# The Hubble Tension: A Particle Physics Perspective

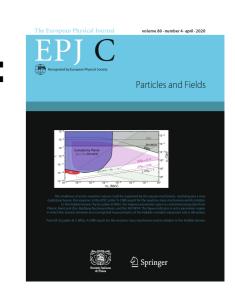
#### Miguel Escudero Abenza

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short review + original work based on:

<u>ArXiv:1909.04044</u>, EPJC 80 (2020) 4, 294 <u>ArXiv:2004.01470</u>, NuPhys19 Proceedings

ArXiv:2103.03249, EPJC 81 (2021) 6, 515



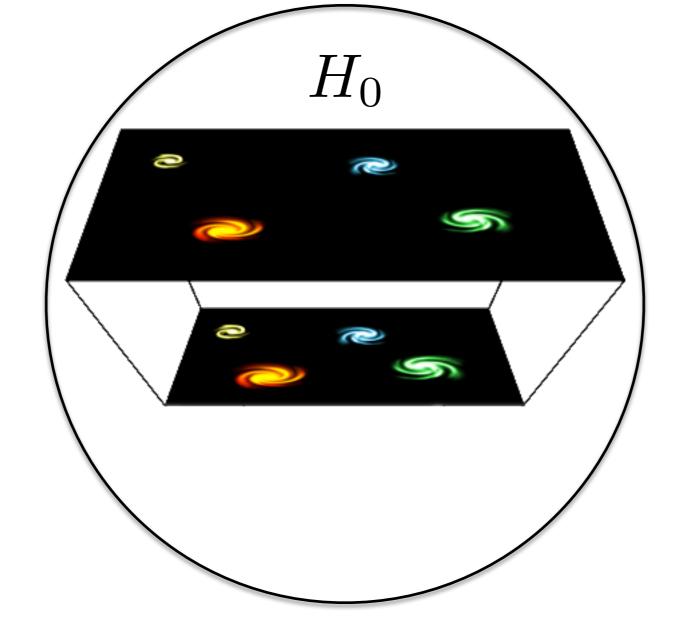
with Sam Witte



**Mainz** 02-11-2021

Unterstützt von / Supported by





Riess et al. 2012.08534

**Local Measurements** 

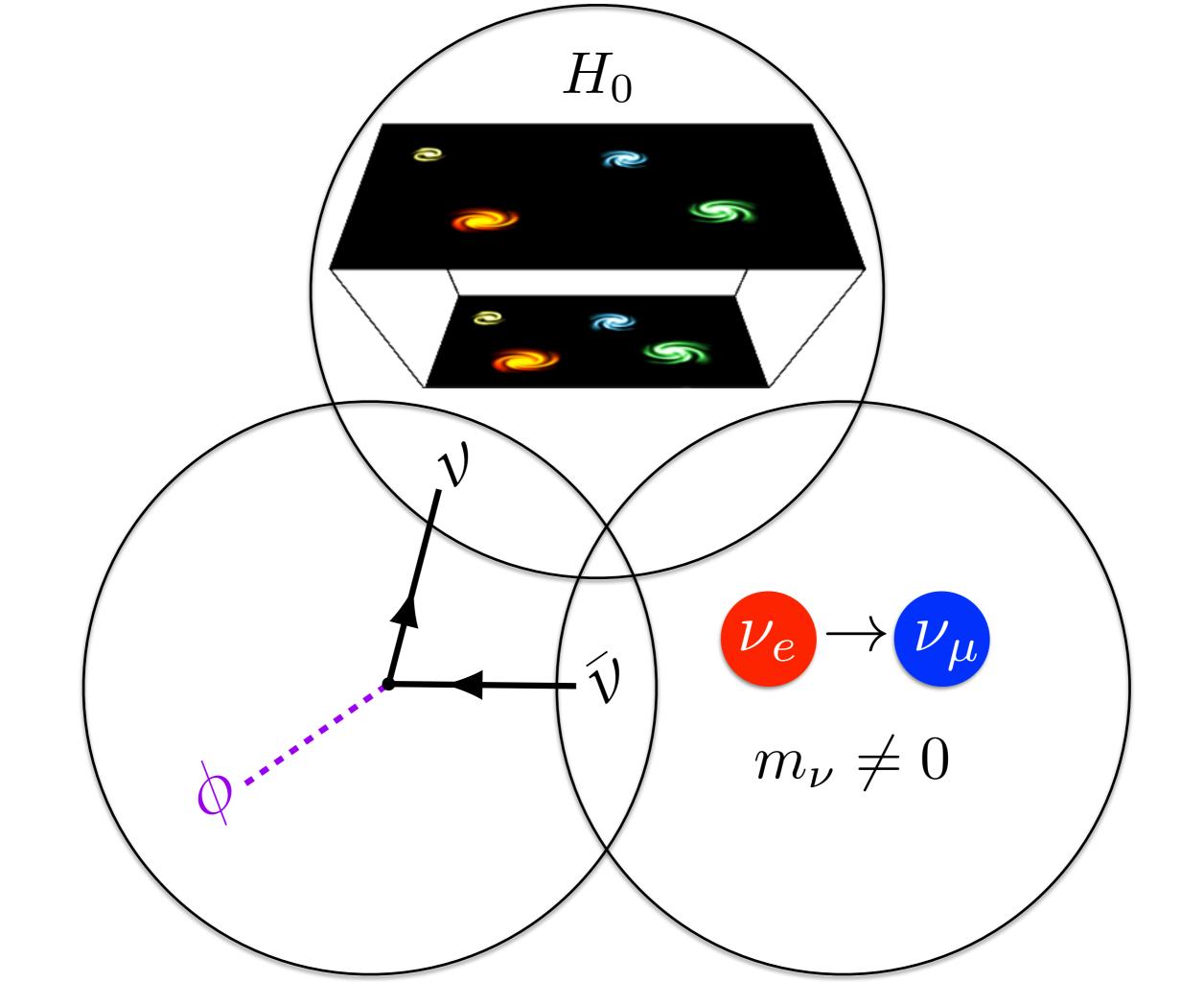
$$H_0 = 73.2 \pm 1.3 \,\text{km/s/Mpc}$$

**ACDM Prediction** 

Planck 2018 1807.06209

$$H_0 = 67.4 \pm 0.5 \,\text{km/s/Mpc}$$

4.2 σ tension within ΛCDM!



#### Outline

1) The Hubble Tension

**Observational Evidence** 

**Brief review of Models** 

2) The Majoron as a solution to the H<sub>0</sub> tension

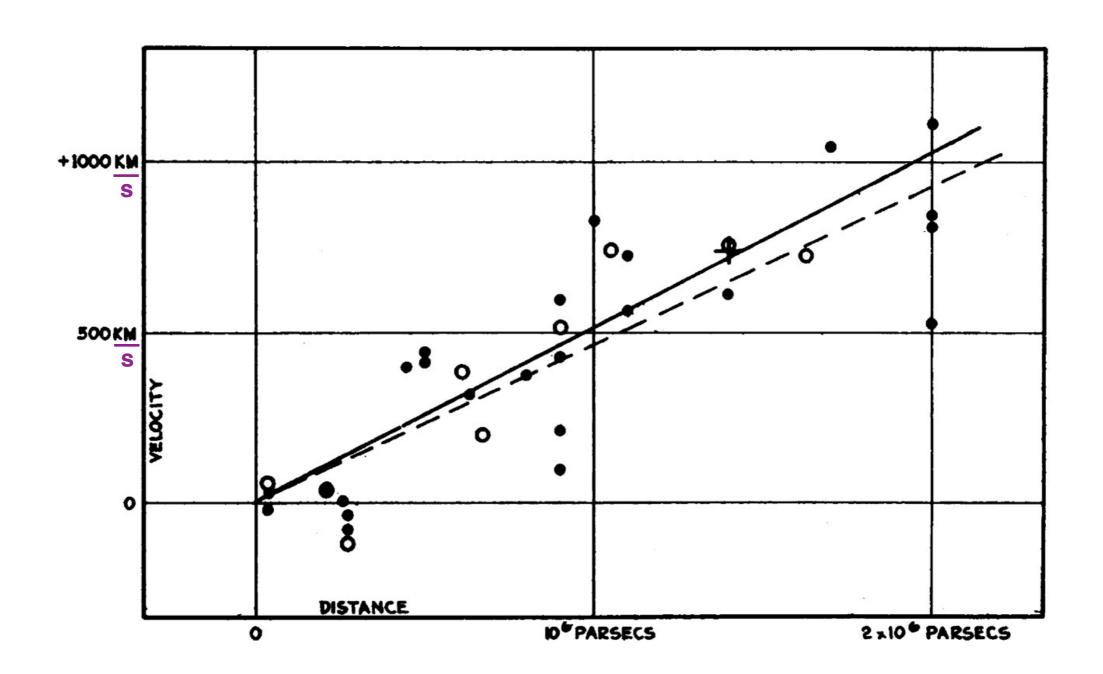
The singlet Majoron model Majoron Cosmology

3) Conclusions and Outlook

### The Hubble Law

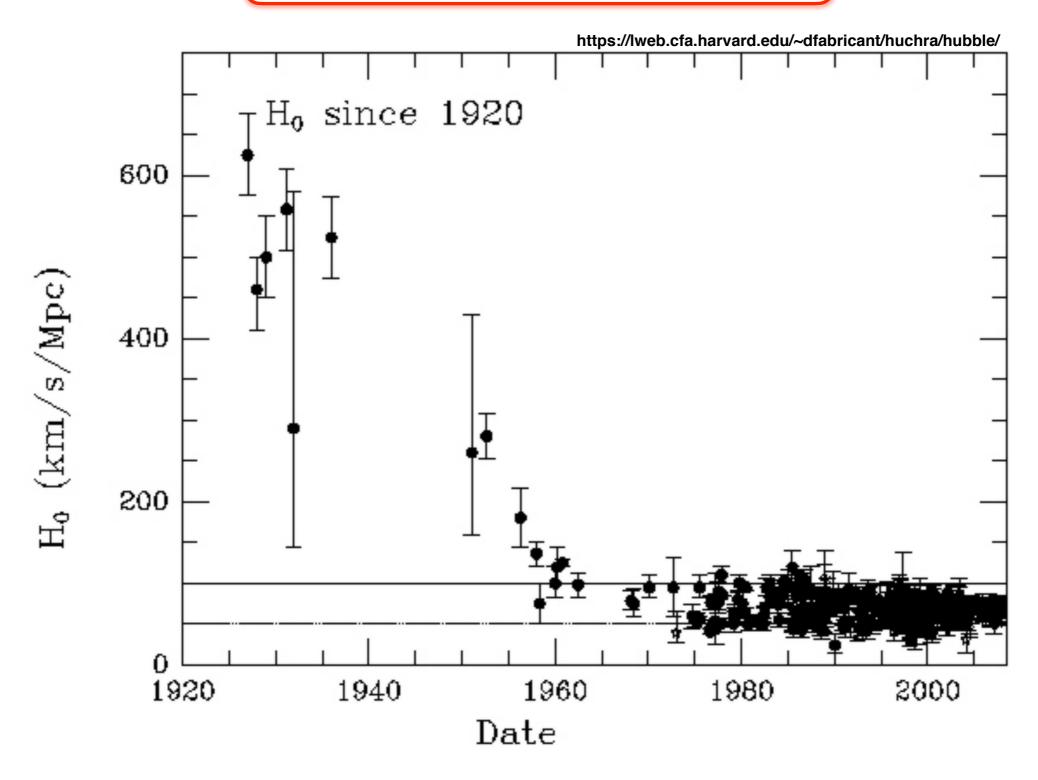
The Universe is expanding!

Hubble (1929): 
$$v = H_0 d$$

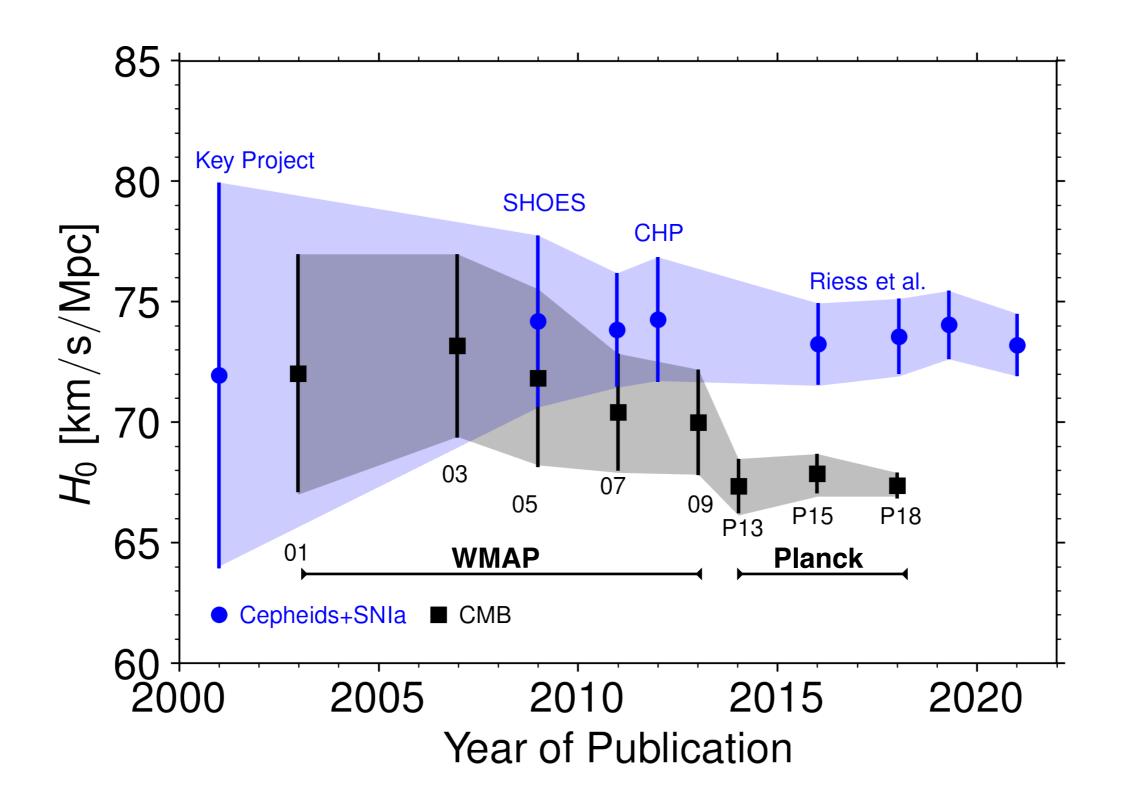


# The Hubble Tension in Perspective

Hubble law (1929):  $v = H_0 d$ 



# The Hubble Tension in Perspective



### The Hubble Tension

#### The Hubble Tension:

$$H_0 = 73.2 \pm 1.3 \,\text{km/s/Mpc}$$

$$H_0 = 67.4 \pm 0.5 \,\text{km/s/Mpc}$$

Riess et al. 2012.08534

Planck 2018 1807.06209

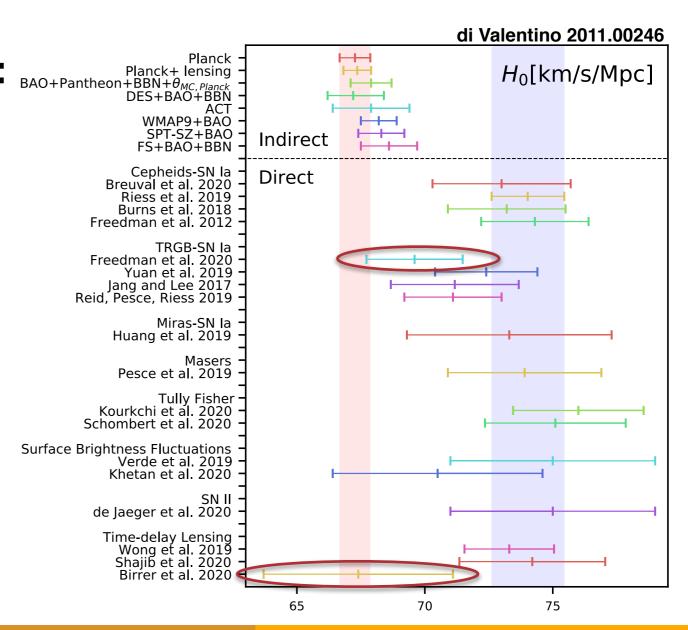
#### 4.2 $\sigma$ tension within $\Lambda$ CDM!

### A pattern has clearly emerged: Planck Planck+ lensing Planck+ lensing Planck+ Planck Planck

4-6 σ tension depending upon the datasets included

see Verde, Treu, Riess 1907.10625 for a review

- Baryon Acoustic Oscillations point to small H<sub>0</sub>
- Cepheids+Type-Ia SN are among the most precise and they point to  $H_0 \sim (73 \pm 2) \, \mathrm{km/s/Mpc}$
- Some analyses do point to smaller values, Freedman et al. 20' and Birrer et al. 20'



### The Hubble Tension

Possible resolutions:

1) Systematics in the CMB data

very unlikely

2) Systematics in local measurements

none so far\*

3) New feature of ΛCDM

Possibilities beyond ΛCDM:

See 2103.01183 by di Valentino et al. for a review (over 1000 references ...)

1) Late Universe Modifications

very unlikely

2) Early Universe Modifications

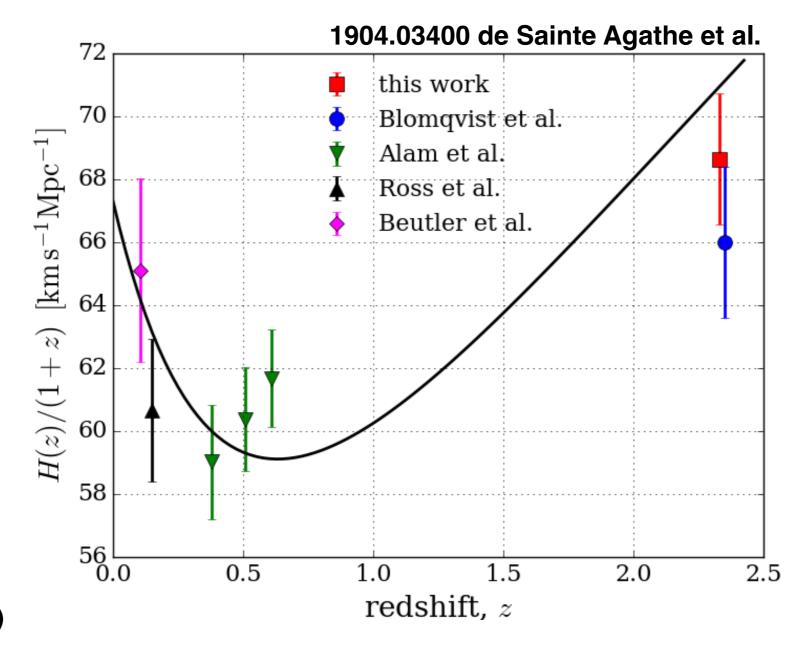
hard but doable

\*After 8 years of intense scrutiny, it has been pointed out recently that perhaps dust could play a relevant role Mortsell et al. 2105.11461

#### Why late Universe modifications do not work? see e

see e.g. 2103.08723 by Efstathiou

Because type Ia SN and Baryon Acoustic Oscillations constrain the expansion history of the Universe at z < 2.5 and they agree with the predictions of  $\Lambda$ CDM



(H₀ is measured locally, at z < 0.15)

Why Early Universe modifications could work?

Because the CMB does not measure  $H_0$  directly!

Planck measures the positions of the peaks:  $\theta_s \equiv r_s/D_M(z_\star)$ (0.03% precision)

$$r_s = \int_{z_\star}^\infty \frac{c_s}{H(z')} \, dz' \qquad \text{Comoving sound horizon}$$
 (Early Universe) 
$$D_M(z) = \int_0^z \frac{1}{H(z')} \, dz' \qquad \text{Comoving angular diameter distance}$$
 (Late Universe) 
$$H_0$$

**Model Building task:** 

The game is to make  $r_s$  smaller by ~8% so that  $H_0$ can be the one reported by Riess. But, not spoiling the fit to ultra precise CMB data from Planck!

simplest:

**Knox and Millea** 1908.03663

Enhance the expansion history of the Universe prior and close to recombination!

### Dark Radiation as a solution?



By far the simplest possibility, we have  $\mathcal{O}(10^3)$  models that can do it:

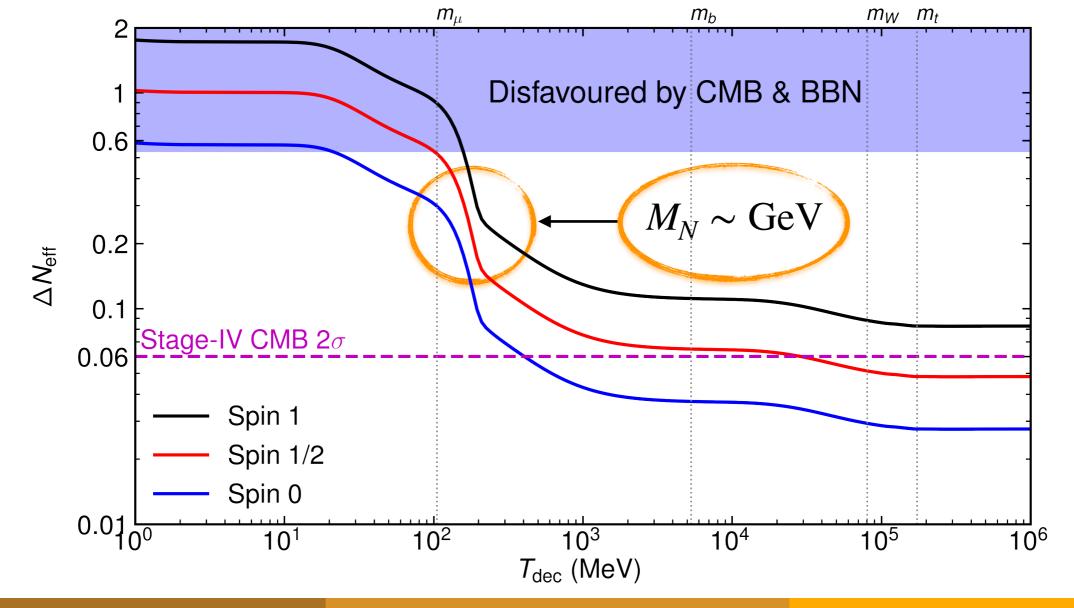
$$\Delta N_{\rm eff} = N_{\rm eff} - N_{\rm eff}^{\rm SM} > 0$$
  $N_{\rm eff}^{\rm SM} \simeq 3.04$ 

Goldstone Bosons as Fractional Cosmic Neutrinos

Steven Weinberg Phys. Rev. Lett. **110**, 241301 – Published 10 June 2013

Editors' Suggestion

#### Typically interpreted as additional massless dark radiation as a relic from the Big Bang



### Neff as a solution to the Ho Tension?

- How large would  $\Delta N_{
m eff}$  need to be to solve the tension?

$$H_0 \simeq \left[67.4 + 6.2 \,\Delta N_{\text{eff}}\right] \,\text{km/s/Mpc}$$

Vagnozzi 1907.07569

- $\Delta N_{
  m eff} \simeq 1$  would yield the value of  $H_0$  reported by Riess
- - Constraints are dominated by Helium measurements (that could suffer from systematics)
  - In many models  $\Delta N_{\mathrm{eff}}^{\mathrm{CMB}} \neq \Delta N_{\mathrm{eff}}^{\mathrm{BBN}}$
- $_{\odot}$  Problem 2) Within the framework of  $\Lambda$ CDM Planck is compatible with  $N_{
  m eff} \simeq 3$

$$N_{\rm eff}^{\rm CMB+BAO} = 2.99 \pm 0.17$$
 Planck 2018

Maybe there is an effect that can compensate a large Neff at the level of the CMB fit?

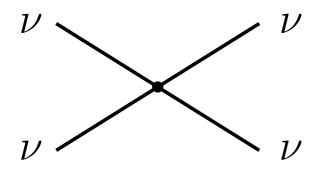
Weutrinos interactions can lead to a relevant impact on the CMB spectra Bashinsky and Seljak astro-ph/0310198

$$\delta G_{\mu\nu} = 8\pi G \, \delta T_{\mu\nu}$$



### **Neutrinos and the Hubble Tension**

Strong Neutrino Scattering + Extra Radiation Kreisch, Cyr-Racine, Doré 1902.00543



H<sub>0</sub> tension solved if TEEE data is ignored



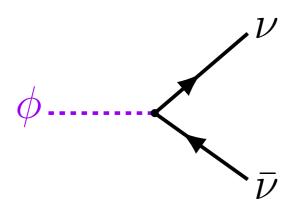
If pol data is included no solution for H<sub>0</sub>



Almost excluded by Lab data (Blinov++1905.02727)



Light Neutrinophilic Scalar + Dark Radiation Escudero & Witte 1909.04044



 $H_0$  tension from 4.4 $\sigma$  to 2.5 $\sigma$ 



CMB fit is not degraded

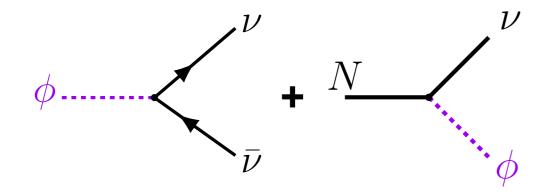


**Direct connection with Seesaw** 



Ad hoc  $\Delta N_{\rm eff} \sim 0.5$ 





Sterile neutrinos can source  $\Delta N_{eff}\sim 0.4$  Sterile neutrinos can lead to Leptogenesis Ho tension from 4.2 $\sigma$  to 2 $\sigma$ 

Hundreds of Models in the market See 2103.01183 by di Valentino et al.

Most of them do not actually work. They either lead to a bad CMB fit or do not shift H₀ enough Schöneberg et al. 2107.10291: *The H₀ Olympics: A fair ranking of proposed models* 

Model	$\Delta N_{ m param}$	$M_B$	Gaussian Tension	$Q_{\rm DMAP}$ Tension		$\Delta\chi^2$	$\Delta { m AIC}$		Finalist
$\Lambda { m CDM}$	0	$-19.416 \pm 0.012$	$4.4\sigma$	$4.5\sigma$	X	0.00	0.00	X	X
			<b>1</b>			•	Ī		

How large is the Hubble tension? small values here are better!

How good is the CMB fit? negative values are good here!

Hundreds of Models in the market See 2103.01183 by di Valentino et al.

Most of them do not actually work. They either lead to a bad CMB fit or do not shift H₀ enough Schöneberg et al. 2107.10291: *The H₀ Olympics: A fair ranking of proposed models* 

	Model	$\Delta N_{ m param}$	$M_B$	Gaussian Tension	$Q_{\mathrm{DMAP}}$ Tension		$\Delta\chi^2$	$\Delta { m AIC}$		Finalist
	$\Lambda \mathrm{CDM}$	0	$-19.416 \pm 0.012$	$4.4\sigma$	$4.5\sigma$	X	0.00	0.00	X	X
	$\Delta N_{ m ur}$	1	$-19.395 \pm 0.019$	$3.6\sigma$	$3.9\sigma$	X	-4.60	-2.60	X	X
	SIDR	1	$-19.385 \pm 0.024$	$3.2\sigma$	$3.6\sigma$	X	-3.77	-1.77	X	X
	DR-DM	2	$-19.413 \pm 0.036$	$3.3\sigma$	$3.4\sigma$	X	-7.82	-3.82	X	X
	mixed DR	2	$-19.388 \pm 0.026$	$3.2\sigma$	$3.7\sigma$	X	-6.40	-2.40	X	X
	$SI\nu+DR$	3	$-19.440 \pm 0.038$	3.70	$3.9\sigma$	X	-3.56	2.44	X	X
2	Majoron	3	$-19.380 \pm 0.027$	$3.0\sigma$	$2.9\sigma$	<b>√</b>	-13.74	-7.74	$\checkmark$	√ ②
3	primordial B	1	$-19.390 \pm 0.018$	$3.5\sigma$	$3.5\sigma$	X	-10.83	-8.83	$\checkmark$	√ ③
5	varying $m_e$	1	$-19.391 \pm 0.034$	$2.9\sigma$	$3.2\sigma$	X	-9.87	-7.87	$\checkmark$	√ ③
5	varying $m_e + \Omega_k$	2	$-19.368 \pm 0.048$	$2.0\sigma$	$1.7\sigma$		-16.11	-12.11	$\checkmark$	✓ •
۲	EDE	3	$-19.390 \pm 0.016$	$3.6\sigma$	$1.6\sigma$	/ <	-20.80	-14.80	$\checkmark$	✓ ②
	NEDE	3	$-19.380 \pm 0.021$	$3.2\sigma$	$2.0\sigma$	$\checkmark$	-17.70	-11.70	$\checkmark$	✓ ②
2	CPL	2	$-19.400 \pm 0.016$	$3.9\sigma$	$4.1\sigma$	X	-4.23	-0.23	X	X
	PEDE	0	$-19.349 \pm 0.013$	$2.7\sigma$	$2.0\sigma$	$\checkmark$	4.76	4.76	X	X
	MPEDE	1	$-19.400 \pm 0.022$	$3.6\sigma$	$4.0\sigma$	X	-2.21	-0.21	X	X
	$\mathrm{DM} \to \mathrm{DR} + \mathrm{WDM}$	2	$-19.410 \pm 0.013$	$4.2\sigma$	$4.4\sigma$	X	-4.18	-0.18	X	X
	$DM \to DR$	2	$-19.410 \pm 0.011$	$4.3\sigma$	$4.2\sigma$	X	0.11	4.11	X	X

- None of them fully solves the Hubble tension!
- Most of those that can ameliorate the tension are not theoretically well motivated

best performance:

#### A critical review of the best performing models

Schöneberg et al. 2107.10291: The H₀ Olympics: A fair ranking of proposed models

Model	$\Delta N_{ m param}$	$M_B$	Gaussian Tension	$Q_{ m DMAP}$ Tension		$\Delta \chi^2$	$\Delta { m AIC}$		Finalist
Majoron	3	$-19.380 \pm 0.027$	$3.0\sigma$	$2.9\sigma$	<b>√</b>	-13.74	-7.74	<b>√</b>	√ 2
primordial B	1	$-19.390 \pm 0.018$	$3.5\sigma$	$3.5\sigma$	X	-10.83	-8.83	$\checkmark$	√ ③
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varying $m_e + \Omega_k$	2	$-19.368 \pm 0.048$	$2.0\sigma$	$1.7\sigma$	$\checkmark$	-16.11	-12.11	$\checkmark$	✓ •
EDE	3	$-19.390 \pm 0.016$	$3.6\sigma$	$1.6\sigma$	$\checkmark$	-20.80	-14.80	$\checkmark$	✓ ②
NEDE	3	$-19.380 \pm 0.021$	$3.2\sigma$	$2.0\sigma$	$\checkmark$	-17.70	-11.70	$\checkmark$	√ ②

lacksquare Primordial magnetic fields &  $m_e(t)+\Omega_k$ 

The idea here is that recombination happens earlier than in ΛCDM by either

- a) using primordial magnetic fields of ~ 1 nGauss on kpc scales [Jedamzik & Pogosian 2004.09487]
- b) enhancing  $m_e(t)$  at recombination by ~ 2% [Hart & Chluba 1912.03986]
- Good exercises, but not much theoretical motivation for either of the two settings ...
- Early Dark Energy Poulin, Smith, Karwal, Kamionkowski 1811.04083 Agrawal, Cyr-Racine, Pinner, Randall 1904.01016

The idea is that there is an early dark energy component just acting right before recombination

This can be done with a very light scalar field with  $m_{\phi} \sim 10^{-27}\,\mathrm{eV}$  and  $f \lesssim M_{\mathrm{Pl}}$  that yields  $f_{\mathrm{EDE}} \sim 10\,\%$  but with a very particular potential:  $V_{\phi} \sim m^2 f^2 \left[1 - \cos\phi/f\right]^3 \sim m^2 \phi^6/f^4$ 

- Highly unclear where such potential could come from and there is a coincidence problem ...
- New Early Dark Energy

Another possibility is to trigger a 1st order phase transition at  $T \sim \mathrm{eV}$  [Niedermann & Sloth 1910.10739]

 $ule{e}$  This seems theoretically more plausible, although requires some idea for why  $T\sim {
m eV}$  but  $m_\phi\sim 10^{-27}\,{
m eV}$ 

### The Hubble Tension

#### 1) Observational evidence

There is strong observational evidence from Cepheids+SNIa

However, it is still just a tension. It needs to be confirmed by other methods

We expect significant improvements in ~3-4 years, particularly with upcoming data from Gaia & the James Webb telescope

#### 2) Theoretical modeling

Despite the strong efforts, we have no perfect model so far

Most of the models lack theoretical motivation

#### My personal view as a particle physicist:

Exciting opportunity to try to learn about fundamental physics

# Mid Seminar Pause =)

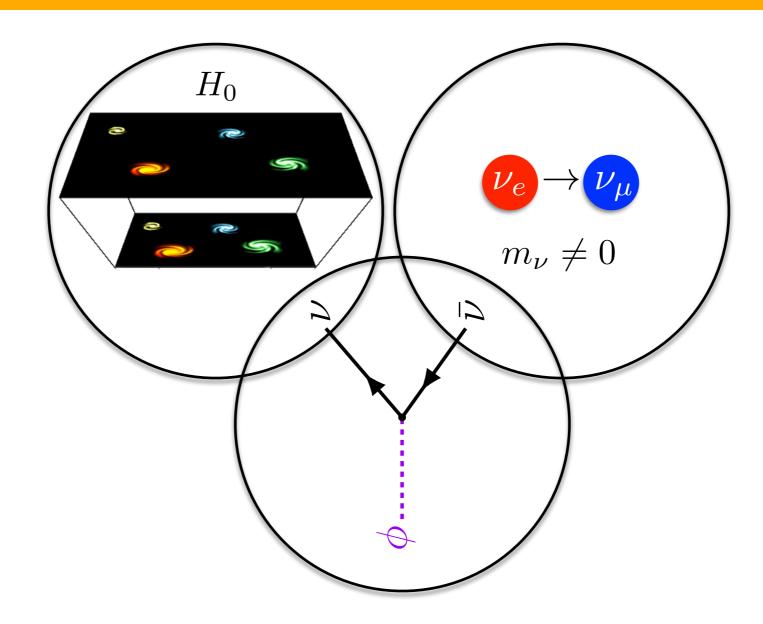
Questions?

Comments?

**Criticism?** 

All are most welcome!

### The Idea



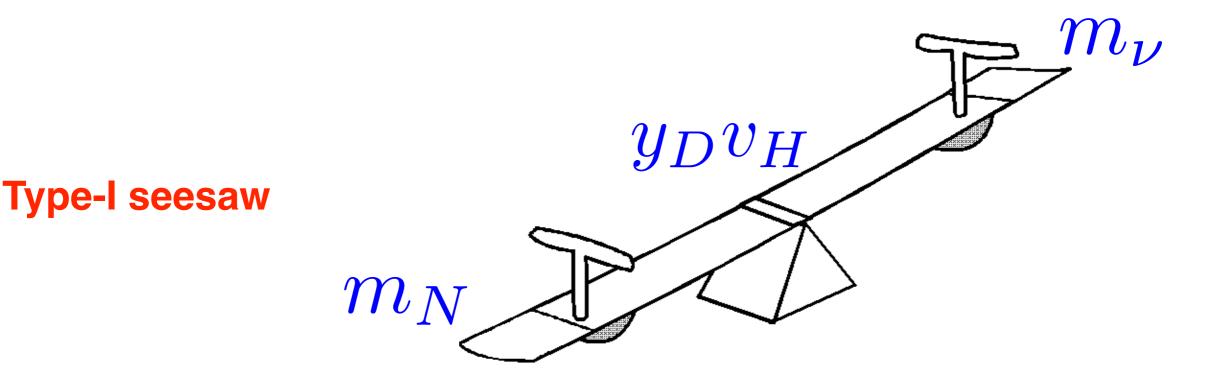
What is the Majoron?

What are its properties and its cosmology?

Why can it address the Hubble tension?

#### The Seesaw Mechanism

Minkowski, Yanagida, Gell-Mann, Ramond, Slansky, Glashow, Mohapatra, Senjanovic, Schechter, Valle



#### Neutrinos are very light Majorana particles:

$$m_{\nu} \simeq 0.03 \,\text{eV} \, \left(\frac{y_D}{10^{-6}}\right)^2 \, \frac{\text{TeV}}{M_N}$$

#### The Scenario

#### Spontaneously Broken Symmetry Global U(1)<sub>L</sub>

Chikashige, Mohapatra, Peccei (1981)

Sterile Neutrinos 
$$\mathcal{L} = - \; \frac{\lambda_{N_{ij}}}{\sqrt{2}} \Phi \; \overline{N}_{R, \, i} N_{R, \, j}^c - h_{\alpha i} \overline{L}_L^{\alpha} H N_{Ri} + \mathrm{h.c.} \; , \quad \begin{array}{l} L[\Phi] = 2 \\ L[N] = 1 \end{array}$$

SSB: 
$$\Phi \rightarrow v_L/\sqrt{2}$$
  $\longrightarrow$   $M_N = \lambda_N v_L$  seesaw  $m_\nu \simeq h^2 v_H^2/(2M_N)$ 

$$V_{\Phi} = -\mu_{\Phi}^2 \Phi^{\dagger} \Phi + \lambda_{\Phi} (\Phi^{\dagger} \Phi)^2 - \lambda_{\Phi H} (H^{\dagger} H) (\Phi^{\dagger} \Phi)$$

$$\Phi = \frac{v_L + \rho}{\sqrt{2}} e^{i\phi/v_L} \qquad \rho \equiv \text{CP-even scalar} \qquad m_\rho^2 = 2\lambda_\Phi v_L^2$$

$$\phi \equiv \text{Majoron} \qquad \text{pseudo-Gold}$$

$$\rho \equiv \text{CP-even scalar}$$

$$m_{\rho}^2 = 2\lambda_{\Phi} v_L^2$$

$$\phi \equiv \text{Majoron}$$

pseudo-Goldstone:  $m_{\phi} \simeq 0$ 

$$\mathcal{L}_{\text{eff}} = -\frac{\lambda_{N}}{2} \left[ \rho \bar{N} N - i \phi \bar{N} \gamma_{5} N \right] \qquad \lambda_{\nu} \ll \lambda_{N\nu} \ll \lambda_{N}$$

$$-\frac{\lambda_{N\nu}}{2} \left[ \rho \bar{N} \nu - i \phi \bar{N} \gamma_{5} \nu \right] + \text{h.c.} \qquad \lambda_{\nu} \simeq |\theta| \lambda_{N\nu} \simeq |\theta|^{2} \lambda_{N}$$

$$-\frac{\lambda_{\nu}}{2} \left[ \rho \bar{\nu} \nu - i \phi \bar{\nu} \gamma_{5} \nu \right] \qquad |\theta|^{2} \simeq 5 \times 10^{-11} \frac{m_{\nu}}{0.05 \text{ eV}} \frac{1 \text{ G}}{M}$$

$$\lambda_{\nu} \ll \lambda_{N\nu} \ll \lambda_{N}$$

$$\lambda_{\nu} \simeq |\theta| \lambda_{N\nu} \simeq |\theta|^2 \lambda_N$$

$$|\theta|^2 \simeq 5 \times 10^{-11} \frac{m_{\nu}}{0.05 \,\text{eV}} \frac{1 \,\text{GeV}}{M_N}$$

### The Scenario

#### Spontaneously Broken Symmetry Global U(1)<sub>L</sub>

Chikashige, Mohapatra, Peccei (1981)

The Majoron:  $\phi$ 

$$\mathcal{L}_{\rm int} = i\lambda \,\phi \,\bar{\nu} \,\gamma_5 \,\nu$$

Very weakly interacting:

$$\lambda \simeq 10^{-13} rac{m_{
u}}{0.05\,\mathrm{eV}} rac{\mathrm{TeV}}{v_L}$$
 (seesaw)

Extremely feebly interacting with matter:  $\lambda_{\phi ee} \sim 10^{-20}$ 

$$\lambda_{\phi ee} \sim 10^{-20}$$

**Dimension-5 Planck suppressed operators:** 

$$m_{\phi} \sim v_L \sqrt{\frac{v_L}{M_{\rm Pl}}} \lesssim 0.1 \,\mathrm{keV}$$

$$\Delta V = \beta \left(\Phi^{\star}\Phi\right)^2 \frac{\Phi^{\star} + \Phi}{M_{\rm Pl}} \qquad \begin{array}{l} \text{Rothstein, Babu, Seckel hep-ph/9301213} \\ \text{Akhmedov, Berezhiani, Mohapatra, Senjanovic hep-ph/9209285} \end{array}$$

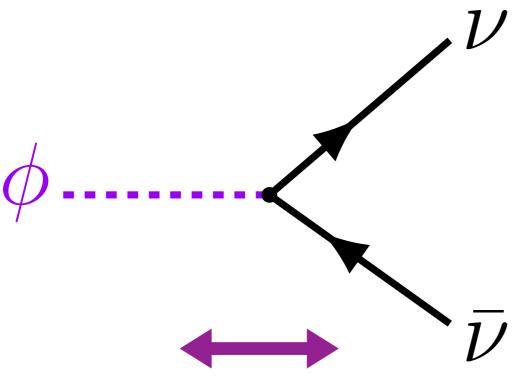
Key coincidence!!

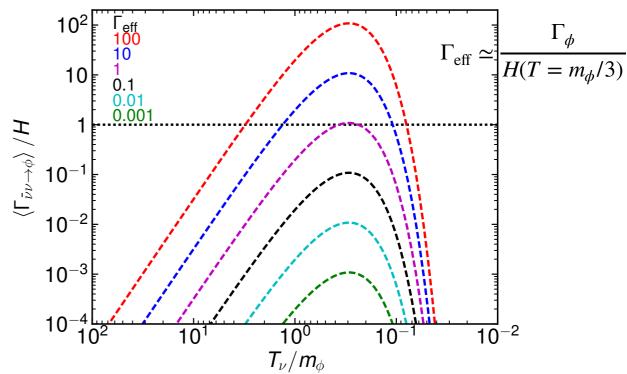
$$\begin{array}{c} \textbf{Retarmetep appareter space} & 10^{-15} \leqslant \lambda \\ 0.1 \, \mathrm{eV} < m_{\phi} \end{array} \leqslant \begin{array}{c} 10^{-3} & \tau_{\phi} \sim t_{\mathrm{rec}}/10 \\ 10^{-3} & \tau_{\phi} \sim t_{\mathrm{rec}}/10 \\ \mathrm{And\ assume\ that\ } n_{\phi} = 0 \text{ at\ BBN} \\ 0.1 \, \mathrm{eV} < m_{\phi} \end{array}$$

# Cosmological Implications

#### **Only Relevant Process:**







#### Two main effects:

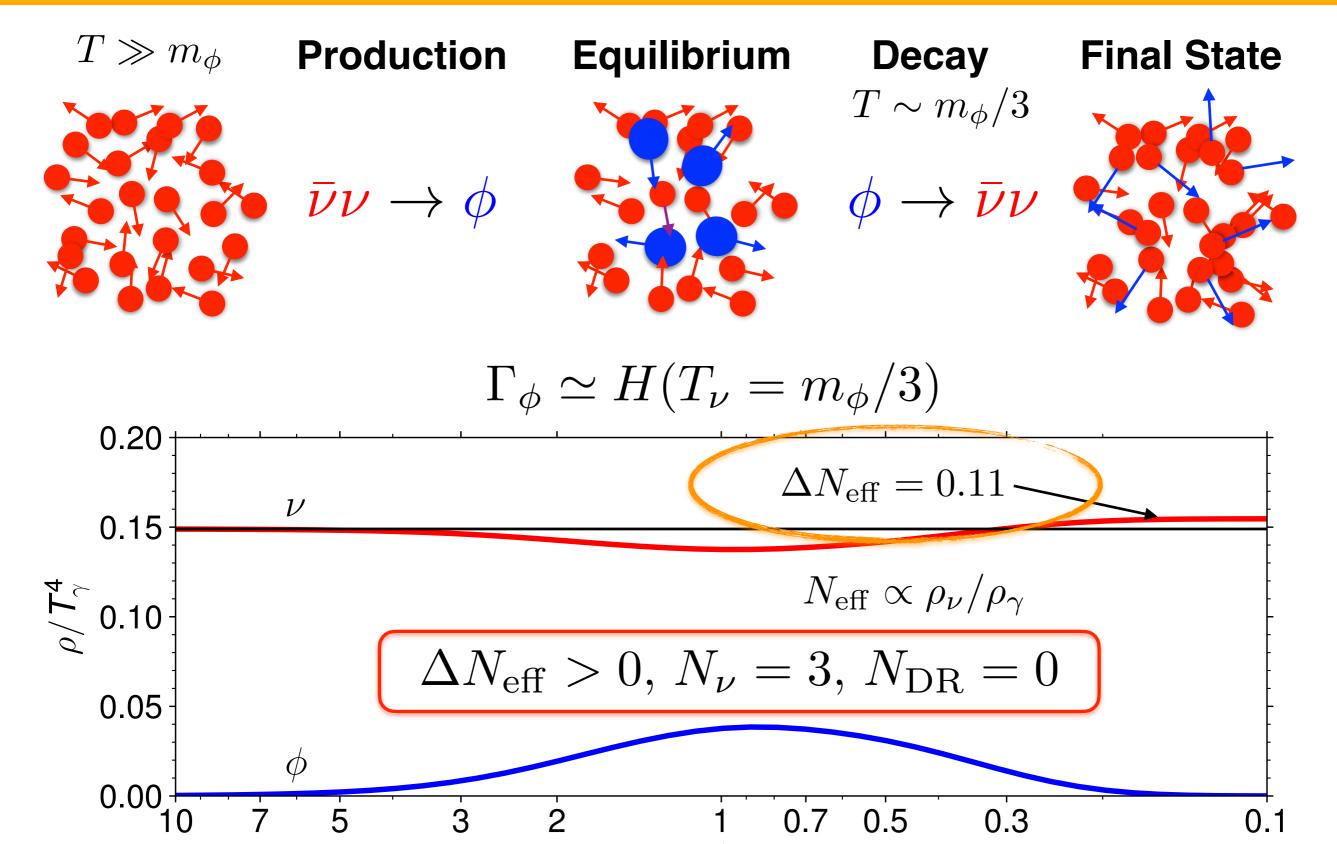
Chacko, Hall, Okui, Oliver hep-ph/0312267

Non-standard expansion history

**Erase the neutrino anisotropic stress** 

- We solve the Boltzmann equation for the background Escudero 1812.05605 & 2001.04466
- We include the neutrino-majoron Boltzmann hierarchy in CLASS

# Cosmological Implications



 $T_{\gamma}/m_{\phi}$ 

0.7

0.5

2

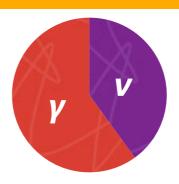
3

0.3

0.1

### **Neutrino Perturbations**

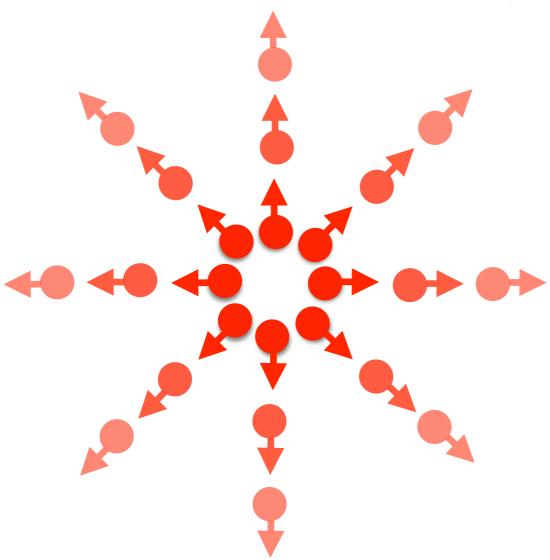
Neutrino perturbations are key:

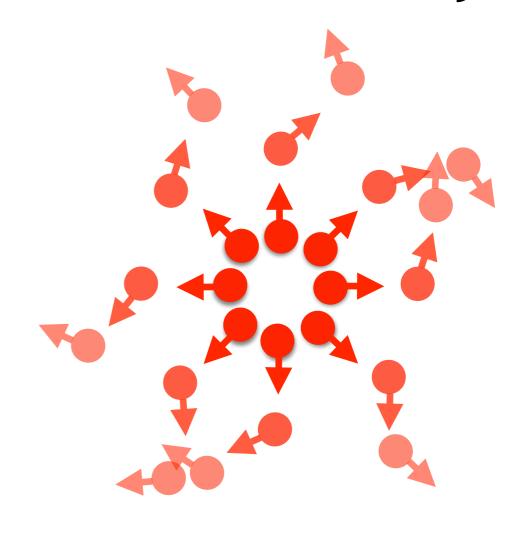


$$\delta G_{\mu\nu} = 8\pi G \delta T_{\mu\nu}$$

Free Streaming Neutrinos  $\sigma_{\!\nu} \neq 0$ 

Interacting Neutrinos  $\sigma_{\!\scriptscriptstyle \mathcal{V}} o 0$ 

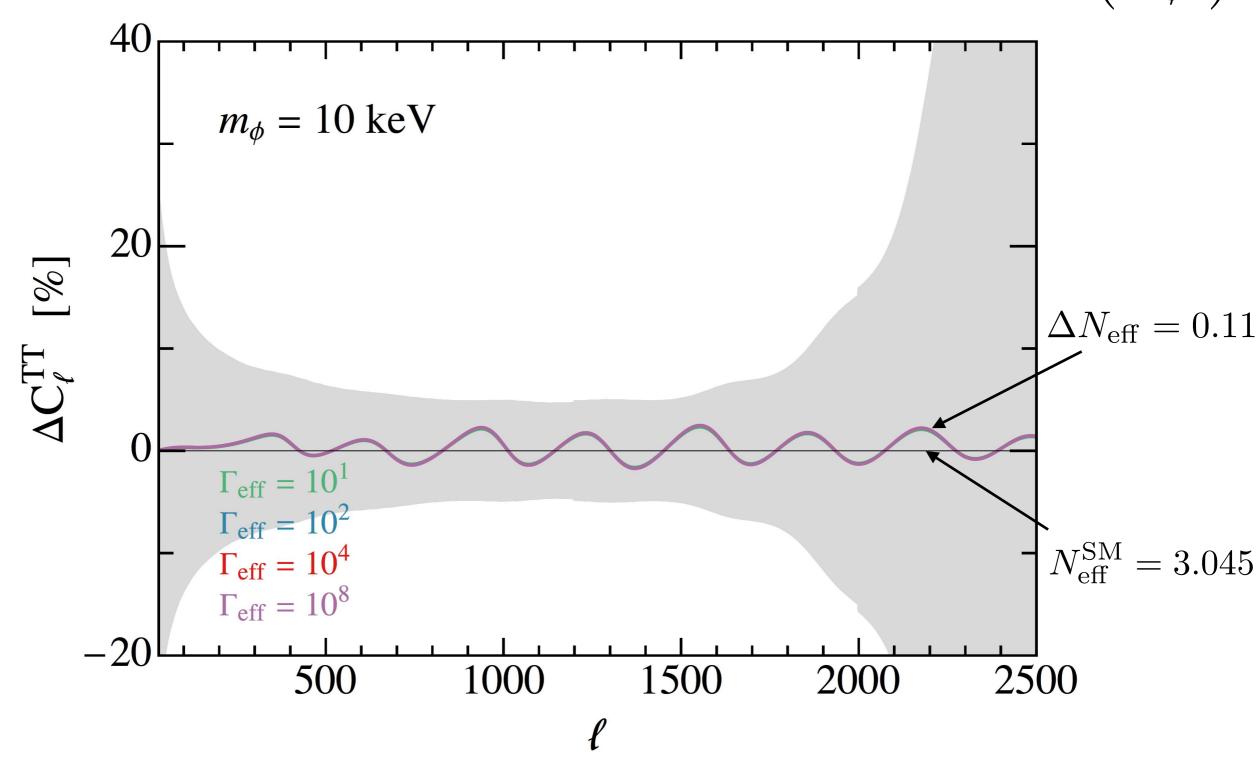




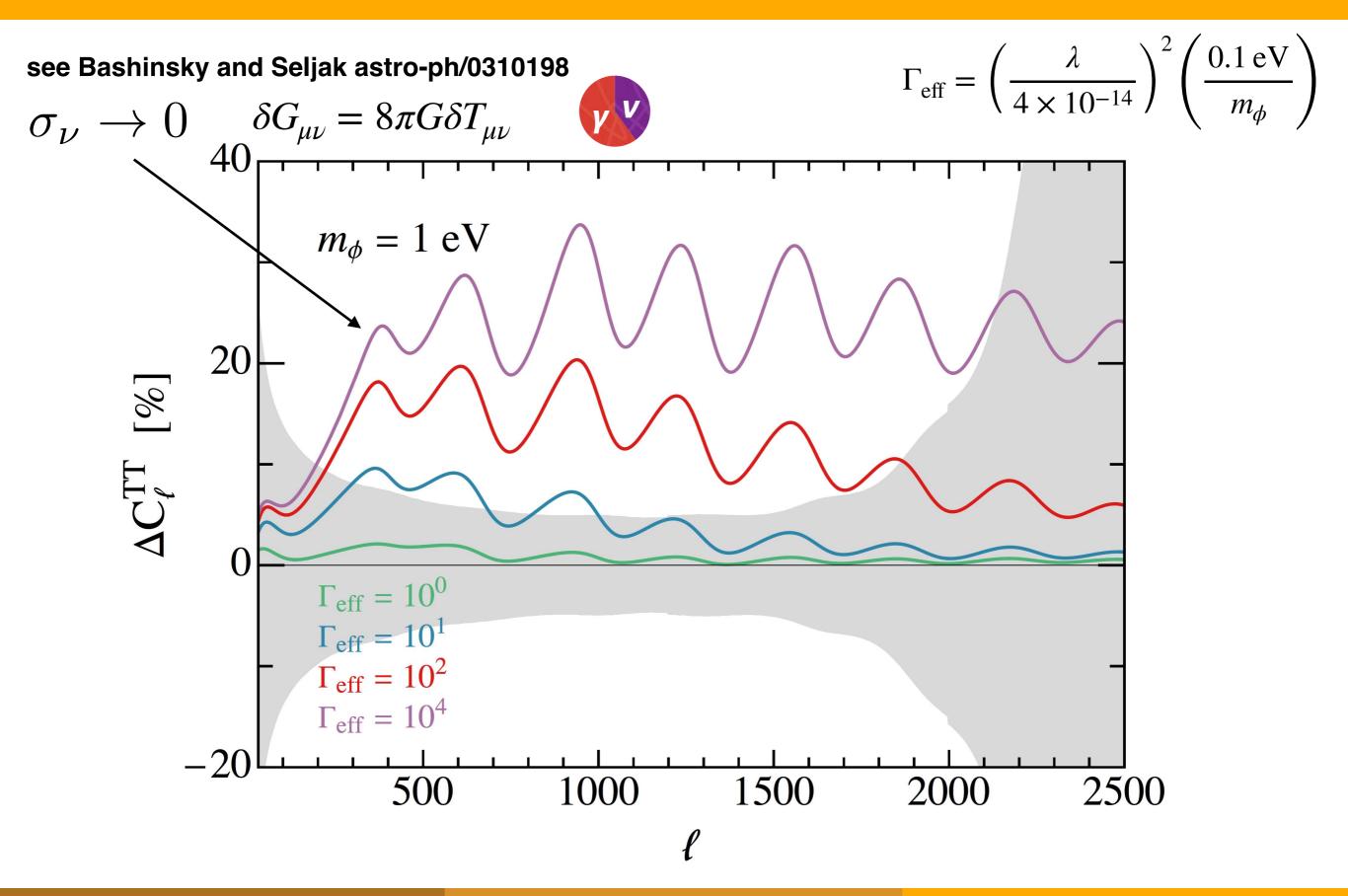
Effect on the CMB is to shift the positions of the peaks! Bashinsky and Seljak astro-ph/0310198

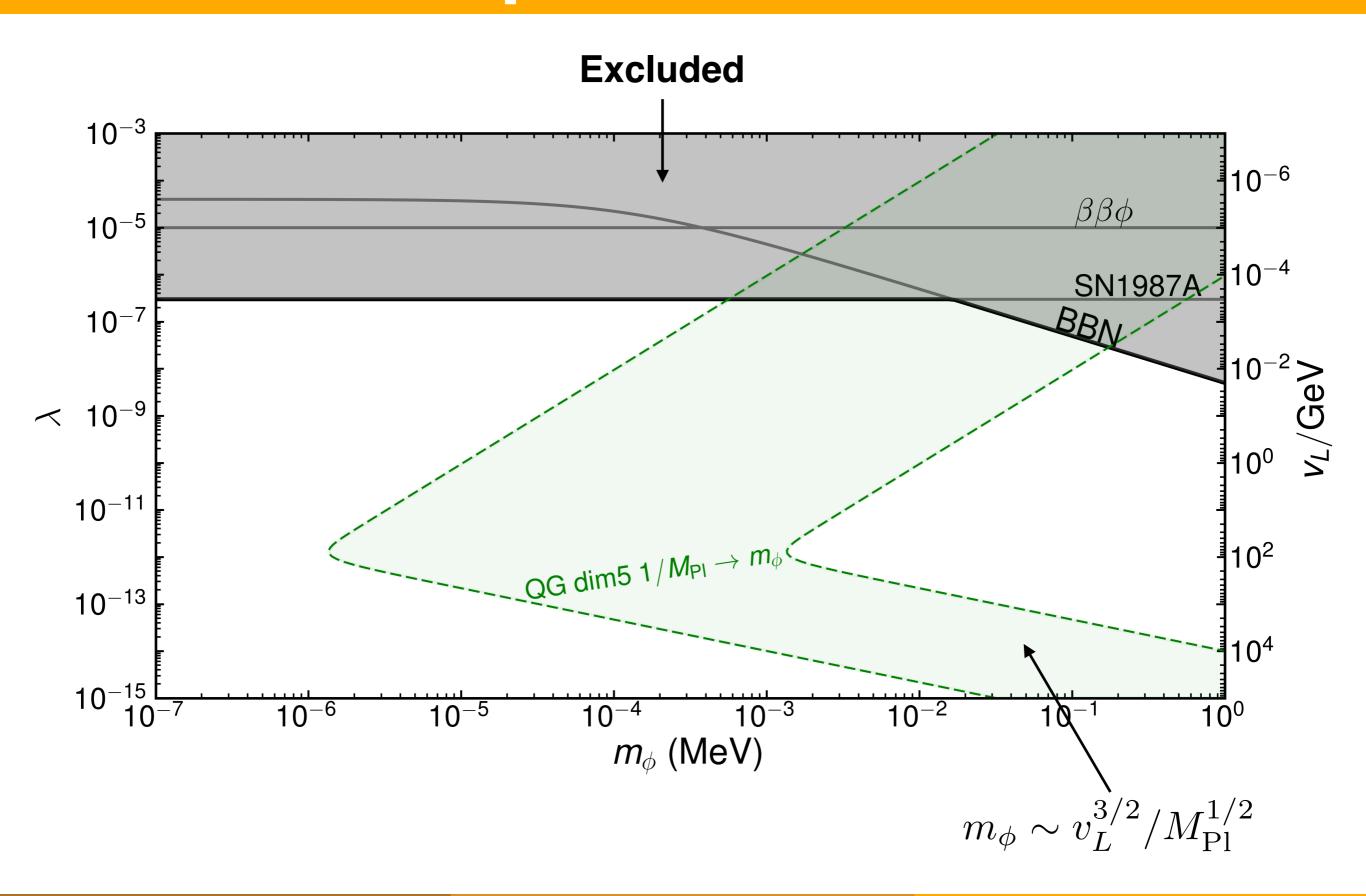
### Effects on the CMB

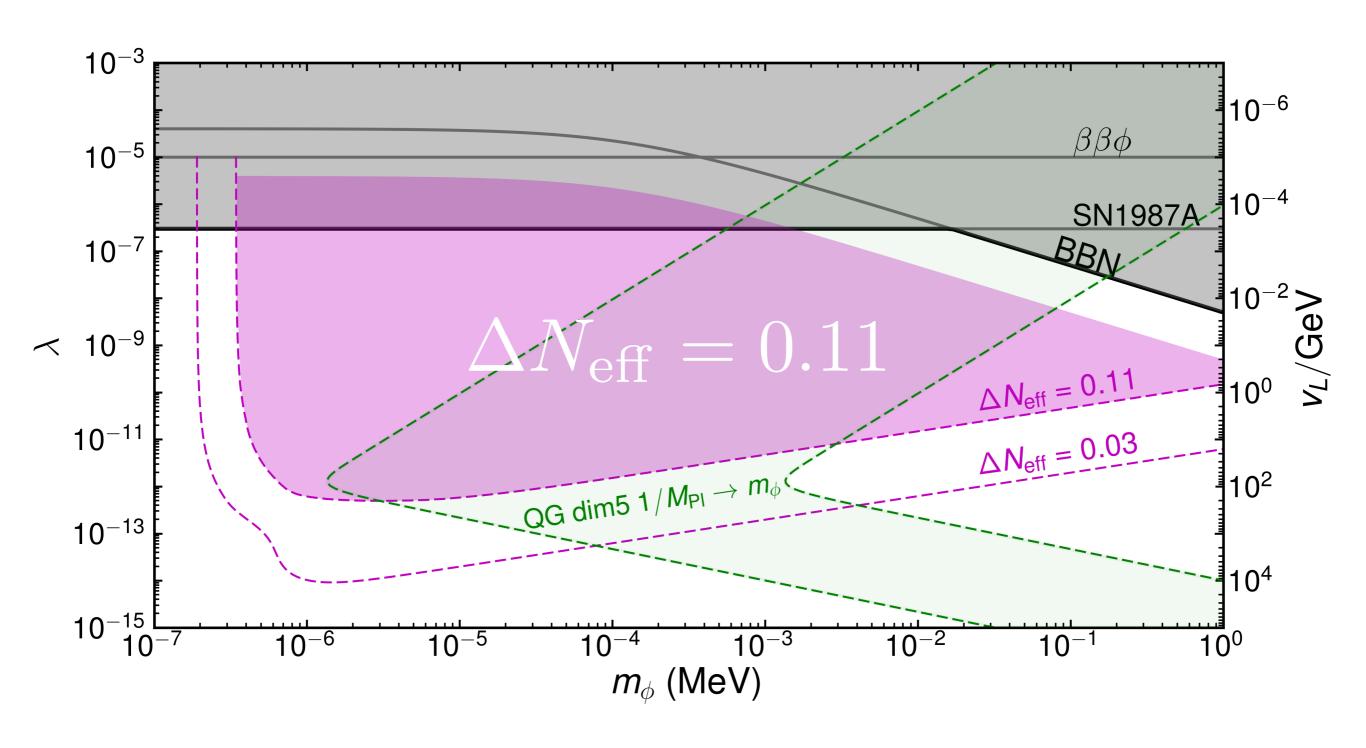
$$\Gamma_{\text{eff}} = \left(\frac{\lambda}{4 \times 10^{-14}}\right)^2 \left(\frac{0.1 \,\text{eV}}{m_{\phi}}\right)$$



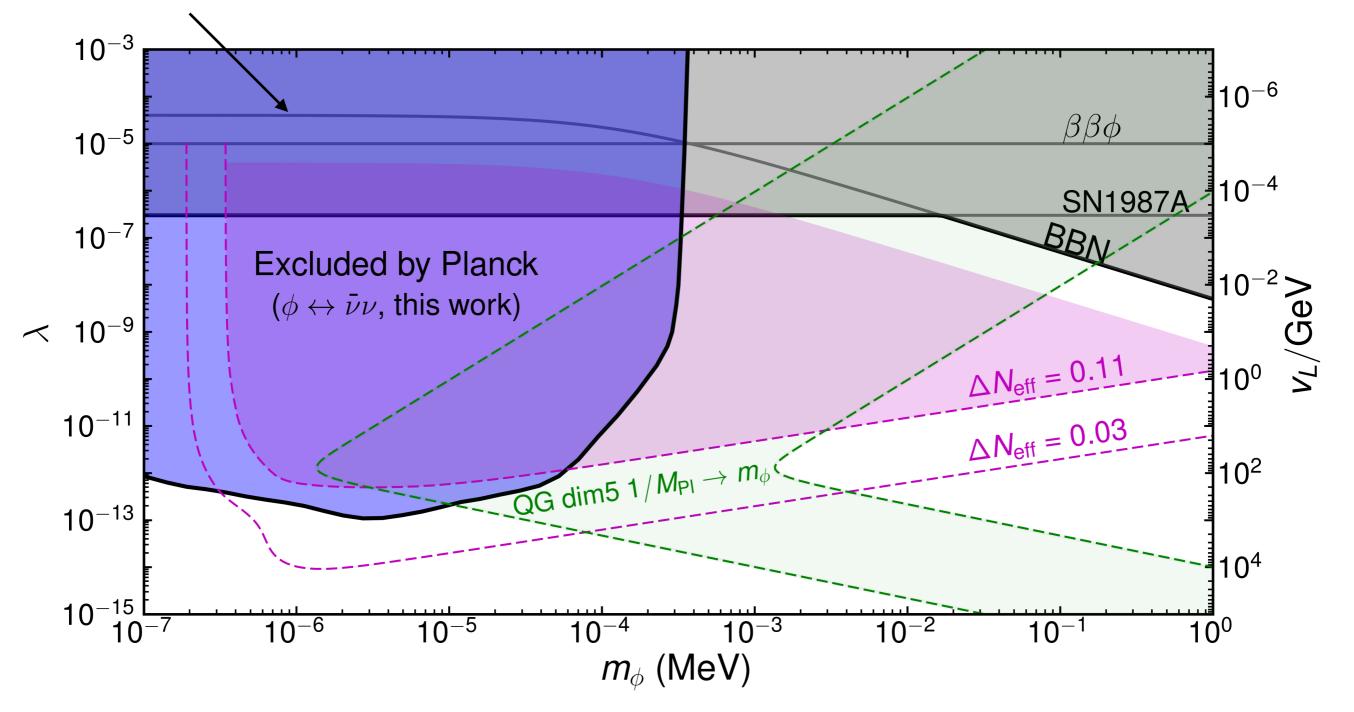
### **Effects on the CMB**

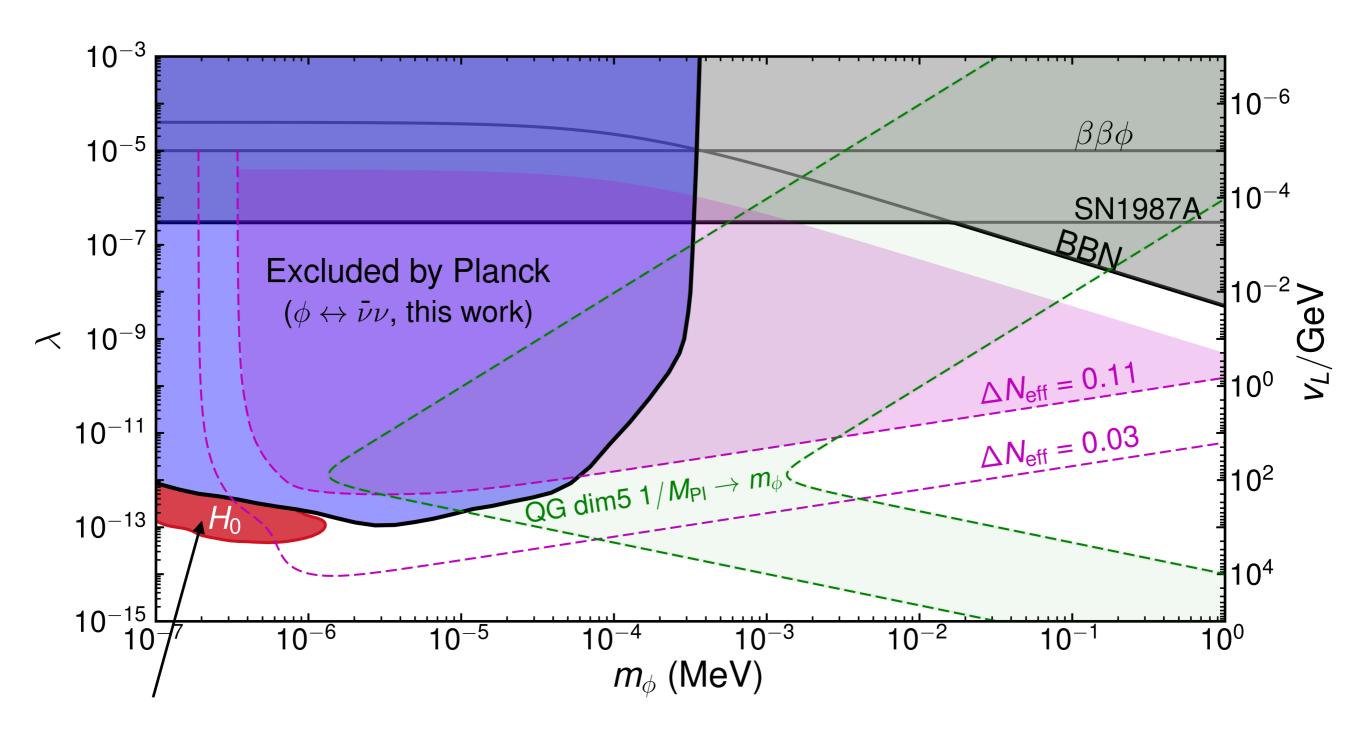






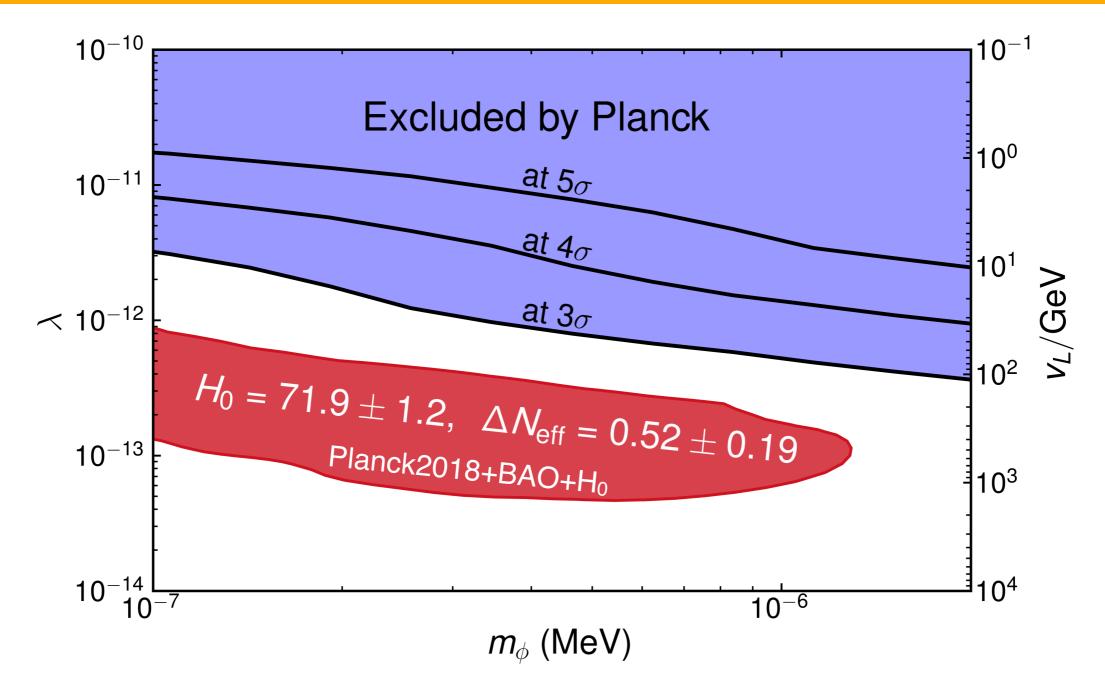
#### Full MCMC to Planck 2018 data





 $1\sigma$  preference when including  $H_0$  in the fit and an additional  $\Delta N_{eff}$ 

# Parameter Space for H<sub>0</sub>



- Requires a positive  $\Delta N_{eff} \sim 0.5$
- $H_0$  Thanks to the  $\nu \phi$  interactions Planck 2018 fit is not degraded wrt ΛCDM
  - lacktriangle Very close to the electroweak scale  $v_L \sim (0.1-1)\,\mathrm{TeV}$

# Summary of Escudero & Witte 19'

- The Majoron and the Hubble tension
  - Planck sets very stringent constraints
  - CMB S4 experiments will test large regions of parameter space since  $\sigma(N_{\rm eff}) \simeq 0.03$
  - Can significantly reduce the tension if:

$$m_{\phi} \sim (0.1 - 1) \,\mathrm{eV}$$
 $v_L \sim (0.1 - 1) \,\mathrm{TeV}$ 
 $\Delta N_{\mathrm{eff}} \sim 0.5$ 

- $\cong$   $\Delta N_{\rm eff} \sim 0.5$  is somewhat ad hoc
- Wow we have a very good reason for it!

# Primordial Majorons

There are sterile neutrinos in the model

since  $v_L \lesssim 1 \, {\rm TeV}$  then we can expect  $\, M_N \sim 1 \, {\rm GeV}$ 

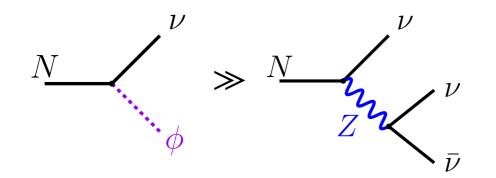
- The decays of GeV-scale sterile neutrinos in the early Universe can lead to  $\Delta N_{\rm eff}^{\rm BBN} \sim 0.3$  since  $T_d \sim M_N/10$
- Neutrino-Majoron interactions can rise it to  $\Delta N_{
  m eff}^{
  m CMB} \sim 0.6$

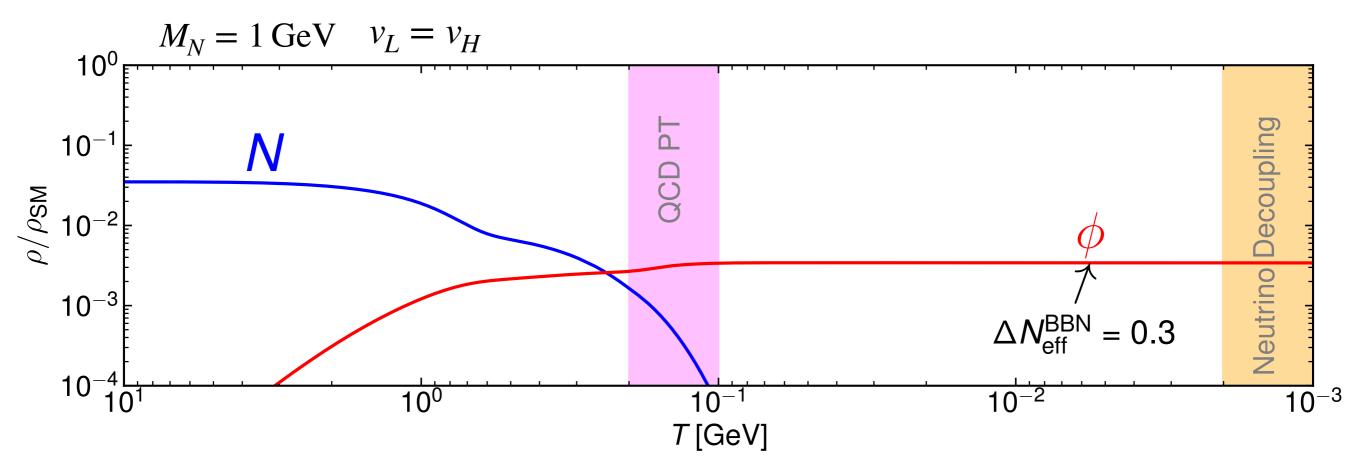
These sterile neutrinos can do ARS Leptogenesis!

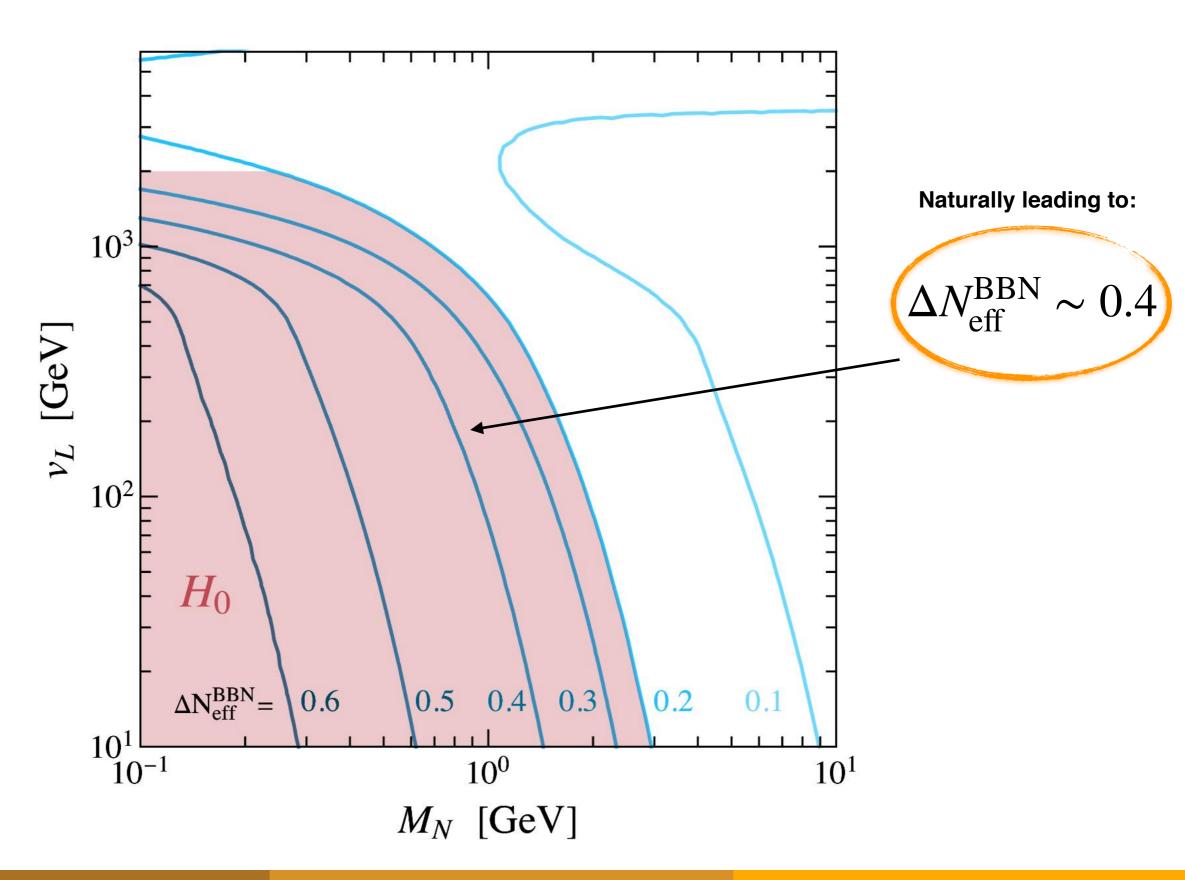
# Production of Majoron population

Sterile neutrinos have masses ~ GeV and interact with the majoron

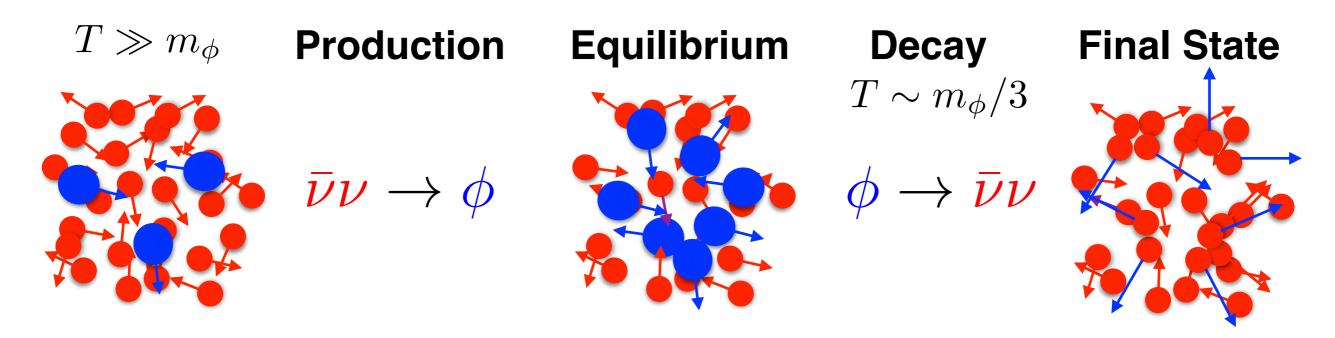
Sterile neutrinos that give mass to the active ones thermalize (Ghiglieri & Laine 1605.07720) In the majoron model sterile neutrinos have a new decay mode



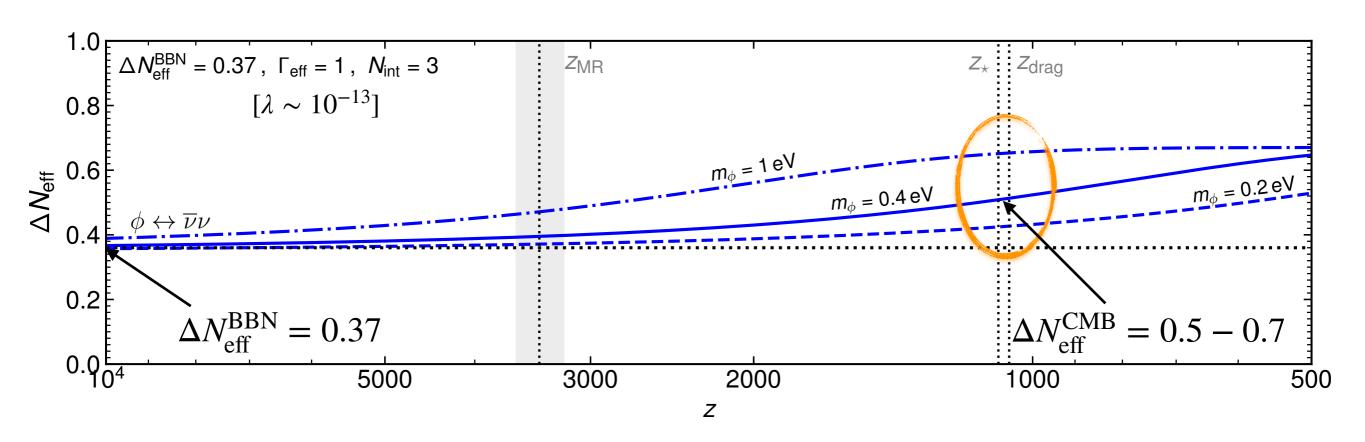




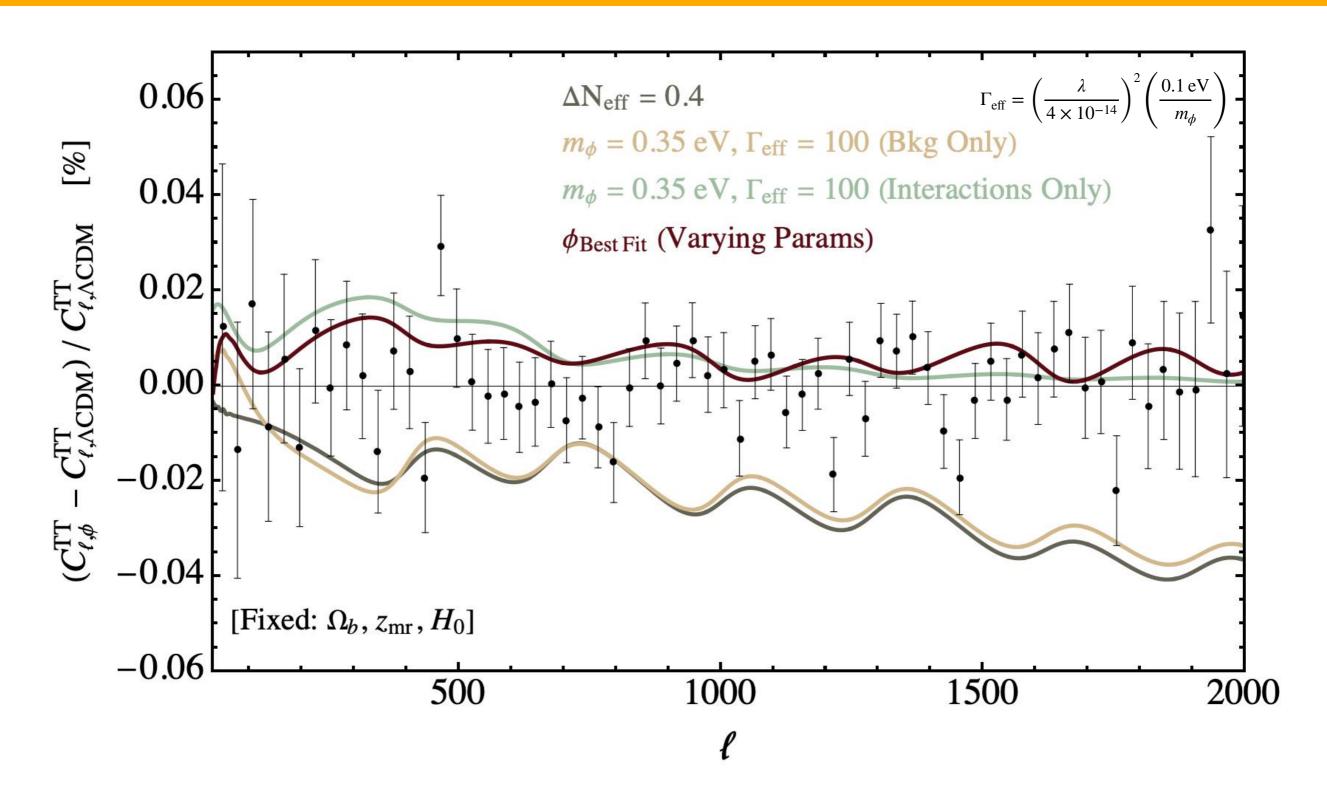
# Cosmological Implications



#### The effect is enhanced if there is a primordial population:



### Effect on the CMB



Neutrino-Majoron interactions can compensate the enhanced expansion history!

# Summary of Escudero & Witte 21'

- Sterile Neutrinos can provide just the right primordial majoron population
- A full Planck Legacy data analysis shows that:

for 
$$\Delta N_{\rm eff}^{\rm BBN} = 0.37$$
 
$$m_{\phi} = (0.1 - 0.8) \text{ eV}$$
 
$$v_L = (0.05 - 2) \text{ TeV}$$
 
$$H_0 = (70.2 \pm 0.6) \text{ km/s/Mpc}$$
 
$$M_N \sim \text{GeV}$$

This makes the tension  $4.2\sigma \rightarrow 2.0\sigma$  but with a better CMB fit than  $\Lambda$ CDM!

We argue that in the parameter space of interest these sterile neutrinos can lead to the baryon asymmetry of the Universe via their CP violating oscillations. ARS-Leptogenesis

Akhmedov, Rubakov & Smirnov, hep-ph/9803255
See also Asaka & Shaposhnikov, hep-ph/0505013

provided 
$$|\lambda_{\phi H}| < 10^{-7} \frac{v_L}{1\,\mathrm{TeV}} \sqrt{\frac{10^5\,\mathrm{GeV}}{T_c}}$$
 which requires some fine tuning but at least is protected under RGE flow

Our expectations have been confirmed by:

2109.10908 Flood, Porto, Schlesinger, Shuve, Thum 2110.14499 Fischer, Lindner, van der Woude

# Final Summary

#### The Majoron as a solution to the Hubble tension

- The Majoron represents can substantially relax the tension and can accommodate  $H_0 \simeq 70\,\mathrm{km/s/Mpc}$  while providing a good CMB fit.
- The Majoron is a well motivated particle:
  - It is predicted within the type-I seesaw with a global Lepton number symmetry
  - Its interaction rate with neutrinos is naturally very feebly
  - $\bullet$  The Majoron mass can be understood from Planck-scale Physics and points to  $m_{\phi} \sim {\rm eV}$
- lacktriangle Parameter space to solve the Hubble tension is very well motivated:  $v_L \sim v_H$
- The sterile neutrinos in the model play a crucial role:
  - By providing by their decays  $\Delta N_{\rm eff}^{\rm BBN} \sim 0.3$
  - In addition, they could be responsible for low-scale Leptogenesis

### Outlook

Collider tests  $K \rightarrow \mu N$  (NA62)

$$K 
ightarrow \mu N$$
 (NA62)

$$\pi \rightarrow e N$$
 (PIENU)

#### Cosmological tests

There are signals for ongoing/upcoming CMB experiments:

**ACT, SPT, Simons Observatory and CMB-S4** 

#### My main conclusion:

Regardless of what happens with the Hubble tension, we will learn about fundamental physics! Probing a well motivated neutrino mass model with  $\Lambda \sim 1 \, {\rm TeV!}$ 

# Acknowledgements

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### Time for Questions and Comments

#### Thank you for your attention!

