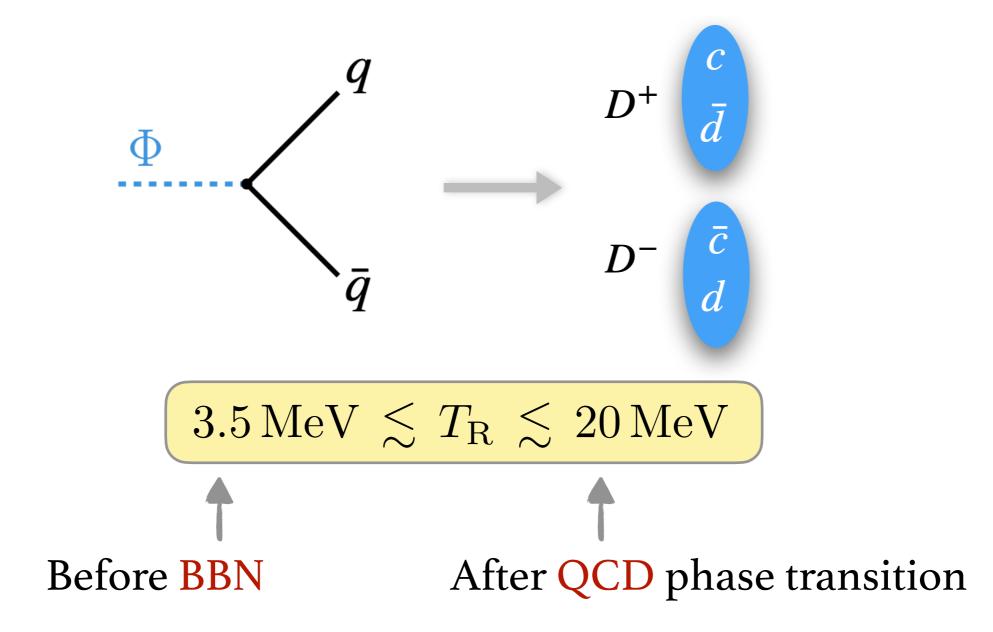
Decays at:  $\Gamma_{\Phi} = 4H(T_R)$  to quarks  $m_{\Phi} \in [5 \text{ GeV}, 100 \text{ GeV}]$ 



### Sakharov I. Out of Equilibrium

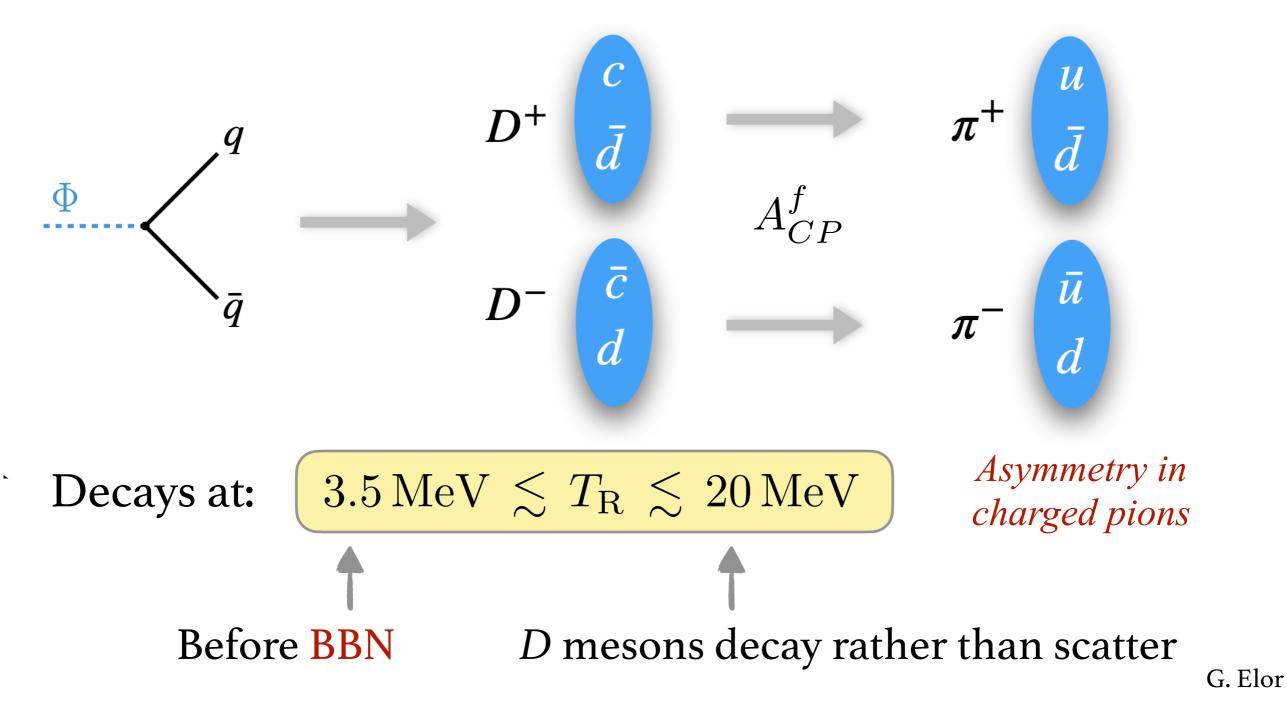
Late decay of an "inflaton-like" field

Decays at:  $\Gamma_{\Phi} = 4H(T_R)$  to quarks  $m_{\Phi} \in [5 \text{ GeV}, 100 \text{ GeV}]$ 



### Sakharov II. CP Violation

D mesons quickly undergo Standard Model decays to pions



## Sakharov III. B Violation?

Need a way to change baryon number



Hide baryon *and lepton* number in a dark sector without violating either.



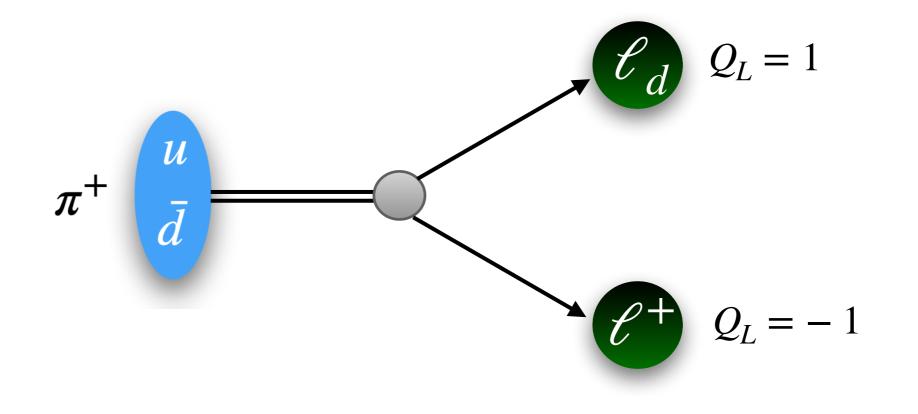
First generate a lepton asymmetry

#### Dark Sector Lepton

Portal Operator:

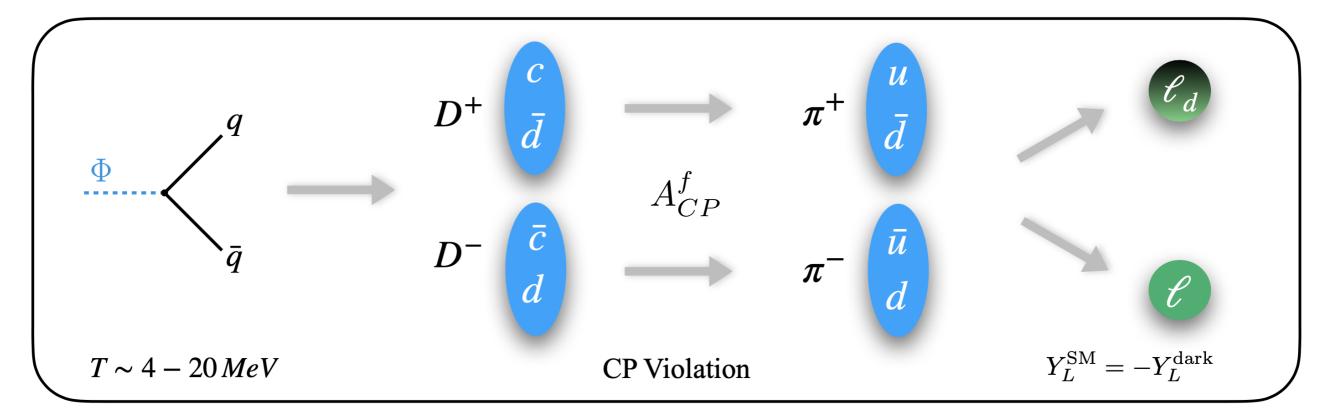
$$\mathcal{O} = \frac{1}{\Lambda^2} \Big[ \bar{d} \Gamma^{\mu} u \Big] \Big[ \bar{\ell}_d \Gamma_{\mu} \ell \Big] + \text{h.c.}$$

Pion Decays:  $\pi^+ \to \ell_d + \ell^+$ ,  $m_{\ell_d} < m_{\pi^+} - m_{\ell}$  Can be light



# Generating a Lepton Asymmetry

#### Equal and opposite dark/visible sector lepton asymmetry



$$Y_L^{\text{dark}} \equiv \left(\frac{n_{\ell_d} - n_{\bar{\ell_d}}}{s}\right) \propto \text{Br}\left(\pi^+ \to \ell_d + \ell^+\right) \sum_f A_{\text{CP}}^f \times \text{Br}\left(D^+ \to f\right)$$

## Generating a Baryon Asymmetry

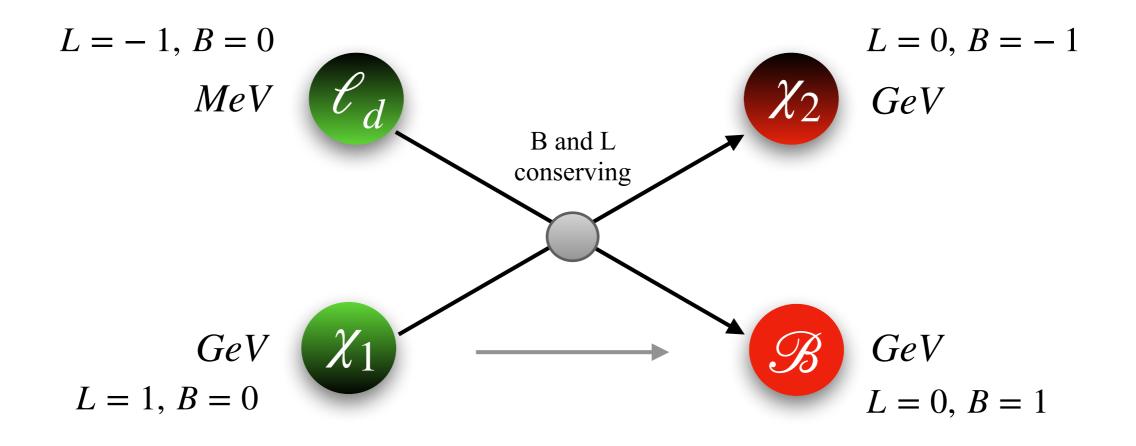
When you make the Universe at 20 MeV, you (of course) can not use Electroweak Sphalerons to transfer a lepton into a baryon asymmetry.

You also don't need them...

#### Freezing-In the Baryon Asymmetry

Scattering off dark sector states charged under SM L and/or B

$$\bar{\ell}_d + \chi_1 \to \chi_2 + \mathcal{B}$$



*Equal and opposite dark and visible sector baryon asymmetry* 

#### **Boltzmann Equations: Lepton Asymmetry**

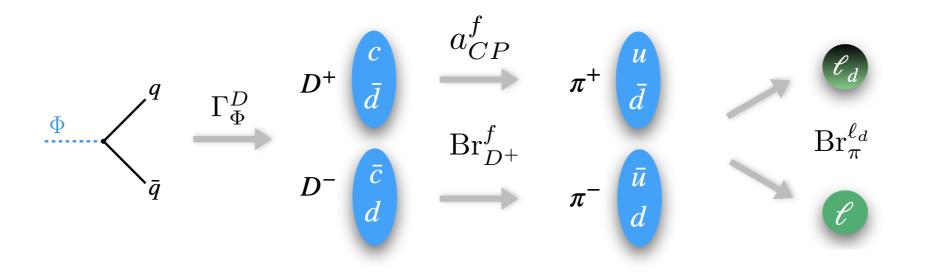
• Inflaton:

$$\frac{dn_{\Phi}}{dt} + 3Hn_{\Phi} = -\Gamma_{\Phi}n_{\Phi}$$

- Radiation:  $\frac{d\rho_{\rm rad}}{dt} + 4H\rho_{\rm rad} = +\Gamma_{\Phi}m_{\Phi}n_{\Phi}$
- Hubble:  $H^2 = \frac{8\pi}{3M_{\rm Pl}^2} \left( \rho_{\rm rad} + m_{\Phi} n_{\Phi} \right) \qquad \Gamma_{\Phi} = 4H \left( T_R \right)$
- The dark lepton asymmetry:

 $\Gamma_{\Phi}^{D} \equiv \Gamma_{\Phi} \mathrm{Br}(\Phi \to c) \mathrm{Br}(c \to D)$ 

$$\frac{d}{dt} \left( n_{\ell_d} - n_{\bar{\ell}_d} \right) + 3H \left( n_{\ell_d} - n_{\bar{\ell}_d} \right) = 2 \Gamma_{\Phi}^D n_{\Phi} \mathrm{Br}_{\pi}^{\ell_d} \sum_f N_{\pi}^f a_{CP}^f \mathrm{Br}_{D^+}^f$$



#### **Boltzmann Equations: Lepton Asymmetry**

• Inflaton:  $\frac{dn_{\Phi}}{dt}$  +

$$\frac{dn_{\Phi}}{dt} + 3Hn_{\Phi} = -\Gamma_{\Phi}n_{\Phi}$$

- Radiation:  $\frac{d\rho_{\rm rad}}{dt} + 4H\rho_{\rm rad} = +\Gamma_{\Phi}m_{\Phi}n_{\Phi}$
- Hubble:  $H^2 = \frac{8\pi}{3M_{\rm Pl}^2} \left(\rho_{\rm rad} + m_{\Phi}n_{\Phi}\right) \qquad \Gamma_{\Phi} = 4H\left(T_R\right)$
- The dark lepton asymmetry:  $\Gamma_{\Phi}^{D} \equiv \Gamma_{\Phi} Br(\Phi \to c) Br(c \to D)$

$$\frac{d}{dt} \left( n_{\ell_d} - n_{\bar{\ell}_d} \right) + 3H \left( n_{\ell_d} - n_{\bar{\ell}_d} \right) = 2 \Gamma_{\Phi}^D n_{\Phi} \mathrm{Br}_{\pi}^{\ell_d} \sum_f N_{\pi}^f a_{CP}^f \mathrm{Br}_{D^+}^f$$

#### **Experimental Observables:**

- SM charged D decays:
- Charged pion decays:

$$\begin{split} a_{CP}^{f} &\equiv A_{CP}^{f} / (1 + A_{CP}^{f}) \approx A_{CP}^{f} & LHCb, B \\ & factories \\ Br_{D^{+}}^{f} &\equiv Br \left( D^{+} \to f \right) \\ Br_{\pi}^{\ell_{d}} &\equiv Br \left( \pi^{+} \to \ell_{d} + \ell^{+} \right) & PIENU, PSI, etc. \\ & G. Elor \end{split}$$

#### **Boltzmann Equations: Lepton Asymmetry**

Numerically Solving for the Lepton asymmetry:

$$\frac{Y_L^{\text{dark}}}{Y_B^{\text{obs}}} \simeq \frac{\text{Br}_{\pi}^{\ell_d}}{10^{-3}} \frac{\sum_f N_{\pi}^f a_{CP}^f \text{Br}_{D^+}^f}{3 \times 10^{-5}} \frac{T_R}{20 \,\text{MeV}} \frac{10 \,\text{GeV}}{m_{\Phi}}$$

$$Y_B^{\text{obs}} = (8.718 \pm 0.004) \times 10^{-11}$$

#### Experimental Observables:

- SM charged D decays:
- Charged pion decays:

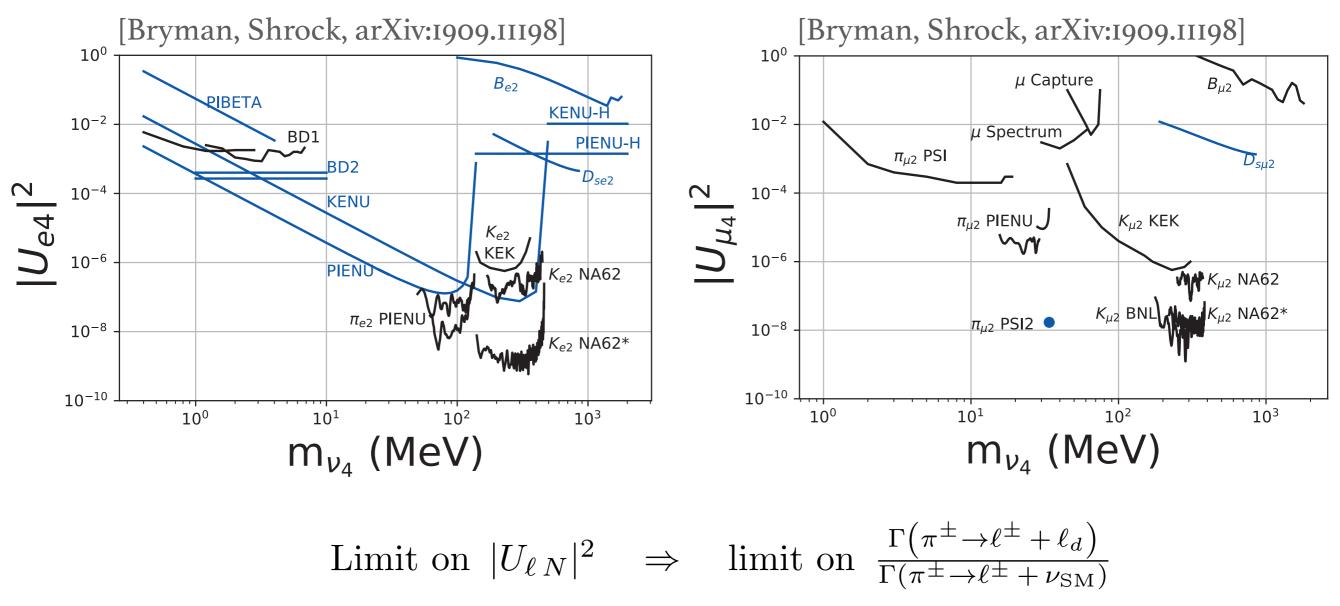
 $a_{CP}^{f} \equiv A_{CP}^{f} / (1 + A_{CP}^{f}) \approx A_{CP}^{f} \qquad LHCb, B$  factories  $Br_{D^{+}}^{f} \equiv Br (D^{+} \to f)$   $Br_{\pi}^{\ell_{d}} \equiv Br (\pi^{+} \to \ell_{d} + \ell^{+}) \qquad PIENU, PSI, etc.$  G. Elor

## Limits on D Decays

$D^+$ decay mode	$A_{CP}^{f}/10^{-2}$	${\rm Br}_{D^+}^f / 10^{-2}$
$K_S^0 \pi^+$	$-0.41 \pm 0.09$	$1.562\pm0.031$
$K^-\pi^+\pi^+$	$-0.18 \pm 0.16$	$9.38\pm0.16$
$K^-\pi^+\pi^+\pi^0$	$-0.3 \pm 0.6 \pm 0.4$	$5.98 \pm 0.08 \pm 0.16^{*}$
$K_S^0 \pi^+ \pi^0$	$-0.1 \pm 0.7 \pm 0.2$	$6.99 \pm 0.09 \pm 0.25^*$
•	•	•

$$\sum_{f} N_{\pi}^{f} a_{CP}^{f} \operatorname{Br}_{D^{+}}^{f} = \left(-9.3 \times 10^{-4}\right)_{-0.0039}^{+0.0031}$$

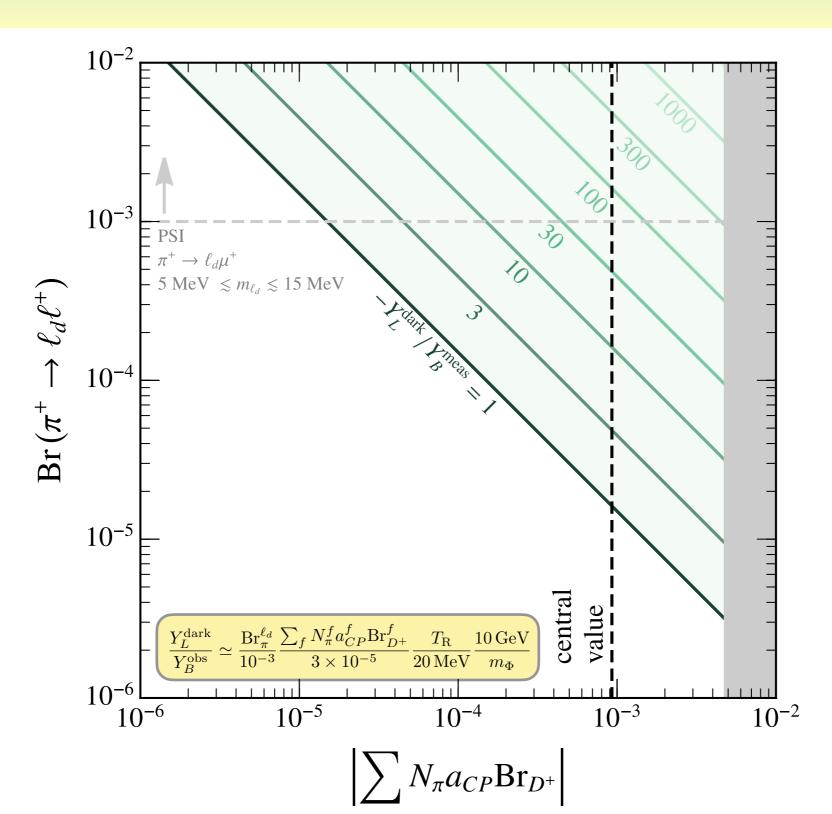
## Limits on Pion Decays



[Shrock, Phys. Rev. D24, 1232 (1981)]

 $Br(\pi^{\pm} \to \mu^{\pm} + MET) \lesssim 10^{-3}$ , for  $5 \,MeV < m_{\ell_d} < 15 \,MeV$ .

### Generating a Lepton Asymmetry



### Freezing-In a Baryon Asymmetry

### Boltzmann Equations with scattering: $\bar{\ell}_d + \chi_1 \rightarrow \chi_2 + B$

- New dark lepton/lepto-baryon:  $m_{\Phi} \gtrsim m_{\chi_1}$   $m_{\Phi} \gtrsim m_{\chi_2} + m_{\mathcal{B}}$  $\frac{dn_{\chi_1}}{dt} + 3Hn_{\chi_1} = \Gamma_{\Phi} n_{\Phi} \operatorname{Br} \left(\Phi \to \chi_1 \bar{\chi}_1\right) - \langle \sigma v \rangle n_{\bar{\ell}_d} n_{\chi_1}$
- Dark lepton:

$$\frac{d}{dt}\left(n_{\ell_d} - n_{\bar{\ell}_d}\right) + 3H\left(n_{\ell_d} - n_{\bar{\ell}_d}\right) = 2\Gamma_{\Phi}^D n_{\Phi} \mathrm{Br}_{\pi}^{\ell_d} \sum_f N_{\pi}^f a_{CP}^f \mathrm{Br}_{D^+}^f - \langle \sigma v \rangle n_{\chi_1} \left(n_{\ell_d} - n_{\bar{\ell}_d}\right)$$

• Baryon asymmetry:

$$\frac{d}{dt}\left(n_{\mathcal{B}} - n_{\overline{\mathcal{B}}}\right) + 3H\left(n_{\mathcal{B}} - n_{\overline{\mathcal{B}}}\right) = -\langle \sigma v \rangle n_{\chi_1}\left(n_{\ell_d} - n_{\overline{\ell}_d}\right)$$

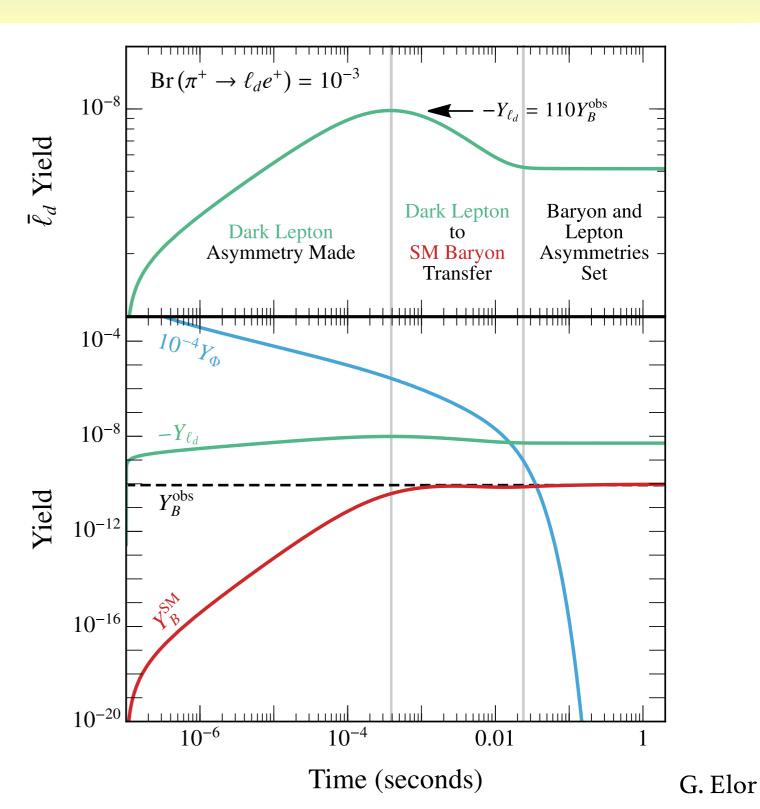
To efficiently transfer the asymmetry

$$\frac{n_{\chi_1} \langle \sigma v \rangle}{H(T)} \Big|_{T=T_R} \gtrsim \frac{Y_B^{\text{obs}}}{Y_L^{\text{dark}}}$$

### Freezing-In a Baryon Asymmetry

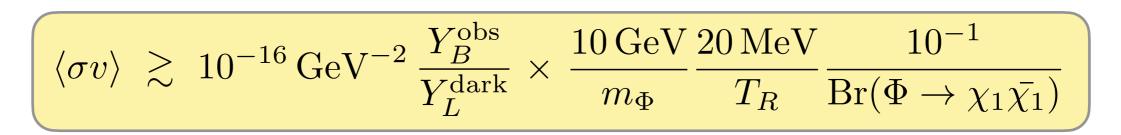
#### Example Benchmark point:

$$T_R = 10 \text{ MeV}, \ m_{\Phi} = 6 \text{ GeV}$$
$$\langle \sigma v \rangle = 1 \times 10^{-15} \text{ GeV}^{-2}$$
$$\operatorname{Br} \left( \Phi \to \chi_1 \bar{\chi}_1 \right) = 0.1$$
$$\sum_f N_{\pi}^f a_{CP}^f \operatorname{Br}_{D^+}^f = \left( -9.3 \times 10^{-4} \right)^{\frac{1}{2}}$$



## Freezing-In a Baryon Asymmetry

### Numerically:



How realistic is this?

# Models

# Proof of concept that what I have told you thus far is not (too) crazy.

• Some example models/dark sector charge assignments.

$$\bar{\ell}_d + \chi_1 \to \chi_2 + \mathcal{B}$$

• Estimation of the scattering cross section to confirm it can be large enough to transfer the asymmetry given current constraints.

$$\langle \sigma v \rangle \gtrsim 10^{-16} \,\mathrm{GeV}^{-2} \, \frac{Y_B^{\mathrm{obs}}}{Y_L^{\mathrm{dark}}} \times \frac{10 \,\mathrm{GeV}}{m_\Phi} \frac{20 \,\mathrm{MeV}}{T_R} \frac{10^{-1}}{\mathrm{Br}(\Phi \to \chi_1 \bar{\chi_1})}$$

## Portal to the Dark Sector

#### Model Build for:

$$\bar{\ell}_d + \chi_1 \to \chi_2 + \mathcal{B}$$

New fields: (Same model as for *B*-Mesogenesis[arXiv:1810.00880])

	Field	Spin	L	В	$\mathbb{Z}_2$	Mass
Color triplet scalar mediator	Y	0	0	-2/3	+1	$\gtrsim 1{ m TeV}$
	$\ell_d$	1/2	1	0	+1	$\mathcal{O}(10-140{ m MeV})$
Dark Baryon	$\psi_B$	1/2	0	-1	+1	$\gtrsim 1.2{ m GeV}$

 $\mathcal{L}_{\text{eff}} = \frac{y^2}{M_{-}^2} \bar{u}_i^c d_j d_k^c \psi_B$ 

Collider bounds (as just discussed)

Stability of matter, neutron star bounds

Allowed Interactions:

$$\mathcal{L} \supset y_{u_i d_j} Y^* \bar{u}_i d_j^c + y_{\psi d_k} Y \bar{\psi}_B d_k^c + h.c.$$

dark baryon-SM

baryon "mixing"

## Dark Possibilities

$$\bar{\ell}_d + \chi_1 \to \chi_2 + \bar{\psi}_B$$

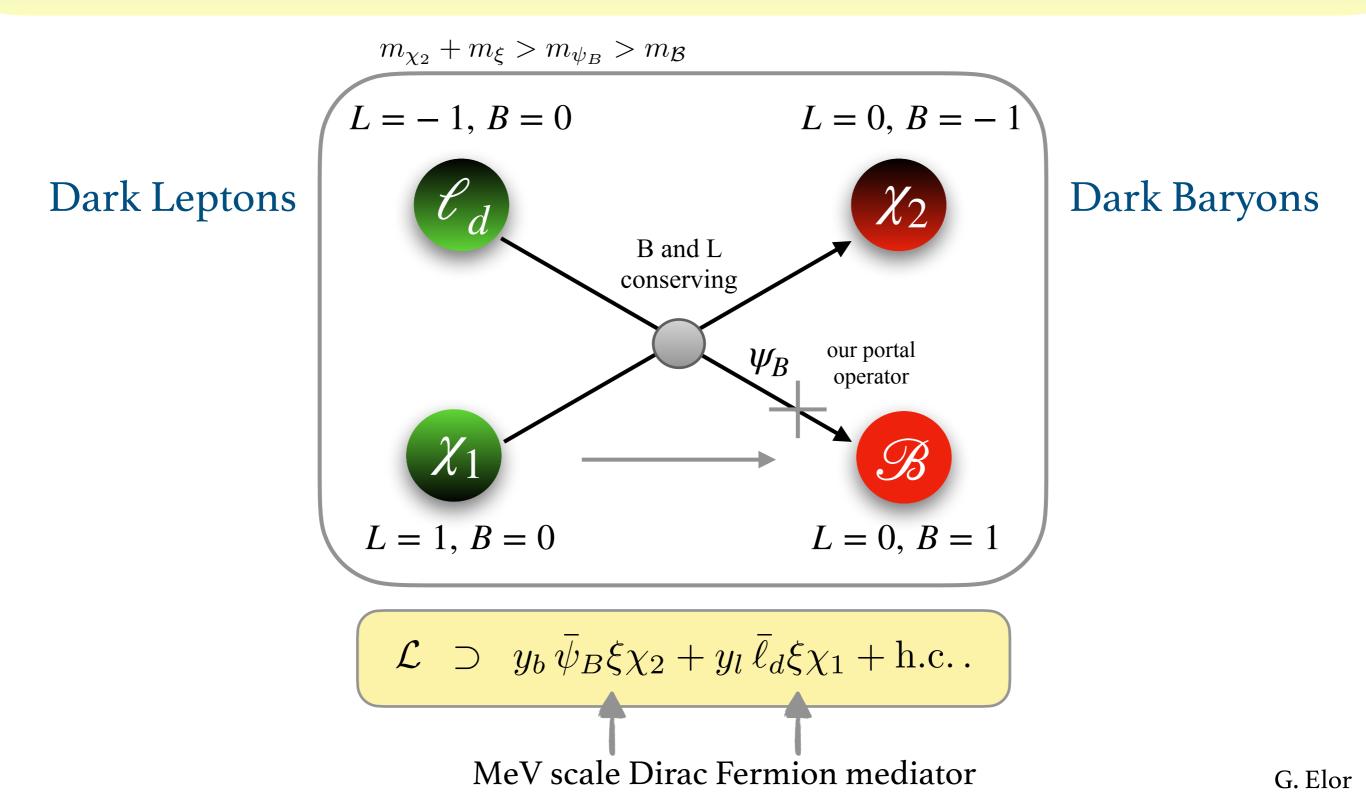
Field	L	В	Field	L	В
$\chi_1$	1	0	$\chi_1$	1	1
$\chi_2$	0	-1	$\chi_2$	0	0
$\chi_1$	0	1	$\chi_1$	0	0
$\chi_2$	1	0	$\chi_2$	-1	-1

## Dark Possibilities

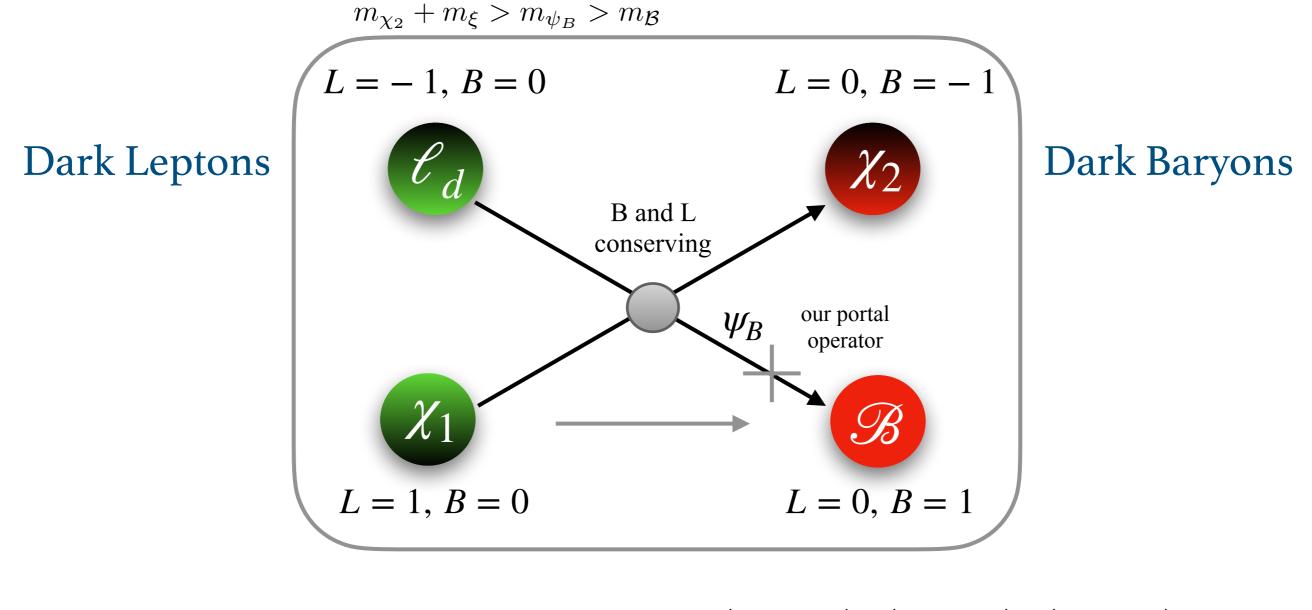
$$\bar{\ell}_d + \chi_1 \to \chi_2 + \bar{\psi}_B$$

Field	L	В	Field	L	В
$\chi_1$	1	0	$\chi_1$	1	1
$\chi_2$	0	-1	$\chi_2$	0	0
$\chi_1$	0	1	$\chi_1$	0	0
$\chi_2$	1	0	$\chi_2$	-1	-1

## Example Charge Assignment



## Example Charge Assignment



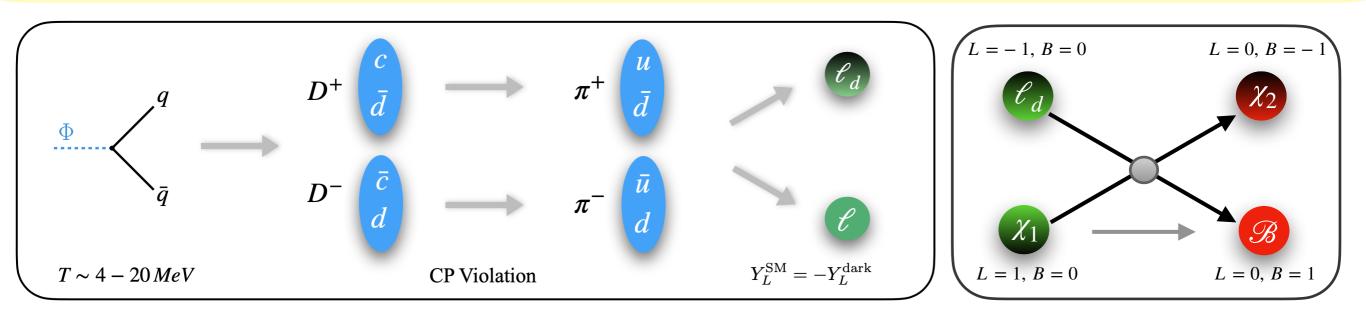
$$\langle \sigma v \rangle \simeq 10^{-15} \,\mathrm{GeV}^{-2} \, (y_l \, y_b)^2 \, \times \left(\frac{10 \,\mathrm{MeV}}{m_{\ell_d}}\right) \left(\frac{20 \,\mathrm{GeV}}{m_{\chi_1}}\right) \left(\frac{10 \,\mathrm{GeV}}{m_{\chi_2}}\right)$$

## Dark Matter Generation

### A few comments

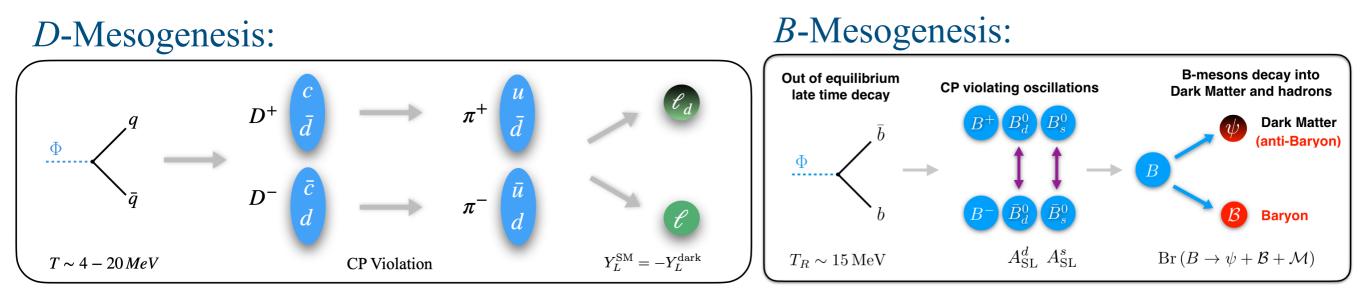
- Details depend on baryon and lepton number charge assignment.
- Symmetric component of any dark baryons will always contribute to some or all of the dark matter.
- Dark states will tend to be overproduced but additional dark sector dynamics can always be at play to achieve the observed dark matter relic abundance.

# **D-Mesogenesis**



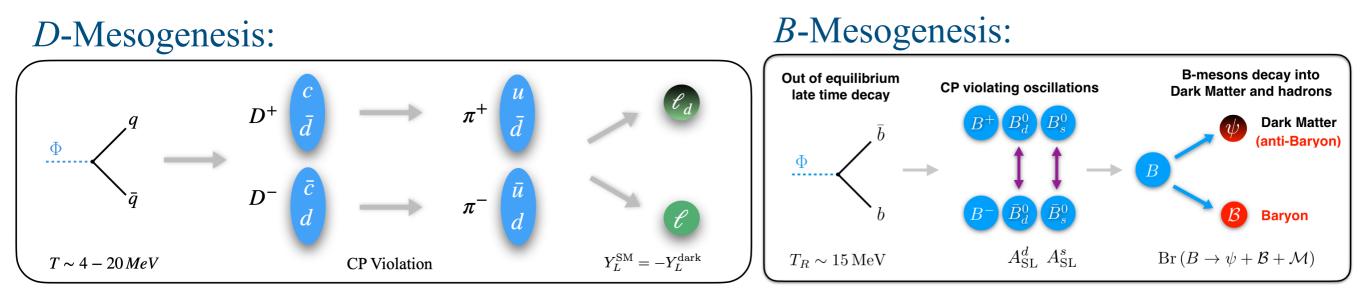
- First generates a lepton asymmetry and then freezes in a baryon asymmetry through dark sector scatterings.
- Baryogenesis and dark matter production are controlled by experimental observables of the charged *D* Mesons system.
- Upcoming experimental probes will better constrain or discover this mechanism.

# Outlook



- Continued support of experimental efforts to probe Mesogenesis e.g. proper form factor calculation in the case of decays relevant for *B*-Mesogenesis.
- Mesogenesis in other meson systems?
- Explore UV embeddings (dark sector? mediators?) of both *B* and *D*-Mesogenesis and associated phenomenology.
- Theory of inflation and flavor.

# Outlook



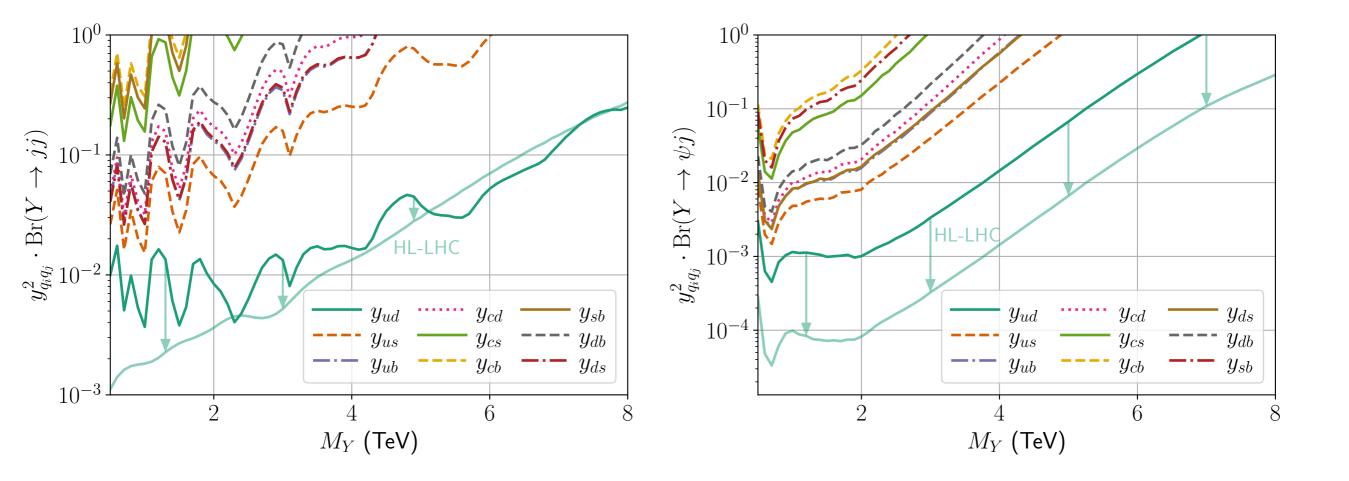
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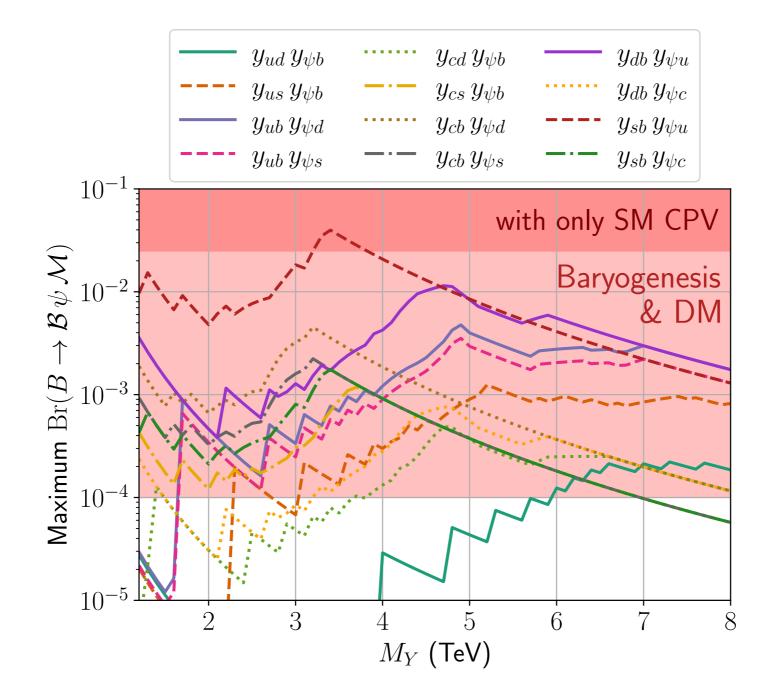
## **Colored Triplet Scalar**

Constraints from LHC squark searches

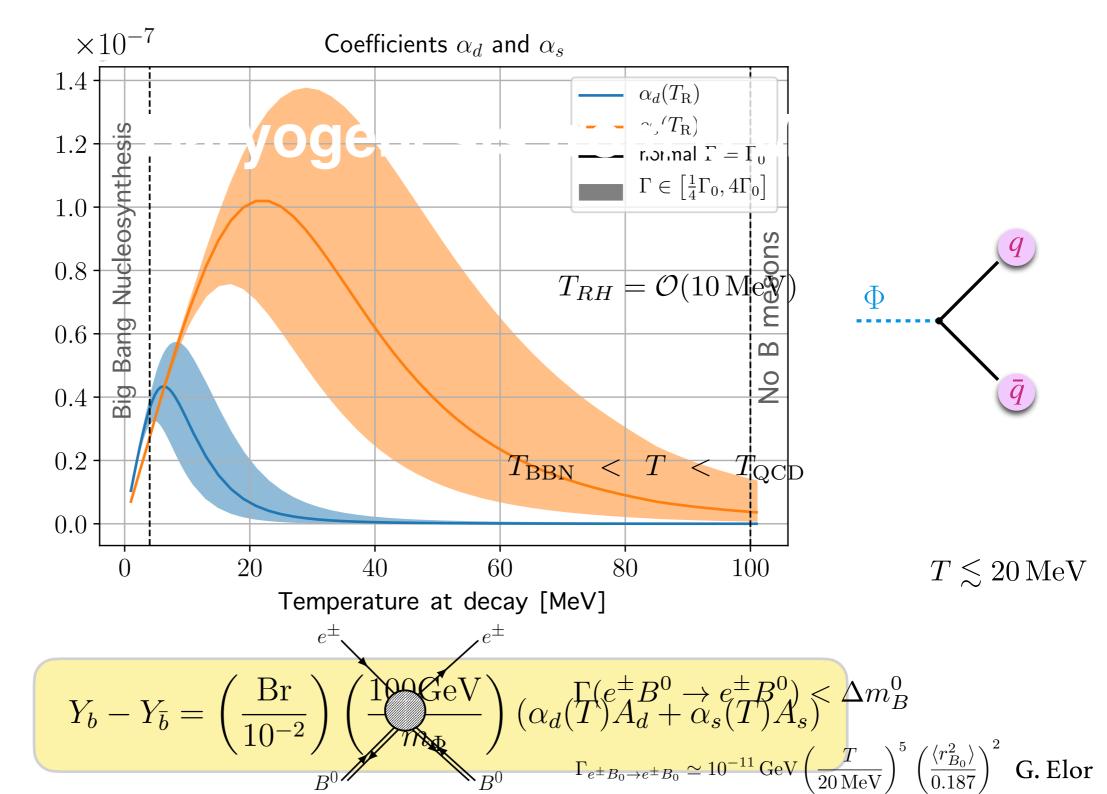


## **Colored Triplet Scalar**

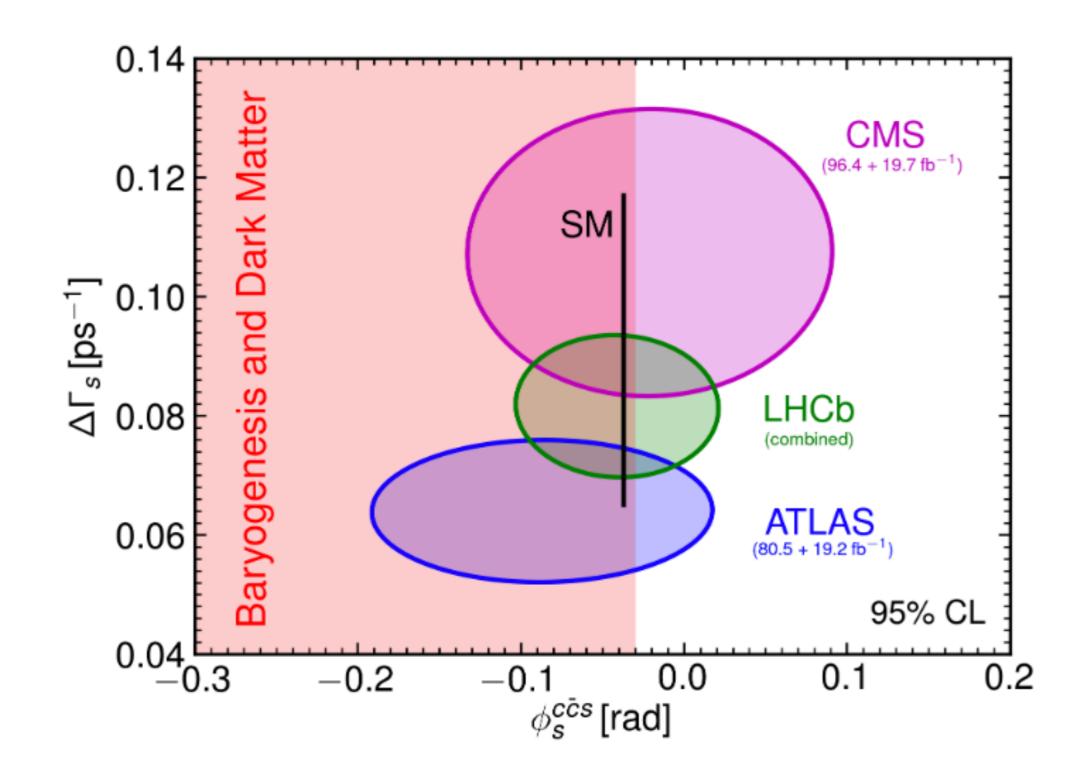
#### Constraints from LHC squark searches



### Baryogenesis and Dark Matter from B Mesons



### **CP Observables**



### The Semi-Leptonic Asymmetry

$$A_{\rm SL}^q = \frac{\Gamma\left(\bar{B}_q^0 \to B_q^0 \to f\right) - \Gamma\left(B_q^0 \to \bar{B}_q^0 \to \bar{f}\right)}{\Gamma\left(\bar{B}_q^0 \to B_q^0 \to f\right) + \Gamma\left(B_q^0 \to \bar{B}_q^0 \to \bar{f}\right)} = -\left|\frac{\Gamma_{12}^q}{M_{12}^q}\right| \sin\left(\phi_{12}^q\right)$$

From SM box diagrams:

 $A_{\rm SL}^d|_{\rm SM} = (-4.7 \pm 0.4) \times 10^{-4}$  $A_{\rm SL}^s|_{\rm SM} = (2.1 \pm 0.2) \times 10^{-5}$ 

Lenz, Tetlalmatzi-Xolocotzi [1912.07621]

World average: As

$${}^{d}_{\rm SL} = (-2.1 \pm 1.7) \times 10^{-3}$$
  
 ${}^{s}_{\rm SL} = (-0.6 \pm 2.8) \times 10^{-3}$ 

HFLAG

Projected Sensitivities:

$$\begin{split} &\delta A_{\rm SL}^s = 10 \times 10^{-4} \; \left[ {\rm LHCb} \left( {33\,{\rm fb}^{-1}} \right) - 2025 \right] & [1812.07638, \\ &\delta A_{\rm SL}^s = 3 \times 10^{-4} \; \left[ {\rm LHCb} \left( {300\,{\rm fb}^{-1}} \right) - 2040 \right] & 1808.08865 \right] \\ &\delta A_{\rm SL}^d = 8 \times 10^{-4} \; \left[ {\rm LHCb} \left( {33\,{\rm fb}^{-1}} \right) - 2025 \right] \\ &\delta A_{\rm SL}^d = 2 \times 10^{-4} \; \left[ {\rm LHCb} \left( {300\,{\rm fb}^{-1}} \right) - 2040 \right] \\ &\delta A_{\rm SL}^d = 5 \times 10^{-4} \; \left[ {\rm Belle} \; {\rm II} \left( {50\,{\rm ab}^{-1}} \right) - 2025 \right] \end{split}$$

## Searching for new b-Hadron Decays

#### Caution: Inclusive vs. Exclusive Rates

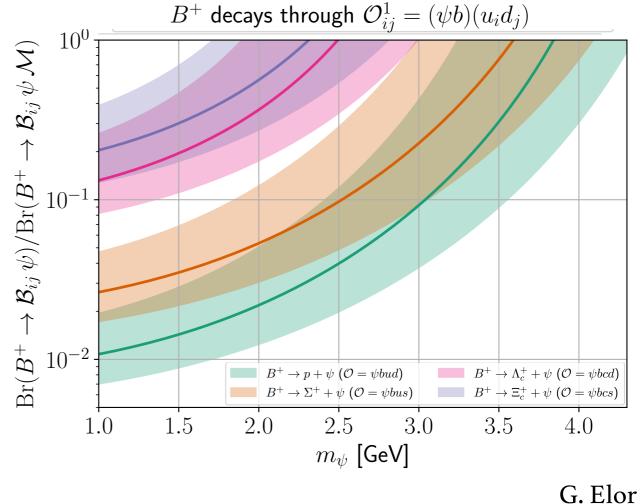
- All decays (and their searches) discussed thus far have been *exclusive*. But, the observable controlling the baryon asymmetry is an *inclusive* rate.
  - $Br(B \to \psi \,\mathcal{B} \,\mathcal{M}) \gtrsim 10^{-4}$
- Need a dedicated calculation using QCD sum rules or lattice techniques etc. to calculate form factors. Beyond my current expertise....
- Phase space method

$$\begin{bmatrix} \text{Bigi, Phys.Lett.B 106, 510 (1981)} \end{bmatrix}$$

$$\bar{b} \to u_i d_j \psi \qquad \gamma(\Lambda) \equiv \int_{(m_{u_i} + m_{d_j})^2}^{\Lambda^2} \frac{\partial \Gamma}{\partial M_{u_i d_j}^2} \, \mathrm{d}M_{u_i d_j}^2$$

$$\frac{\mathrm{Br}(B \to \mathcal{B}_{ij} + \psi)}{\mathrm{Br}(B \to \mathcal{B}_{ij} + \psi + \mathcal{M})} \simeq \frac{\gamma(m_{\mathcal{B}_{ij}})}{\gamma(m_b - m_{\psi})}$$

$$\frac{\mathrm{Br}(B \to \psi \mathcal{B})}{\mathrm{Br}(B \to \psi \mathcal{B} \mathcal{M})} \gtrsim (1 - 10) \%.$$



## Flavorful Variations

No a priori reason to expect a particular flavor structure.

Most general interactions:

$$\mathcal{L}_{-1/3} = -\sum_{i,j} y_{u_i d_j} Y^* \bar{u}_{iR} d_{jR}^c - \sum_k y_{\psi d_k} Y d_{kR}^c \bar{\psi} + \text{h.c.}$$
$$\mathcal{L}_{2/3} = -\sum_{i,j} y_{d_i d_j} Y^* \bar{d}_{iR} d_{jR}^c - \sum_k y_{\psi u_k} Y u_{kR}^c \bar{\psi} + \text{h.c.}$$

Possible operators:

$$\mathcal{O}_{ud} = \psi \, b \, u \, d$$
$$\mathcal{O}_{us} = \psi \, b \, u \, s$$
$$\mathcal{O}_{cd} = \psi \, b \, c \, d$$
$$\mathcal{O}_{cs} = \psi \, b \, c \, s$$

#### New meson and b-flavored hadron decays:

Operator/Decay	Initial State	Final state	$\Delta M \ ({\rm MeV})$
$\mathcal{O} = \psi  b  u  d$	$B_d$ $B_s$	$egin{aligned} \psi + n  (udd) \ \psi + \Lambda  (uds) \end{aligned}$	$\begin{array}{c} 4340.07 \\ 4251.21 \end{array}$
$\overline{b} \to \psi  u  d$	$B^+$ $\Lambda_b$	$\psi + p \left( duu  ight) \ ar{\psi} + \pi^0$	4341.05 5484.5
	$B_d$	$\psi + \Lambda(usd)$	4163.95
$\mathcal{O} = \psi  b  u  s$	$B_s$ $B^+$	$\psi + \Xi^0 (uss)$	4025.03
$b \rightarrow \psi  u  s$	$\frac{B}{\Lambda_b}$	$\frac{\psi + \Sigma^+ (uus)}{\bar{\psi} + K^0}$	4089.95 5121.9
	$B_d$	$\psi + \Lambda_c + \pi^- \left( c d d \right)$	2853.60
$\mathcal{O} = \psi  b  c  d$	$B_s$	$\psi + \Xi_c^0  (cds)$	2895.02
$b \rightarrow \psi  c  d$	$B^+$	$\psi + \Lambda_c (dcu)$	2992.86
	$\Lambda_b$	$ar{\psi}+\overline{D}^0$	3754.7
	$B_d$	$\psi + \Xi_c^0  (csd)$	2807.76
$\mathcal{O} = \psi  b  c  s$	$B_s$	$\psi + \Omega_c \ (css)$	2671.69
$\overline{b}  o \psi  c  s$	$B^+$	$\psi + \Xi_c^+ (csu)$	2810.36
	$\Lambda_b$	$\bar{\psi} + D^- + K^+$	3256.2

## **Colored Triplet Scalar**

#### Constraints from LHC squark searches

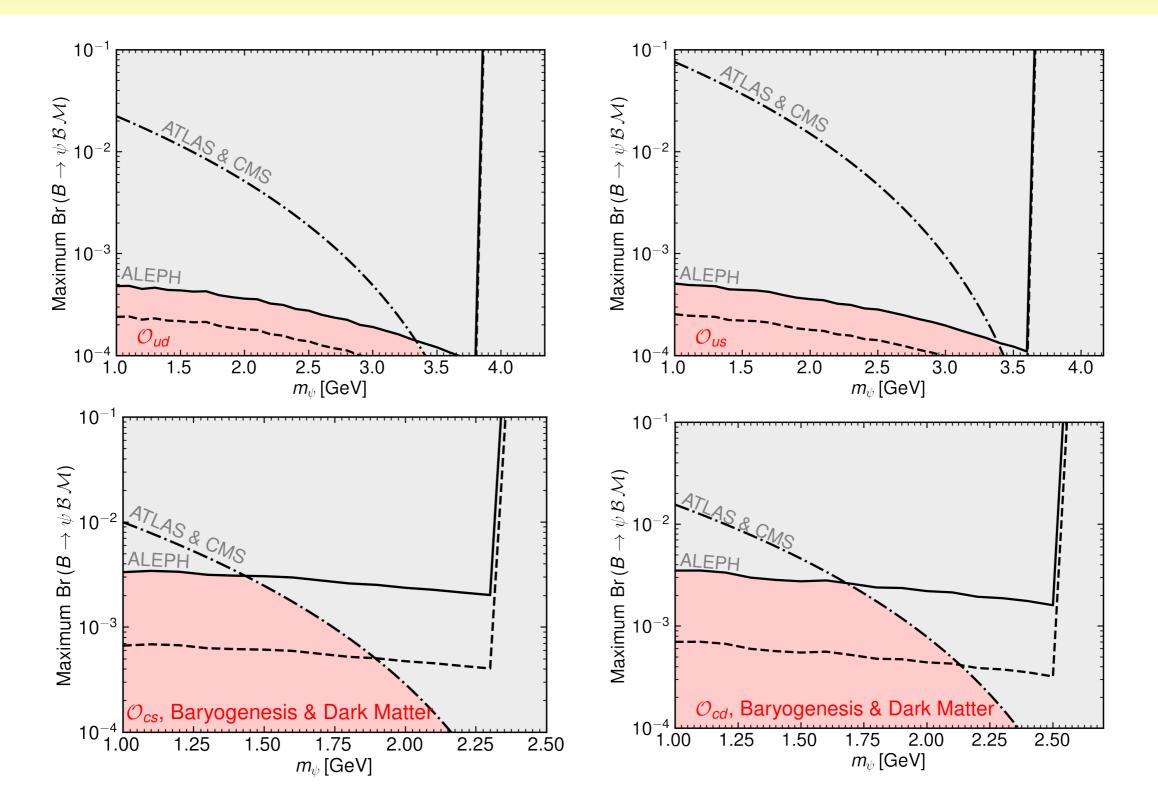
	Max:	Upper limit on $m_{\psi}$				
Operator	$\left[ \sqrt{y^2}  \underline{1.5  \mathrm{TeV}} \right]^4$	Inclusive $\operatorname{Br}(B \to \psi \mathcal{B} \mathcal{M})$				
	$\left\lfloor \overline{0.53} \ \overline{M_Y} \right\rfloor$	$10^{-4}$	$10^{-3}$	$10^{-2}$	2.5%	
$\mathcal{O}_{ud}^1 = (\psi  b)  (u  d)$	0.3	2.0	_	-	-	
$\mathcal{O}_{ud}^2 = (\psi  d)  (u  b)$	6.7	3.2	2.3	-	-	
$\mathcal{O}_{ud}^3 = (\psi  u)  (d  b)$	16.2	3.4	2.8	1.6	-	
$\mathcal{O}_{us}^1 = (\psi  b)  (u  s)$	2.4	2.7	1.6	-	-	
$\mathcal{O}_{us}^2 = \left(\psi  s\right) \left(u  b\right)$	6.7	3.2	2.3	-	-	
$\mathcal{O}_{us}^3 = (\psi  u)  (s  b)$	75.8	3.5	3.0	2.2	1.8	
$\mathcal{O}_{cd}^1 = (\psi  b)  (c  d)$	10.4	2.0	1.2	-	-	
$\mathcal{O}_{cd}^2 = (\psi  d)  (c  b)$	96.6	2.4	1.9	1.2	-	
$\mathcal{O}_{cd}^{3} = \left(\psi  c\right) \left(d  b\right)$	16.2	2.1	1.4	-	-	
$\mathcal{O}_{cs}^1 = (\psi  b)  (c  s)$	50.9	2.0	1.5	-	-	
$\mathcal{O}_{cs}^{2} = \left(\psi  s\right) \left(c  b\right)$	96.6	2.4	1.9	1.2	-	
$\mathcal{O}_{cs}^{3} = \left(\psi  c\right) \left(s  b\right)$	75.8	2.1	1.7	-	-	

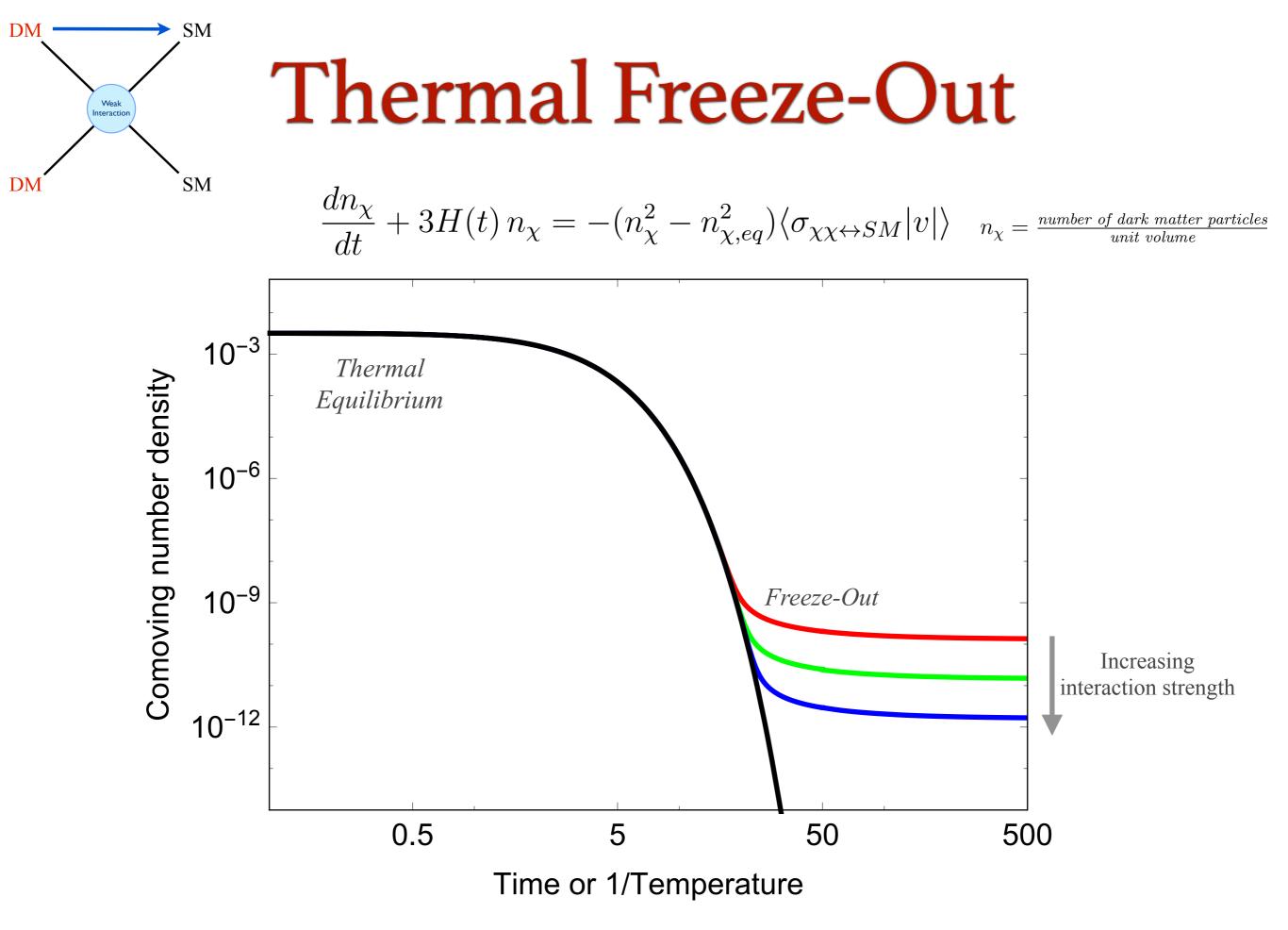
Bounds will depend on the product of coupling

$$\sqrt{y^2}/M_Y$$

Which in tern constrain branching fraction of relevance to the baryon asymmetry G. Elor

### **Colored Triplet Scalar**





## A Supersymmetric Theory

#### MSSM, R Symmetry, and Dirac Gauginos and Sterile Neutrios

Superfield	R-Charge	L no.
$\mathbf{U}^c, \mathbf{D}^c$	2/3	0
Q	4/3	0
$\mathbf{H}_{u},\mathbf{H}_{d}$	0	0
$\mathbf{R}_u, \mathbf{R}_d$	2	0
S	0	0
L	1	1
$\mathbf{E}^{c}$	1	-1
$\mathbf{N}_{R}^{c}$	1	-1

"RPV" 
$$\mathbf{W} = y_u \mathbf{Q} \mathbf{H}_u \mathbf{U}^c - y_d \mathbf{Q} \mathbf{H}_d \mathbf{D}^c - y_e \mathbf{L} \mathbf{H}_d \mathbf{E}^c + \frac{1}{2} \lambda_{ijk}^{"} \mathbf{U}_i^c \mathbf{D}_j^c \mathbf{D}_k^c$$
  
  $+ \mu_u \mathbf{H}_u \mathbf{R}_d + \mu_d \mathbf{R}_u \mathbf{H}_d$   
  $+ \lambda_u^t \mathbf{H}_u \mathbf{T} \mathbf{R}_d + \lambda_d^t \mathbf{R}_u \mathbf{T} \mathbf{H}_d + \lambda_d^s \mathbf{S} \mathbf{R}_u \mathbf{H}_d$ .  
  $\boldsymbol{\mathcal{L}} \stackrel{:}{=} \lambda_{113}^{"} \left( \tilde{d}_R^* u_R^\dagger b_R^\dagger + \tilde{u}_R^* d_R^\dagger b_R^\dagger + \tilde{b}_R^* u_R^\dagger d_R^\dagger \right) ,$   
Gauge:  
  $\mathcal{L}_{gauge} = -\sqrt{2}g(\phi T^a \psi^\dagger) \lambda^{a\dagger} + h.c.$ 

 $\Rightarrow -\sqrt{2}g(\tilde{d}_R^* d_R \tilde{B}^\dagger) - \sqrt{2}g(\tilde{d}_L d_L^\dagger \tilde{B}^\dagger) + \text{h.c.}$ 

Neutrio:

$$\mathbf{W} = \frac{\lambda_N}{4} \mathbf{S} \mathbf{N}_R^c \mathbf{N}_R^c + \mathbf{H}_u \mathbf{L}^i y_N^{ij} \mathbf{N}_R^{c,j} + \frac{1}{2} \mathbf{N}_R^c M_M \mathbf{N}_R^c + \text{h.c.},$$
$$\mathbf{W} = \frac{\lambda_N}{4} \mathbf{N}_R (\lambda_s \nu_R^\dagger \tilde{\nu}_R^* + \phi_s \nu_R^\dagger \nu_R^\dagger) + \text{h.c.}$$

Parameter space: "RPV" couplings and squark mass mixing

## A Supersymmetric Theory

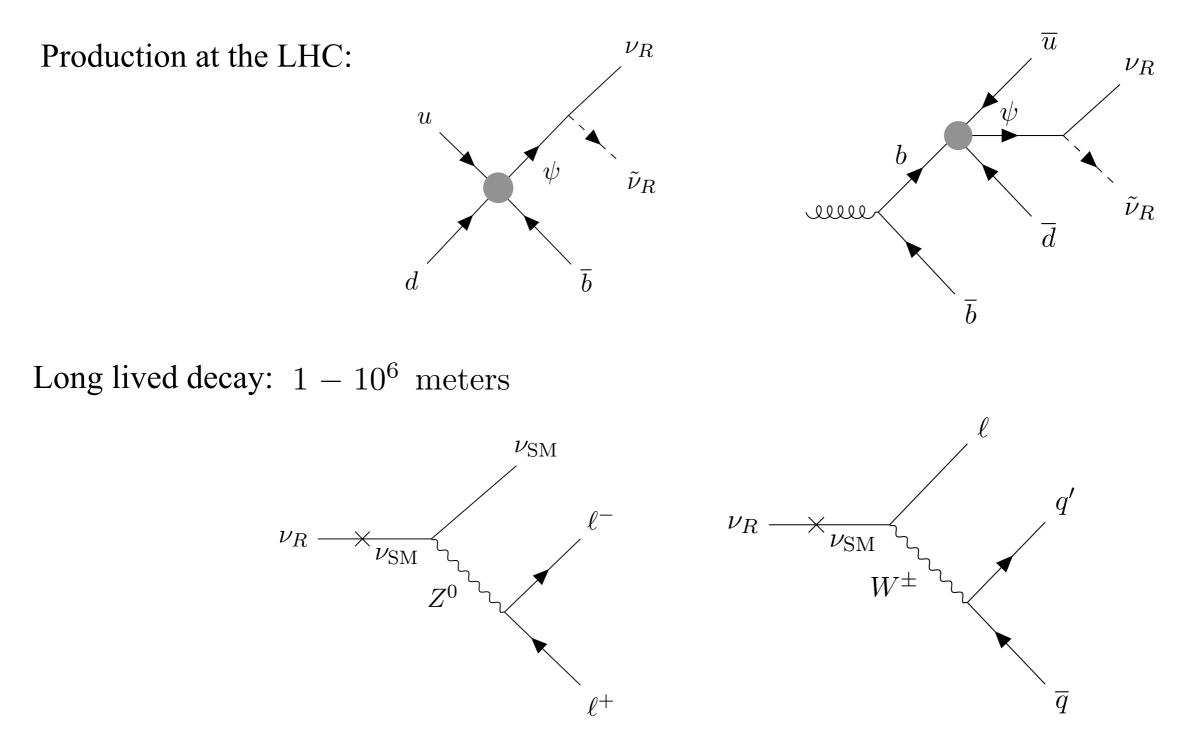
Superpartners and SM particles have different charge under an unbroken R-symmetry. We can identify this with Baryon number.

-----> Super

Superpartners as dark baryons.

	Field	Spin	$Q_{EM}$	Baryon no.	$\mathbb{Z}_2$	Mass
	Φ	0	0	0	+1	$11 - 100 \mathrm{GeV}$
MSSM Squark	$\tilde{d}_R$	0	-1/3	-2/3	+1	$\mathcal{O}({ m TeV})$
Dirac Bino	$\left[\begin{array}{c} \tilde{B} \\ \lambda_s^{\dagger} \end{array}\right]$	1/2	0	-1	+1	$\mathcal{O}({ m GeV})$
Right handed	$ u_R $	1/2	0	0	-1	$\mathcal{O}({ m GeV})$
neutrino multiplet	$\tilde{ u}_R$	0	0	-1	-1	$\mathcal{O}({ m GeV})$

## Model Specific Signals



[G. Alonso-Alvarez, G. Elor, A. E. Nelson, H. Xiao JHEP [arXiv:1907.10612]] [G. Alonso-Alvarez, G. Elor, M. Escudero, D. McKeen [*in preparation*]]